

Supercritical Impurities: Atomic Collapse in Graphene



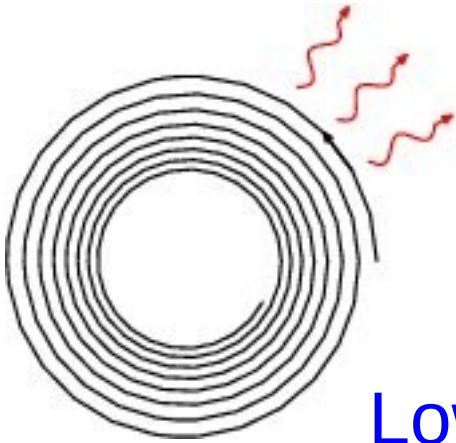
Leonid Levitov

(picture courtesy Mike Crommie)

Stability of a planetary atom

Classical physics: unstable
(energy unbounded)

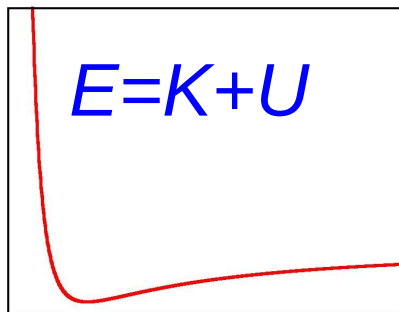
Quantum theory (Bohr, 1913):
stable orbits, Rydberg's formula



$$E_n = -\frac{me^4 Z^2}{2\hbar^2 n^2}$$



Lower bound: $E_1 = -\frac{me^4 Z^2}{2\hbar^2}$



Heisenberg (1926): uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

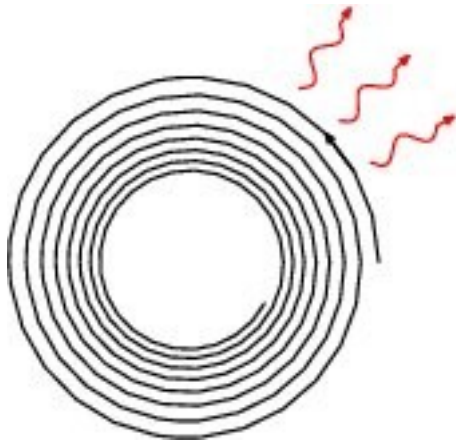
$$K_{\text{nr}} = \frac{p^2}{2m} \sim \frac{\hbar^2}{2mr^2} \gg U = -\frac{Ze^2}{r}$$



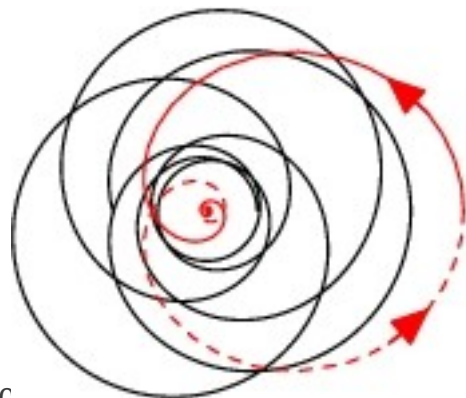
07/21/2013 **Planetary atom stabilized by QM (zero-point motion)**

Stability of a Relativistic Atom

Classical physics: unstable
(energy unbounded)



Relativity: collapsing orbits



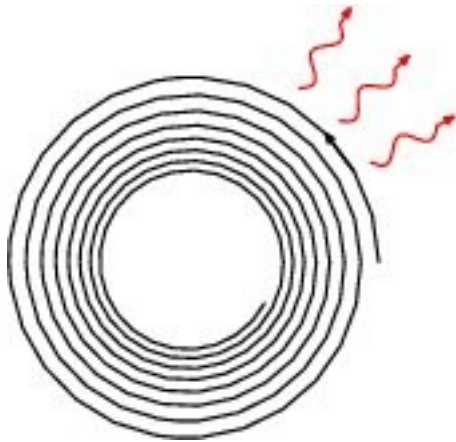
$$v < c$$

$$|Mc| > |Ze^2|$$

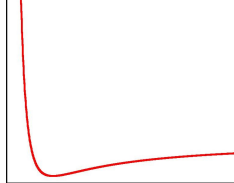
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Stability of a Relativistic Atom

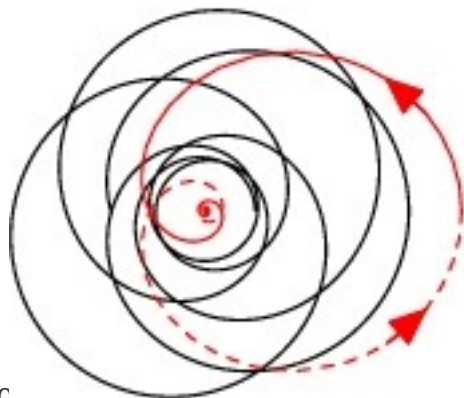
Classical physics: unstable
(energy unbounded)



QM orbitals stabilized by
zero point motion

$$K_{\text{nr}} = \frac{p^2}{2m} \sim \frac{\hbar^2}{2mr^2} \quad E=K+U$$
$$U = -\frac{Ze^2}{r}$$
A graph showing the potential energy U as a function of distance r. The curve starts at a very high negative value for small r and levels off to a constant negative value as r increases, representing the Coulomb potential.

Relativity: collapsing orbits



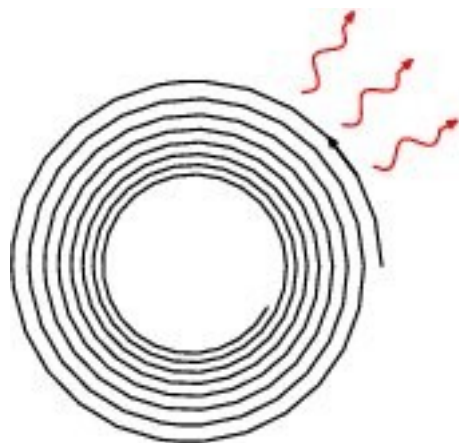
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Stability of a Relativistic Atom

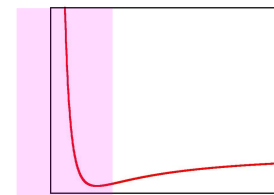
Classical physics: unstable
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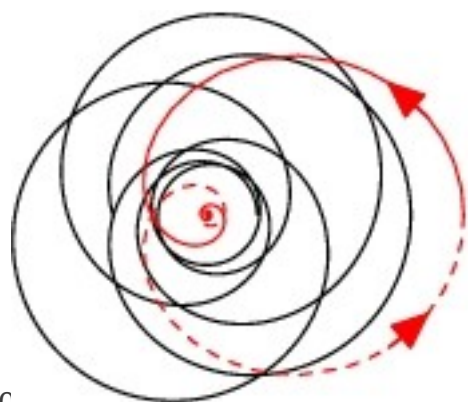
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Relativity: collapsing orbits



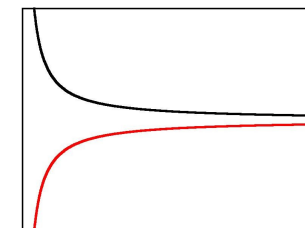
$$v < c$$

$$|Mc| > |Ze^2|$$

$$|Mc| < |Ze^2|$$

Relativity + QM:

$$K = cp \sim \frac{\hbar c}{r} \stackrel{?}{>} U$$



Collapse?

Dirac atoms can implode:

Subcritical ($Z < 137$)

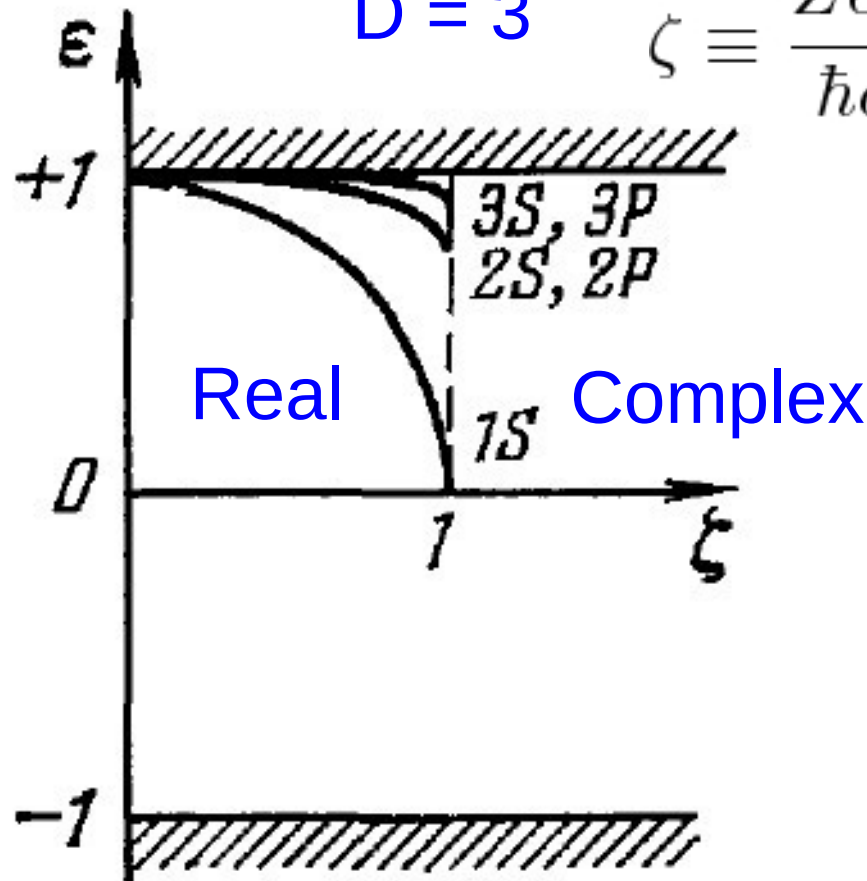


$m \neq 0$ $\epsilon_1 = m\sqrt{1 - \zeta^2}$
 $D = 3$ $\zeta \equiv \frac{Ze^2}{\hbar c}$

Dirac (1929)

Complex energies at

$\zeta > 1$



Dirac atoms can implode:

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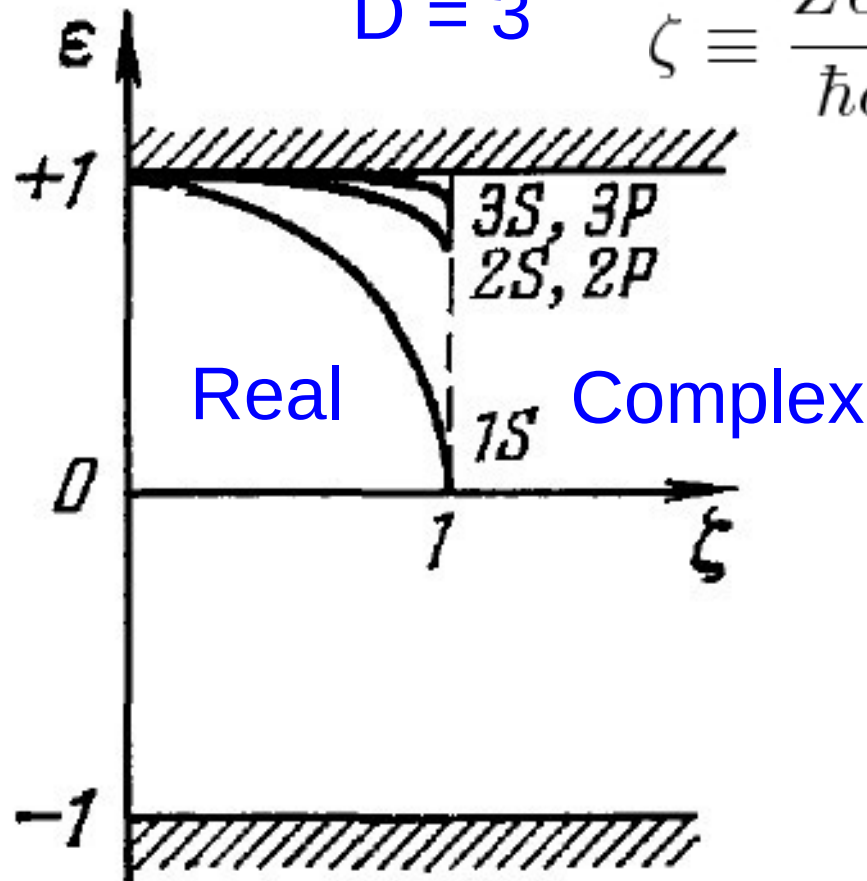


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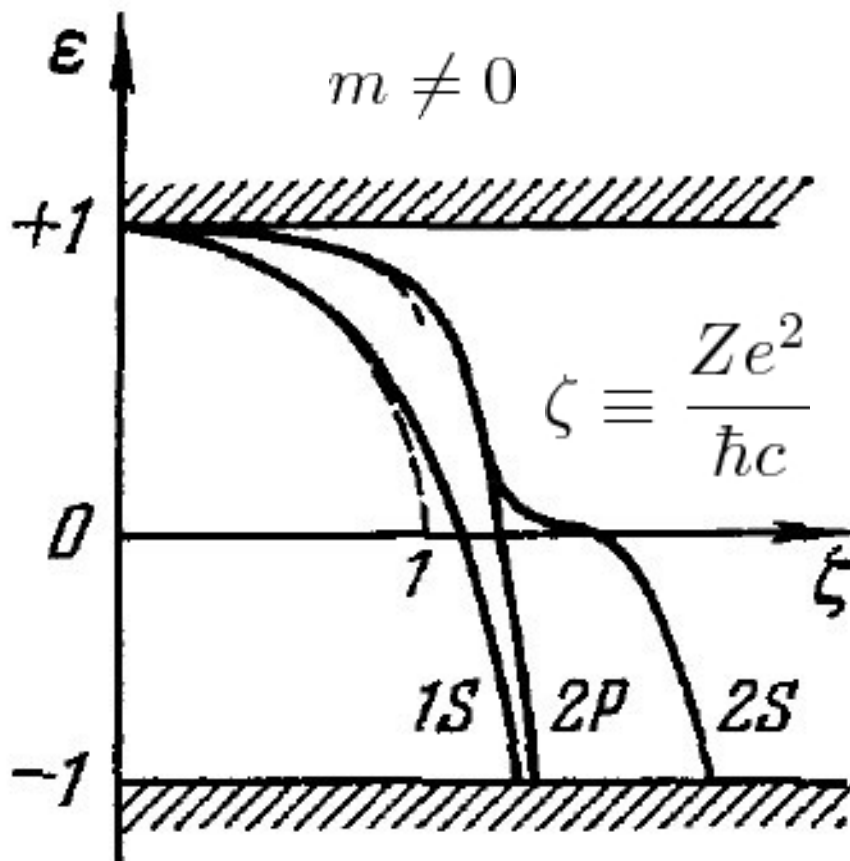
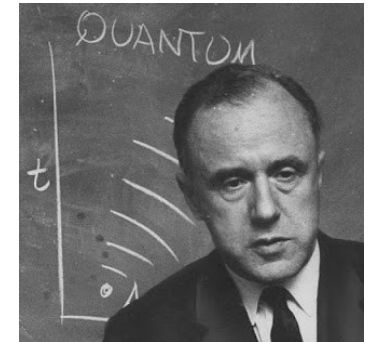


What happens at $Z > 137$?

Supercritical atom

Pre-collapse ($137 < Z < 170$)

*Pomeranchuk & Smorodinskii (1945);
Werner and Wheeler (1957)*



Finite size of nucleus
Pomeranchuk nuclear
formfactor

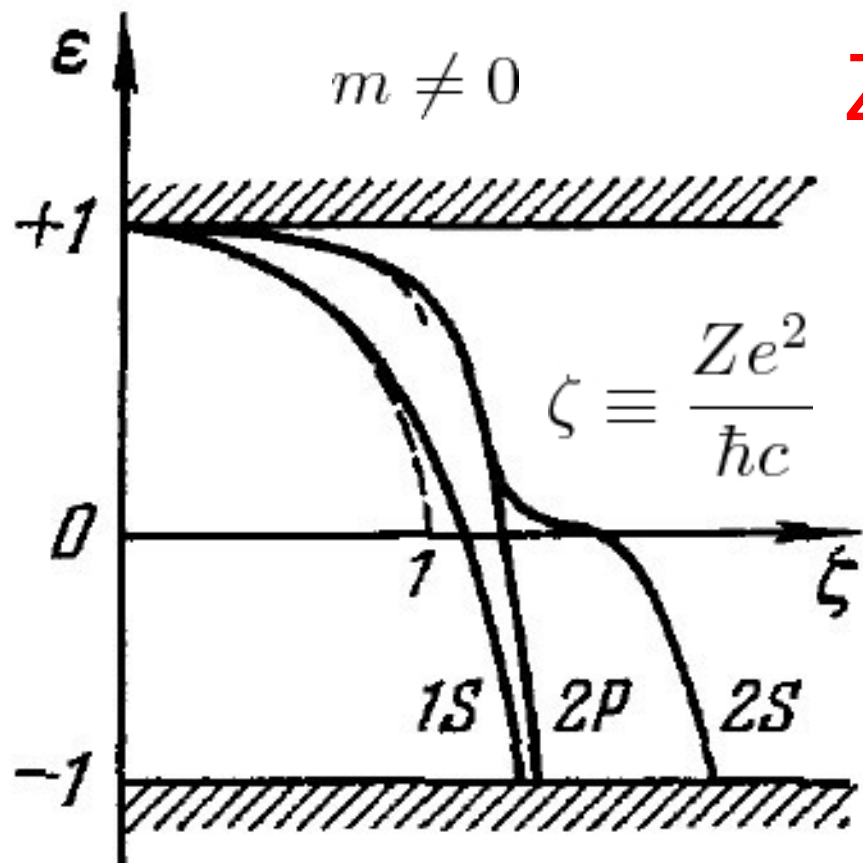
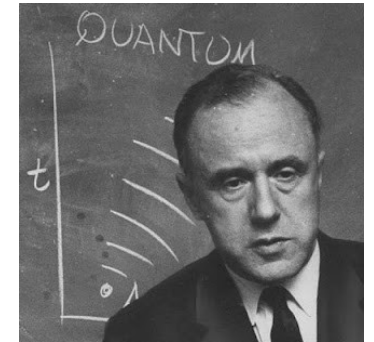
$$r_0 \approx 1.2 \cdot 10^{-12} \text{ cm}$$

1S level dives into
Dirac sea at $Z = 170$

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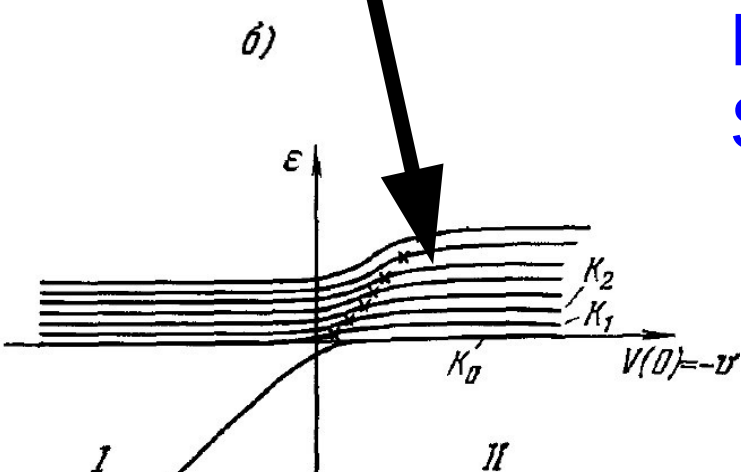
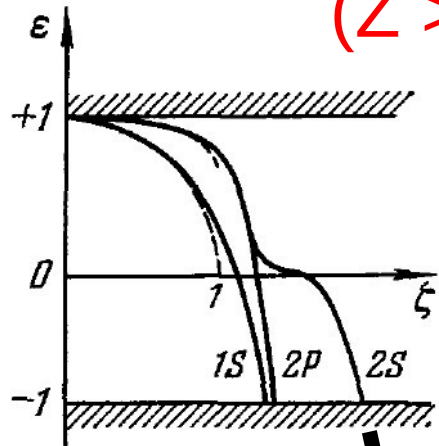
Supercritical atom

Collapse, vacuum reconstruction
($Z > Z_c = 170$)



Gershteyn, Zeldovich (1969)
Popov (1970)

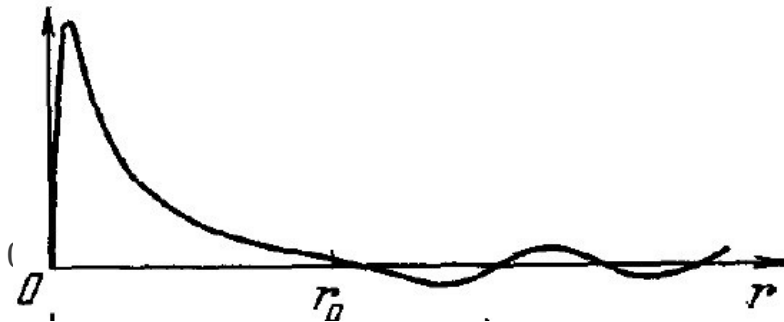
Resonance state in the Dirac sea
Screening by pair production?



$$\epsilon = \epsilon_0 - i\gamma \quad \epsilon_0 = -m - a(Z - Z_c)$$

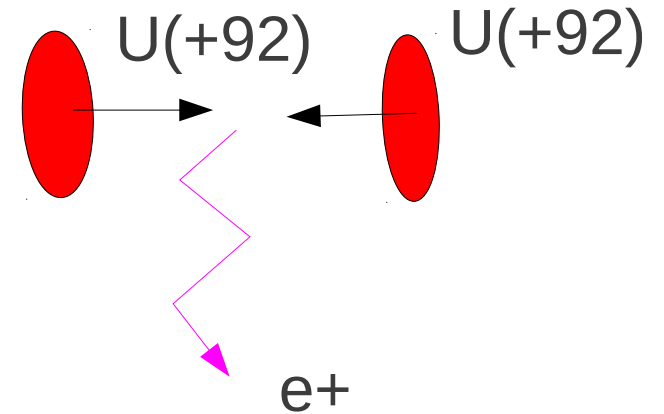
$$\gamma \sim \exp\left(-\frac{b}{\sqrt{Z - Z_c}}\right)$$

Quasilocalized spatial structure
of the resonance state



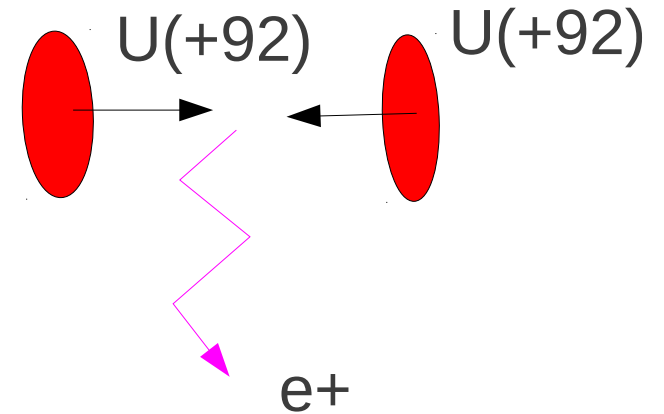
Signatures: supercritical e^+ emission

Darmstadt experiment (1980s,
revisited in 1990s)
UNILAC (EPOS, ORANGE)
3-6 MeV collisions



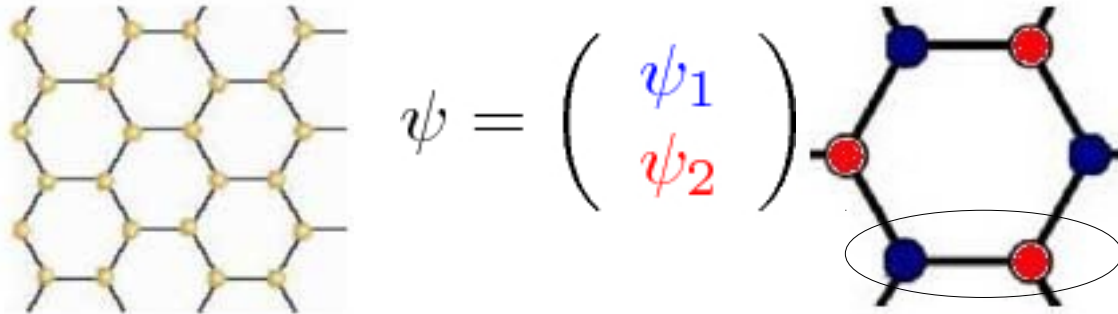
Signatures: supercritical e^+ emission

Darmstadt experiment (1980s,
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UNILAC (EPOS, ORANGE)
3-6 MeV collisions
No signatures of supercritical emission



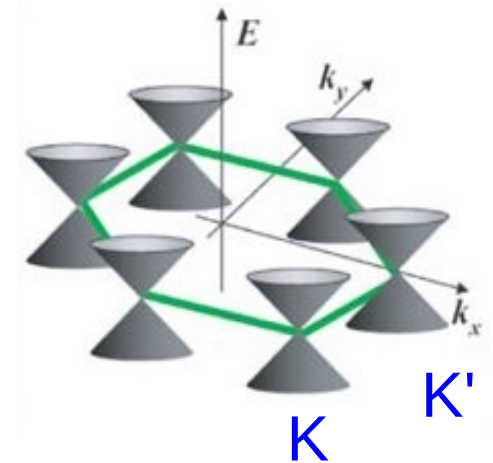
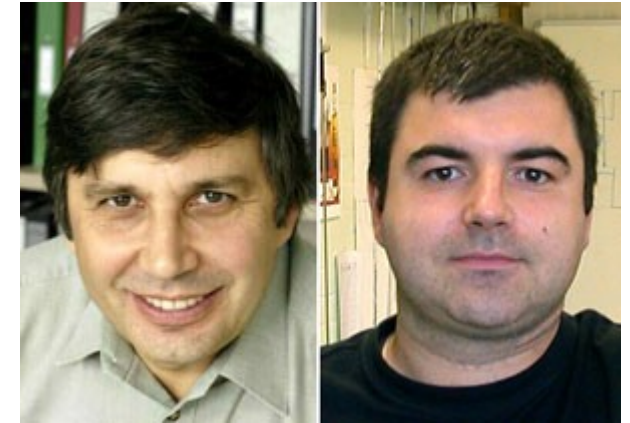
Relativistic massless electrons in graphene

Two sublattices



pseudo-spin (sublattice)

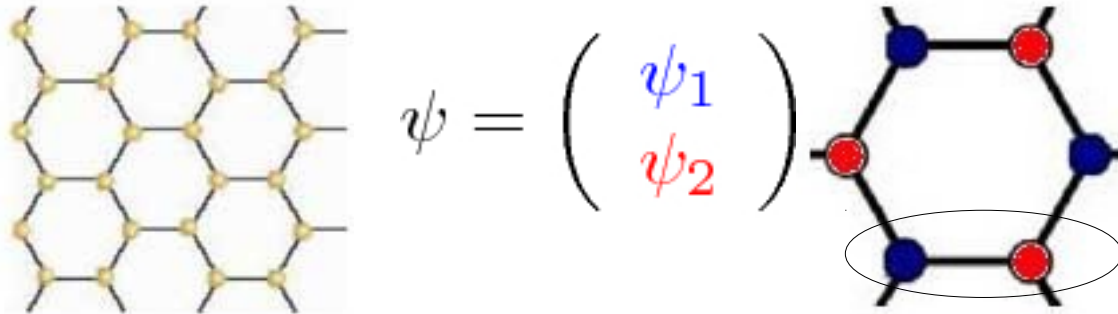
$$\hat{H} = v_F \hat{\sigma} \mathbf{p}$$



4-fold
degeneracy
spin&valley

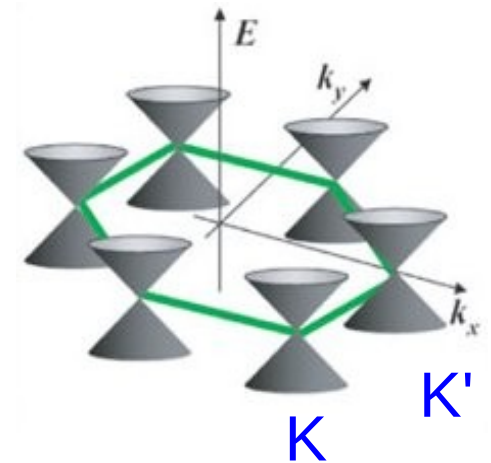
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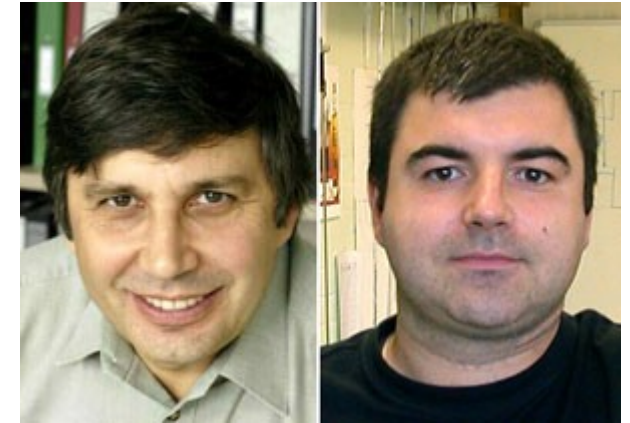
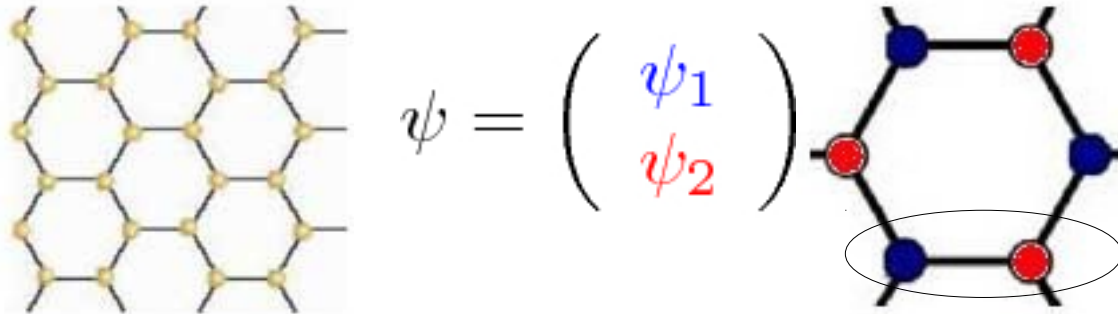
$$\hat{H} = v_F \hat{\sigma} \mathbf{p} \quad E = \sqrt{c^2 p^2 + m^2 c^4}$$



4-fold
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Relativistic massless electrons in graphene

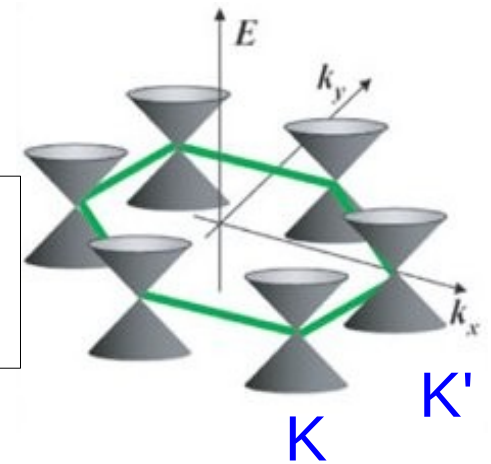
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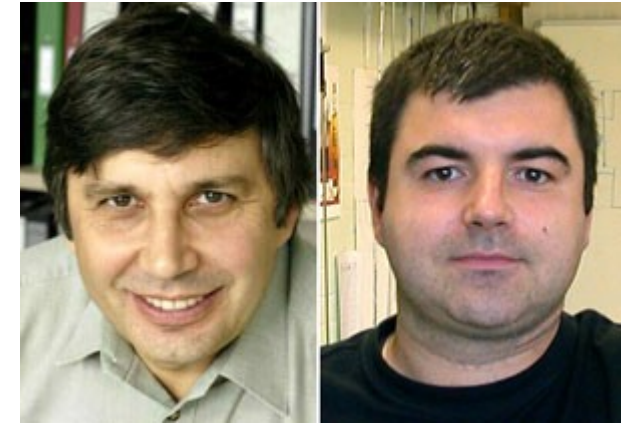
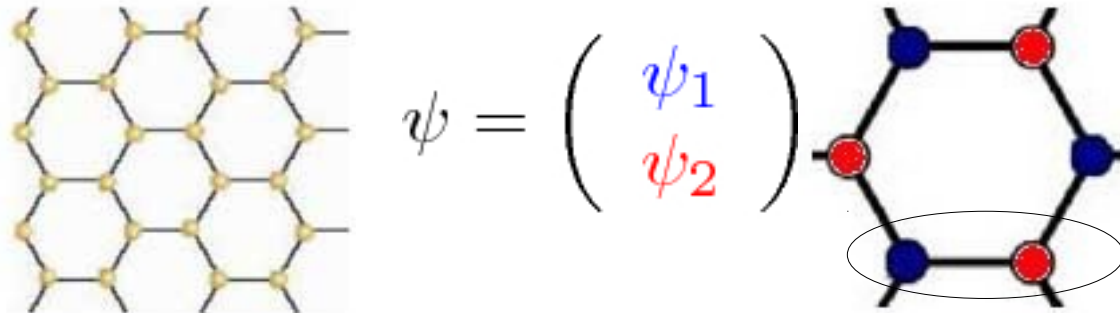
No gap (cond-mat) \equiv No mass (hep)



4-fold
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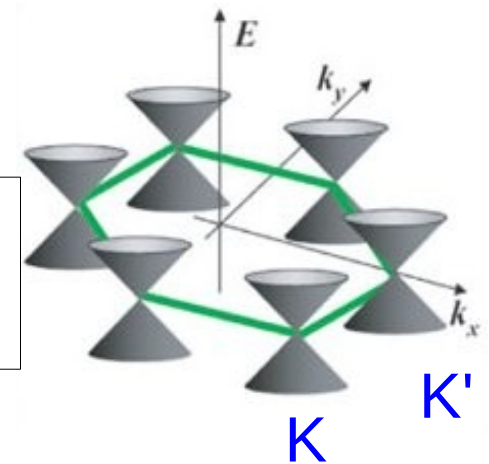
pseudo-spin (sublattice)

$$\hat{H} = v_F \hat{\sigma} \mathbf{p} \quad E = \sqrt{c^2 p^2}$$

No gap (cond-mat) \equiv No mass (hep)

$$v_F \approx 10^6 \text{ m/s} = \frac{c}{300} \quad \alpha = e^2 / \hbar v \approx 2.5$$

Slow, but ultra-relativistic
Dirac fermions



4-fold
degeneracy
spin&valley

Atomic collapse for massless fermions

- Large fine structure constant, $\alpha = \frac{e^2}{\hbar v} \approx 2.5$, low collapse threshold $Z \sim 1$
- Massless Dirac equation: continuum spectrum only, no discrete spectrum
- Manifestation of collapse: formation of resonances (infinite family, quasilocalized spatial structure)
- Strong effects in vacuum polarization: screening cloud may extend indefinitely (*cf.* cutoff at Compton wavelength for massive Dirac electron)

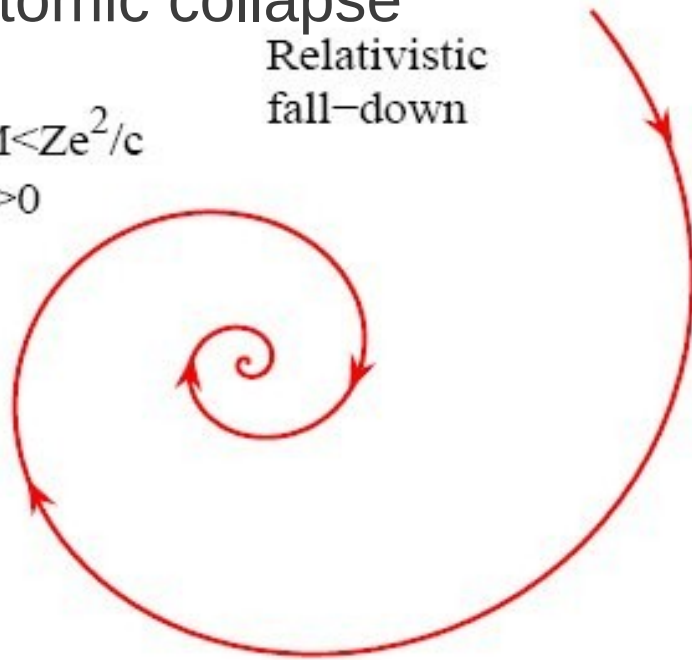
A. V. Shytov, M. I. Katsnelson, and L. S. Levitov,
PRL 99, 236801 (2007), PRL 99, 246802 (2007)

Quasistationary states

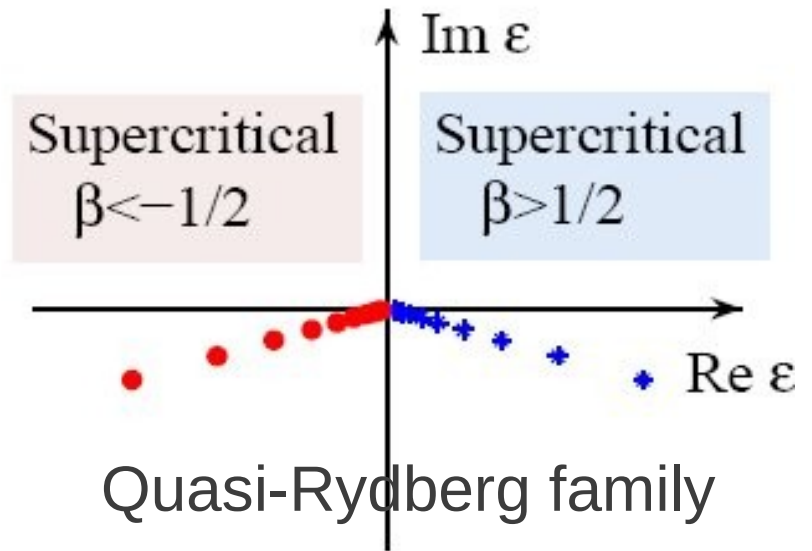
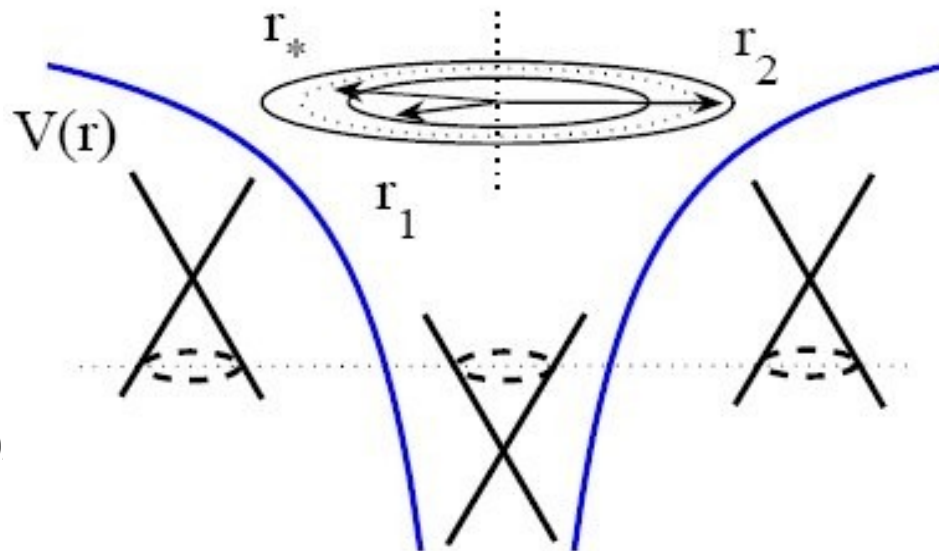
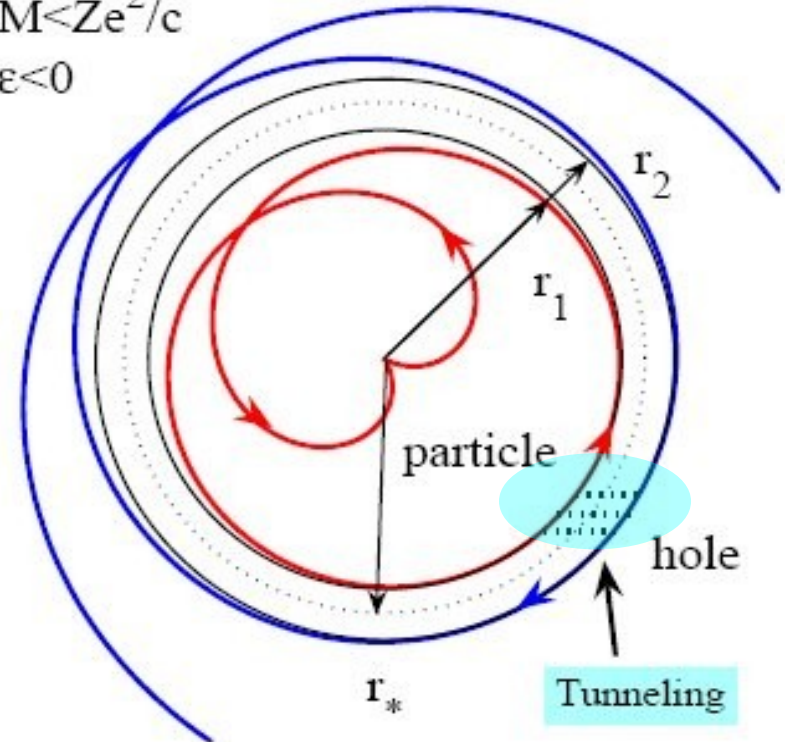
Atomic collapse

Relativistic
fall-down

$M < Ze^2/c$
 $\epsilon > 0$

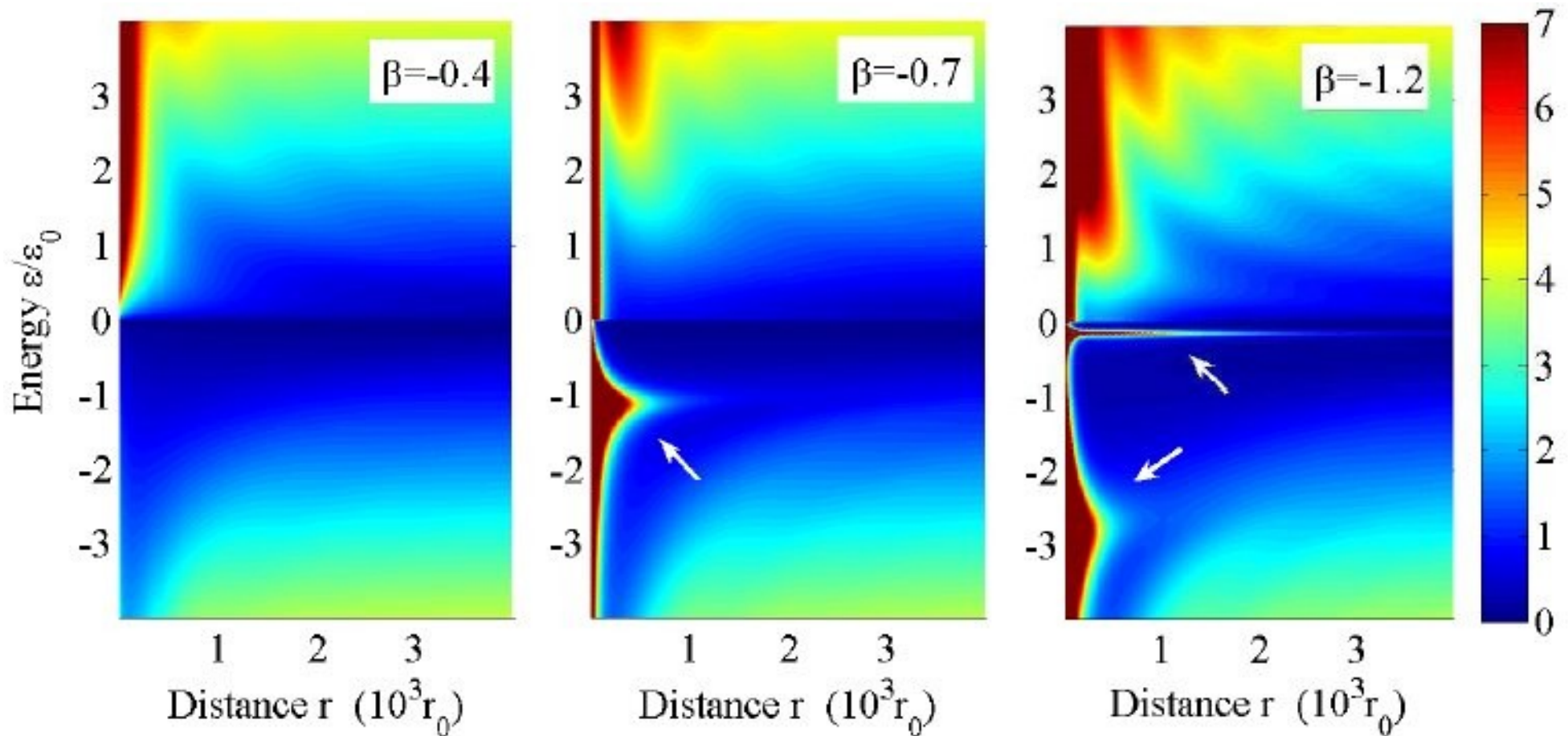


$M < Ze^2/c$
 $\epsilon < 0$



Resonances in the local density of states, can be probed by STM

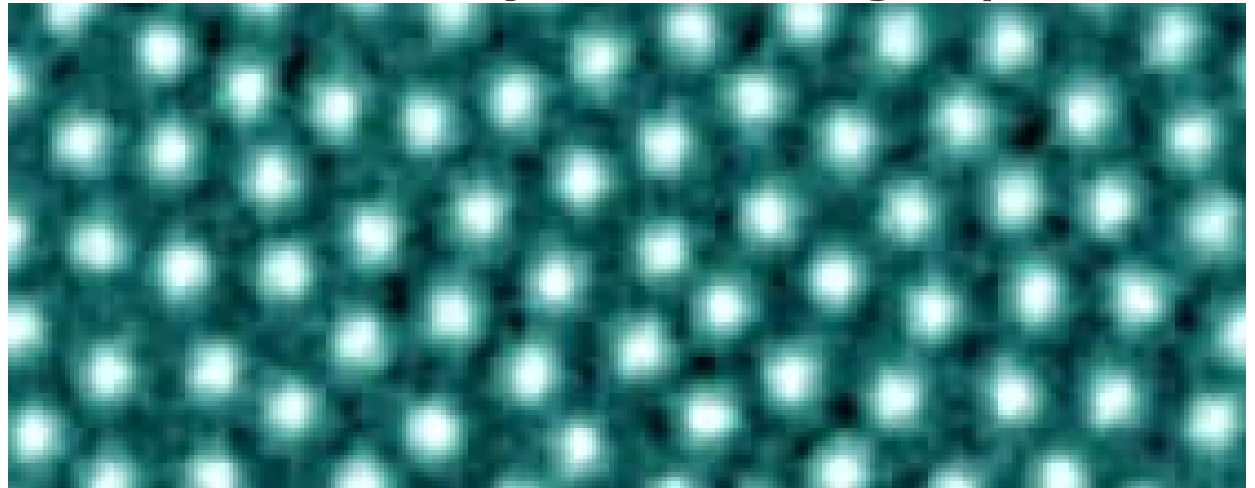
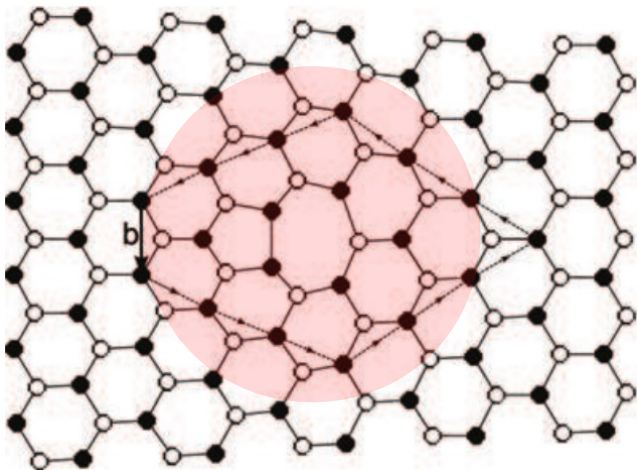
Tunneling spectroscopy $\beta = Ze^2 / \kappa \hbar v$



09/20/2012 Energy scales as the width Γ and as $1/(\text{localization radius})$

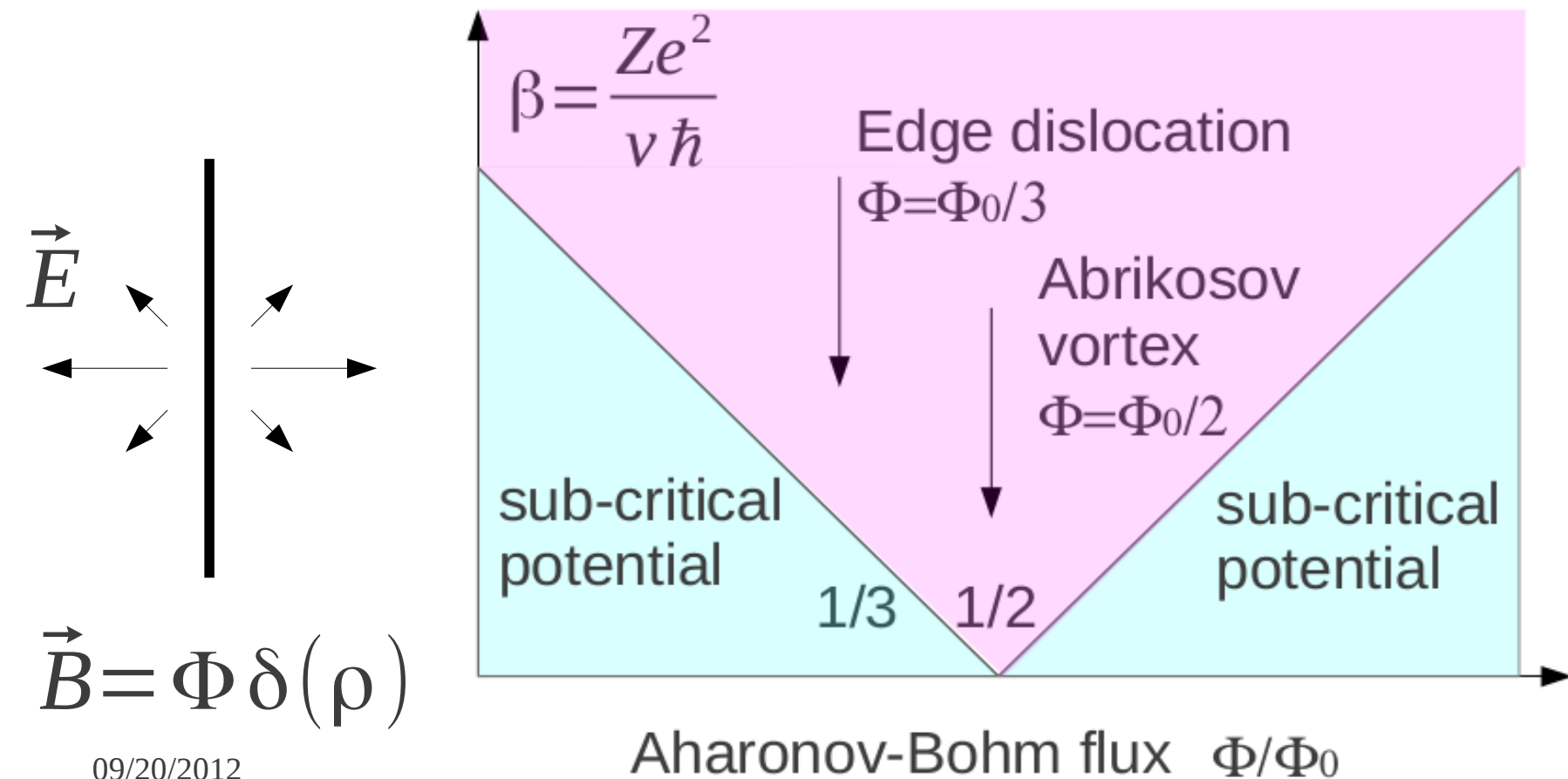
Reaching collapse (critical Z values)

- Vary Z. Critical value affected by intrinsic screening. Accounting for self-consistent nonlinear screening challenging. Optimistic estimates yield $1 < Z < 2$.
- Experiment (Crommie): Co trimers, Ca dimers
- Recent work: AB flux as a vehicle to control collapse, $Z \ll 1$
- Can be realized for dislocations (pseudo-B field), or vortices in a superconductor adjacent to graphene



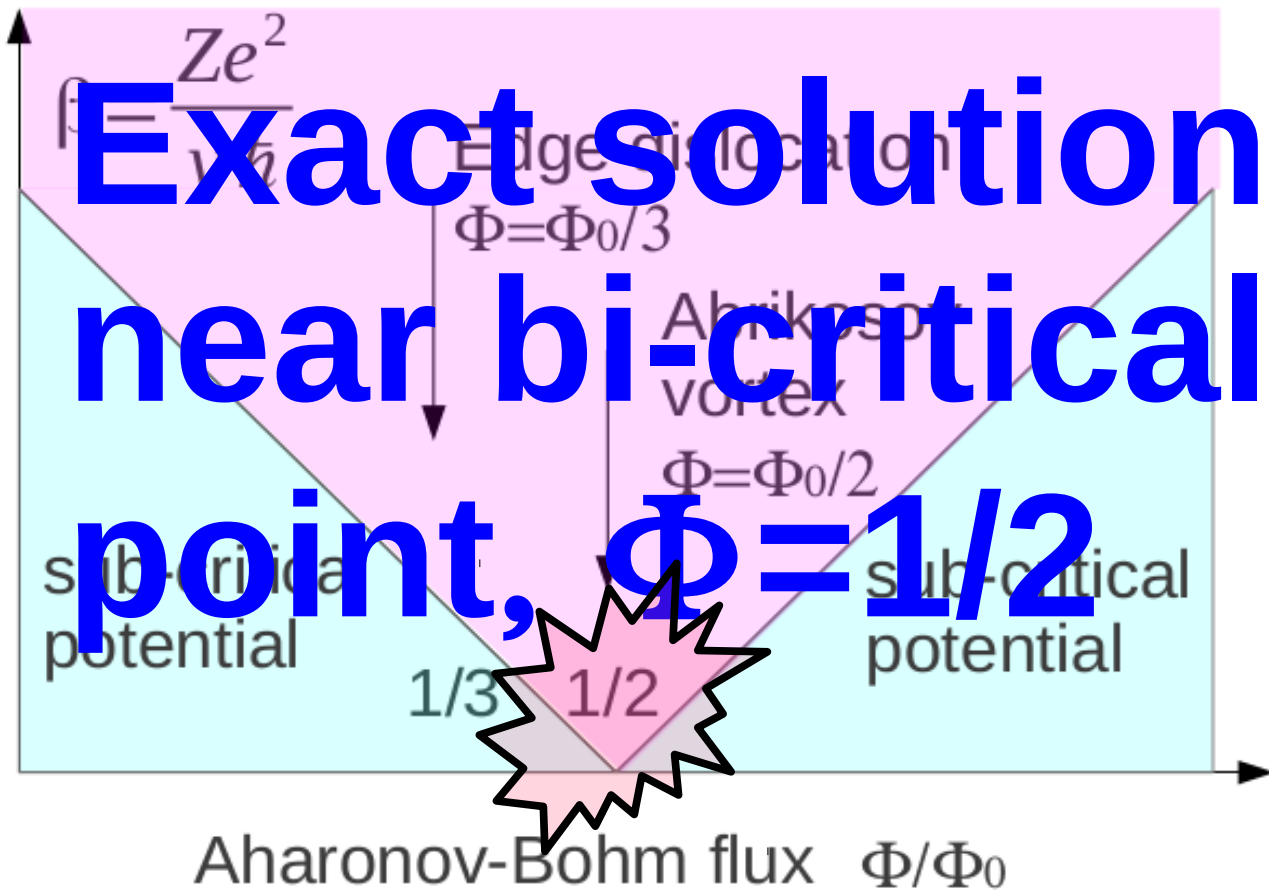
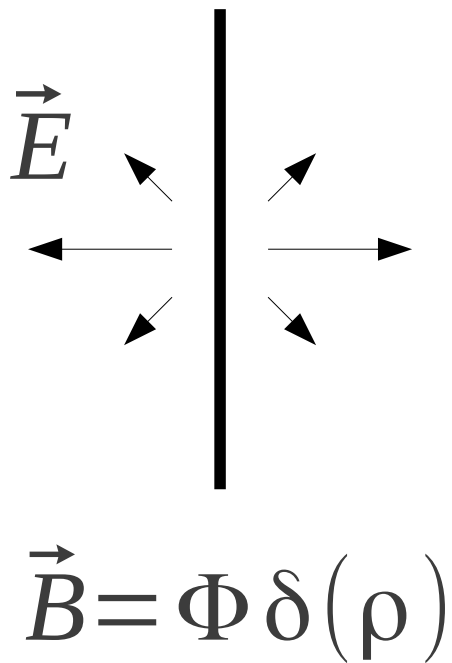
Atomic collapse in the presence of an Aharonov-Bohm solenoid

$$H = v \sigma (\vec{p} - \vec{A}(\rho)) - \frac{Z e^2}{\rho} \quad \vec{A}(\rho) = \frac{\Phi}{2\pi\rho} (-y, x)$$



Atomic collapse in the presence of an Aharonov-Bohm solenoid

$$H = v \sigma (\vec{p} - \vec{A}(\rho)) - \frac{Z e^2}{\rho} \quad \vec{A}(\rho) = \frac{\Phi}{2\pi\rho} (-y, x)$$



Exact solution near bi-critical point, $\Phi = 1/2$

Screening cloud: transition between dielectric and metallic behavior

- Subcritical regime: dielectric behavior, $1/r$ potential, screening charge concentrated on the lattice scale
- Supercritical regime: metallic behavior, large-scale screening cloud with a power law tail
- Predict scaling exponent value from exact solution

$$V(\rho \gg a) \sim \rho^{-\eta} \quad 1 < \eta < 2$$

Take-home message

- AB flux a knob to tune collapse (in addition to Z)
- A bi-critical point, two regimes: “dielectric” and “metallic”
- Exact solution for vacuum polarization and the screening cloud structure
- Experimental search for supercritical potential.

Take-home message

- AB flux a knob to tune collapse (in addition to Z)
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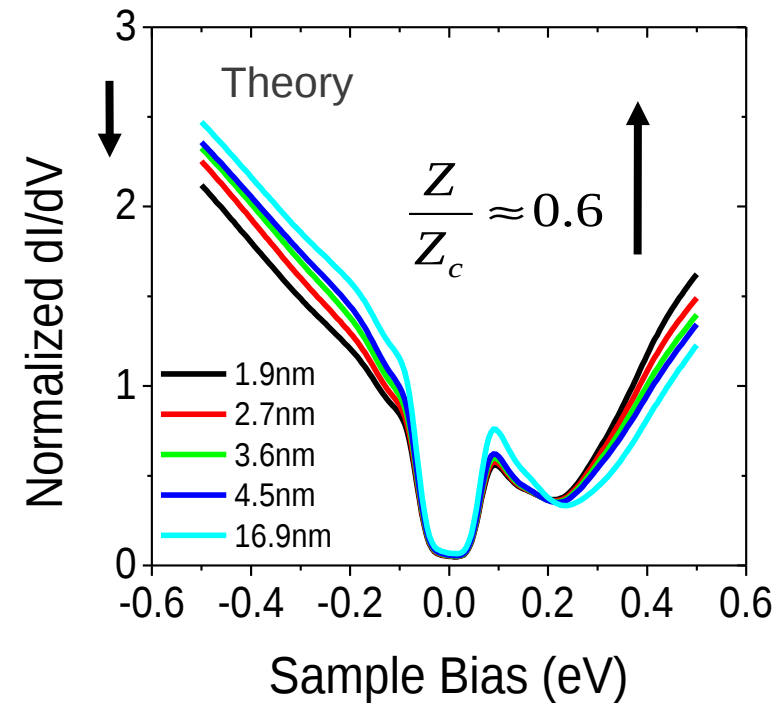
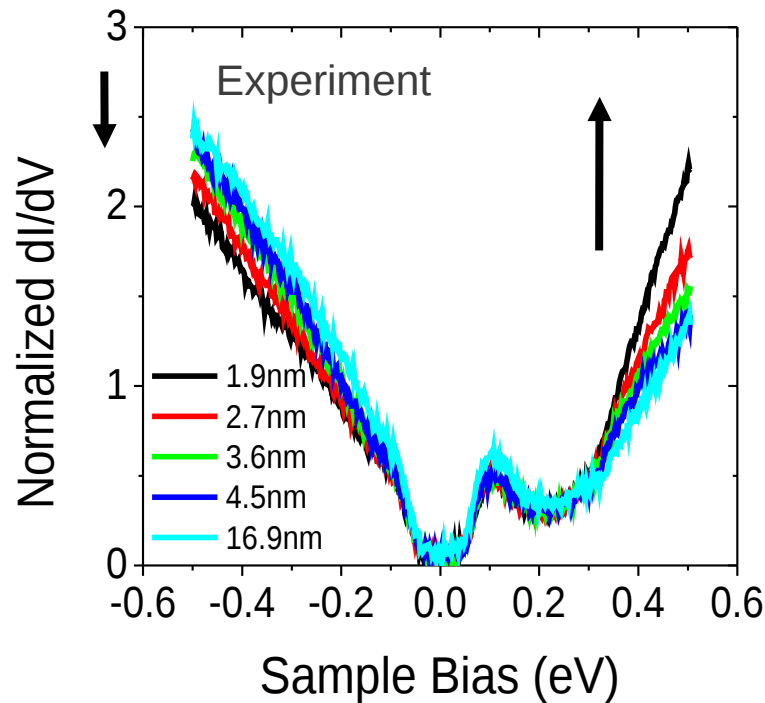
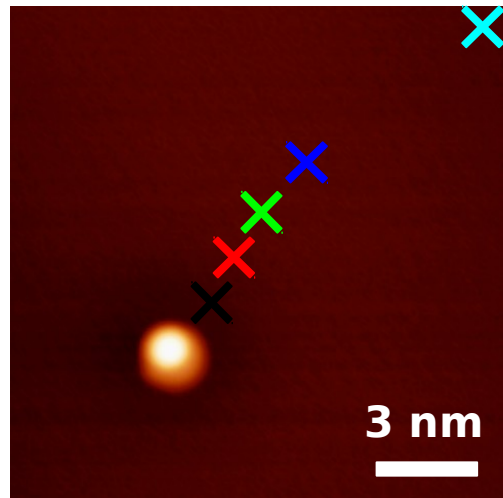
Just found it

How do we make a supercritical potential? Ca dimers on G/BN, Crommie's Lab (Berkeley)

Ca dimer

STS for 1 Dimer

Simulation for 1 Dimer



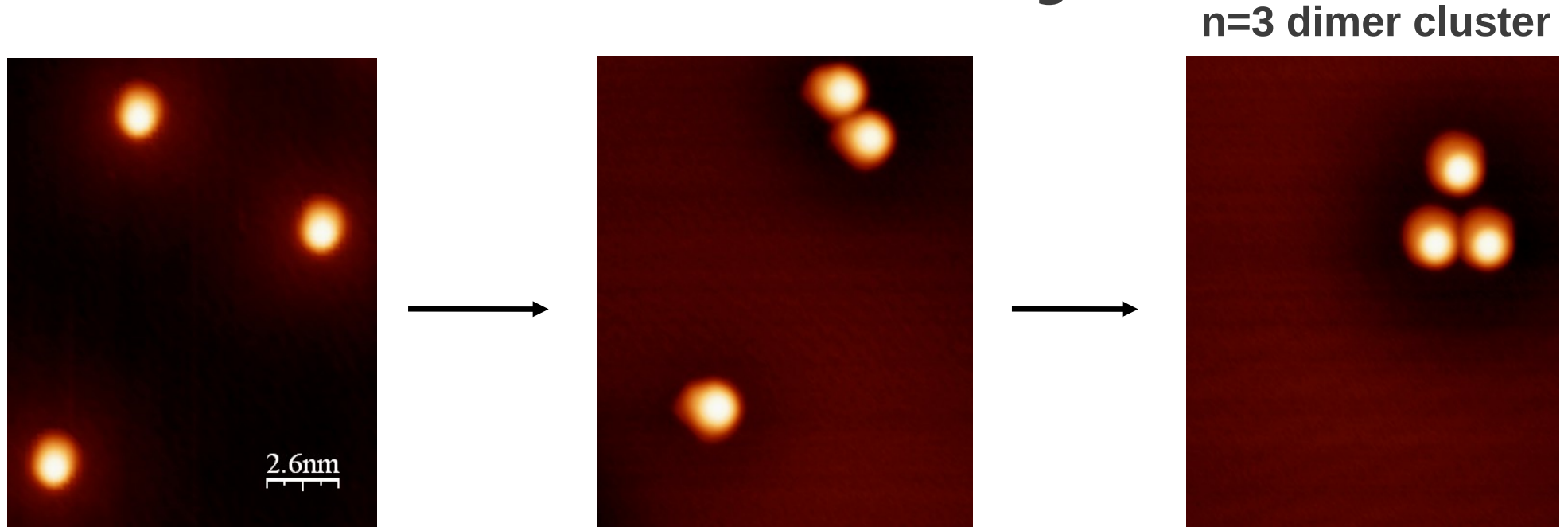
$$Z < Z_c$$

Subcritical...

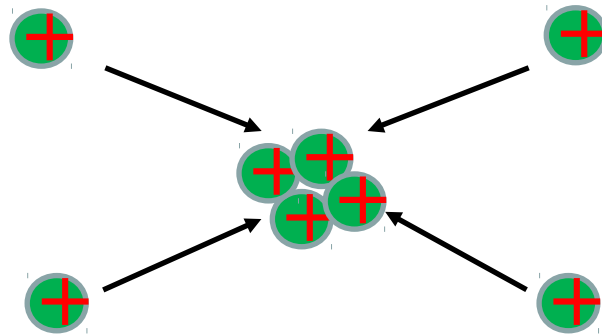
Previous work (probed intrinsic screening for Co trimers)
Wang, Brar, Shytov, Wu, Regan, Tsai, Zettl, Levitov, Crommie,
Nat Phys 8, 653 (2012)

Ca Dimer Has an Advantage: Can Manipulate it

Ca Dimers are Moveable Charge Centers

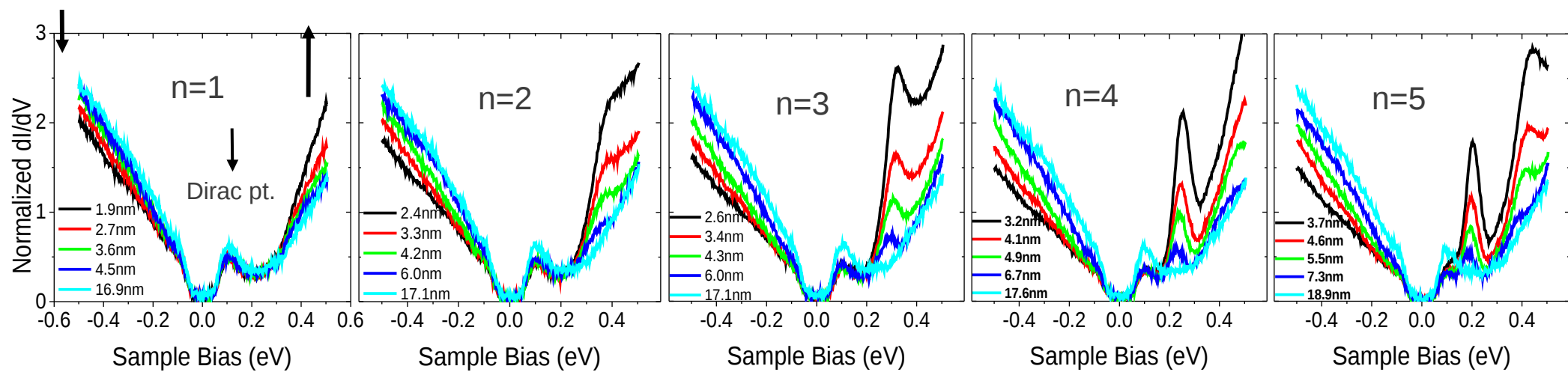
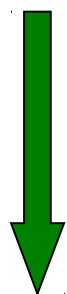
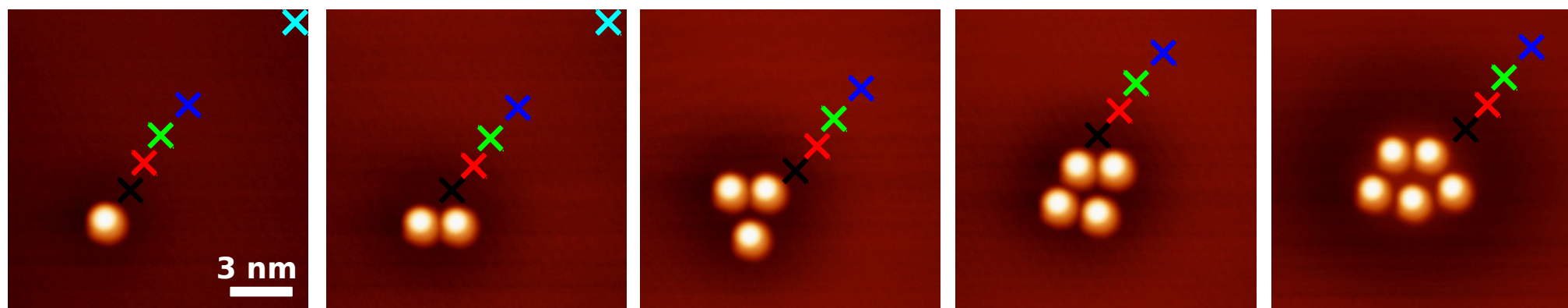


Fabricate Artificial Nuclei

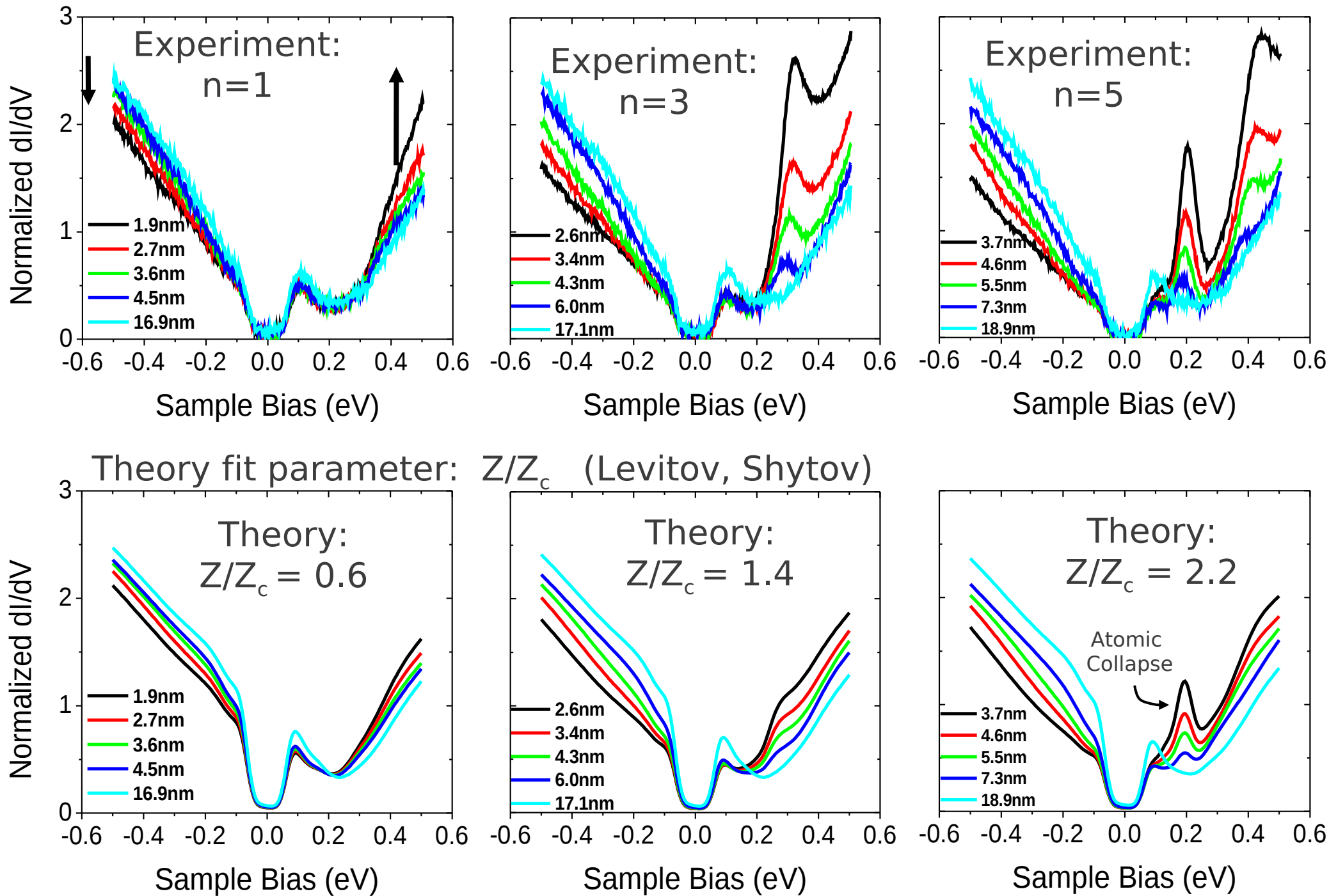



Tune Z above Z_c

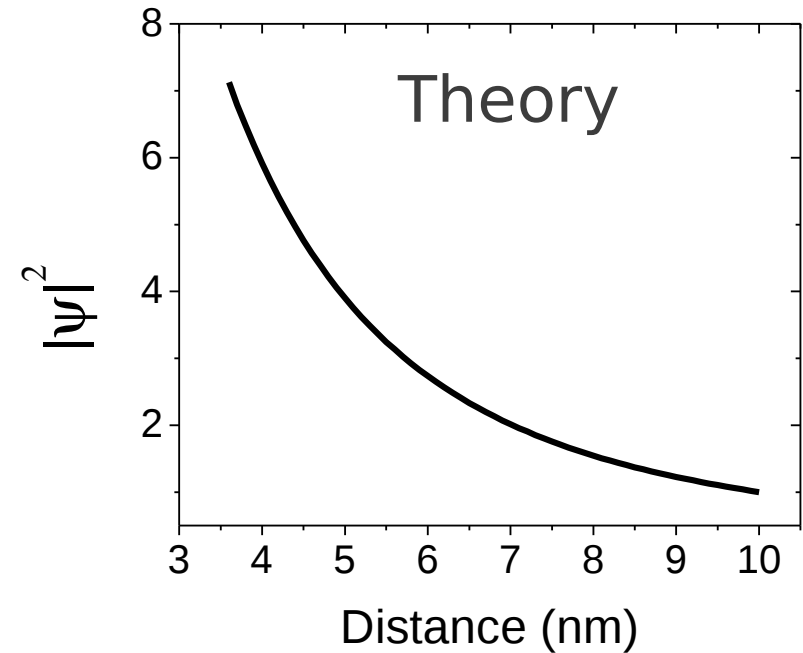
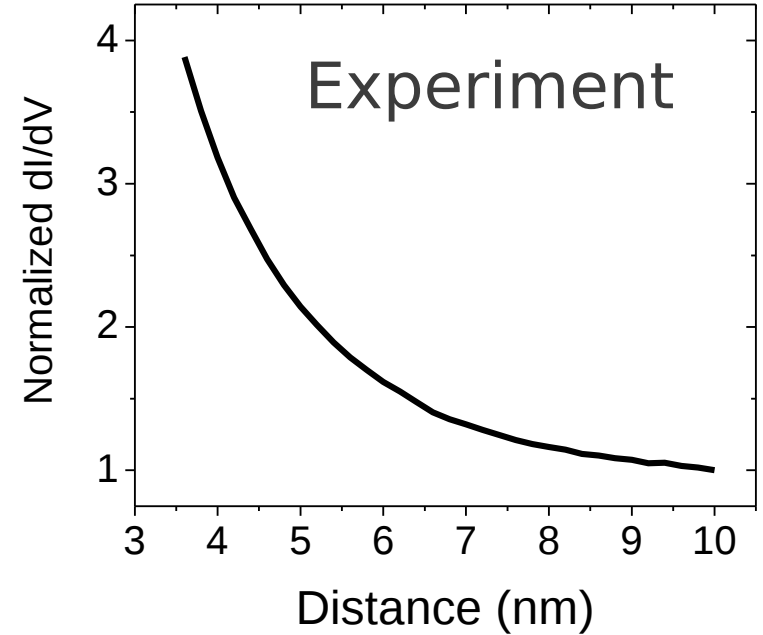
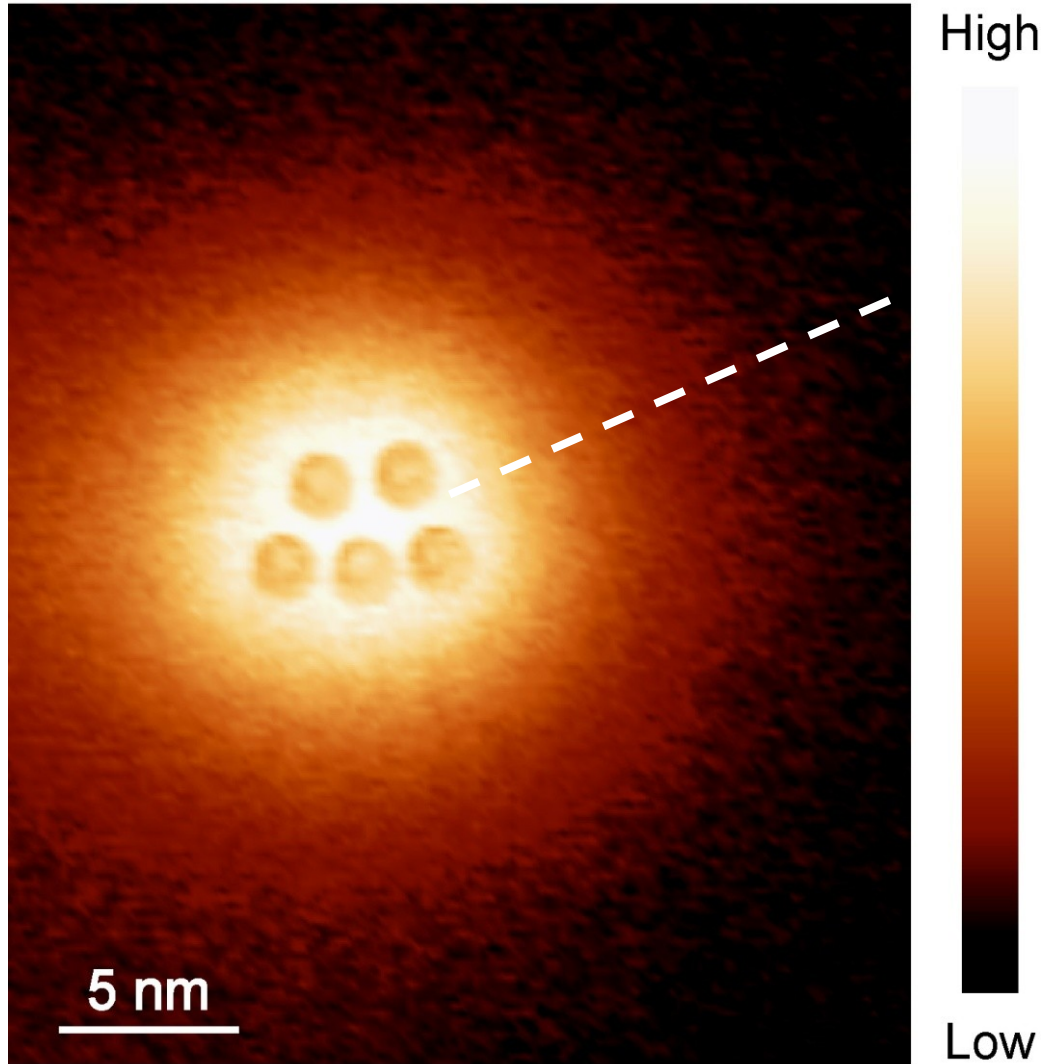
Tuning Z by Building Artificial Nuclei from Ca Dimers



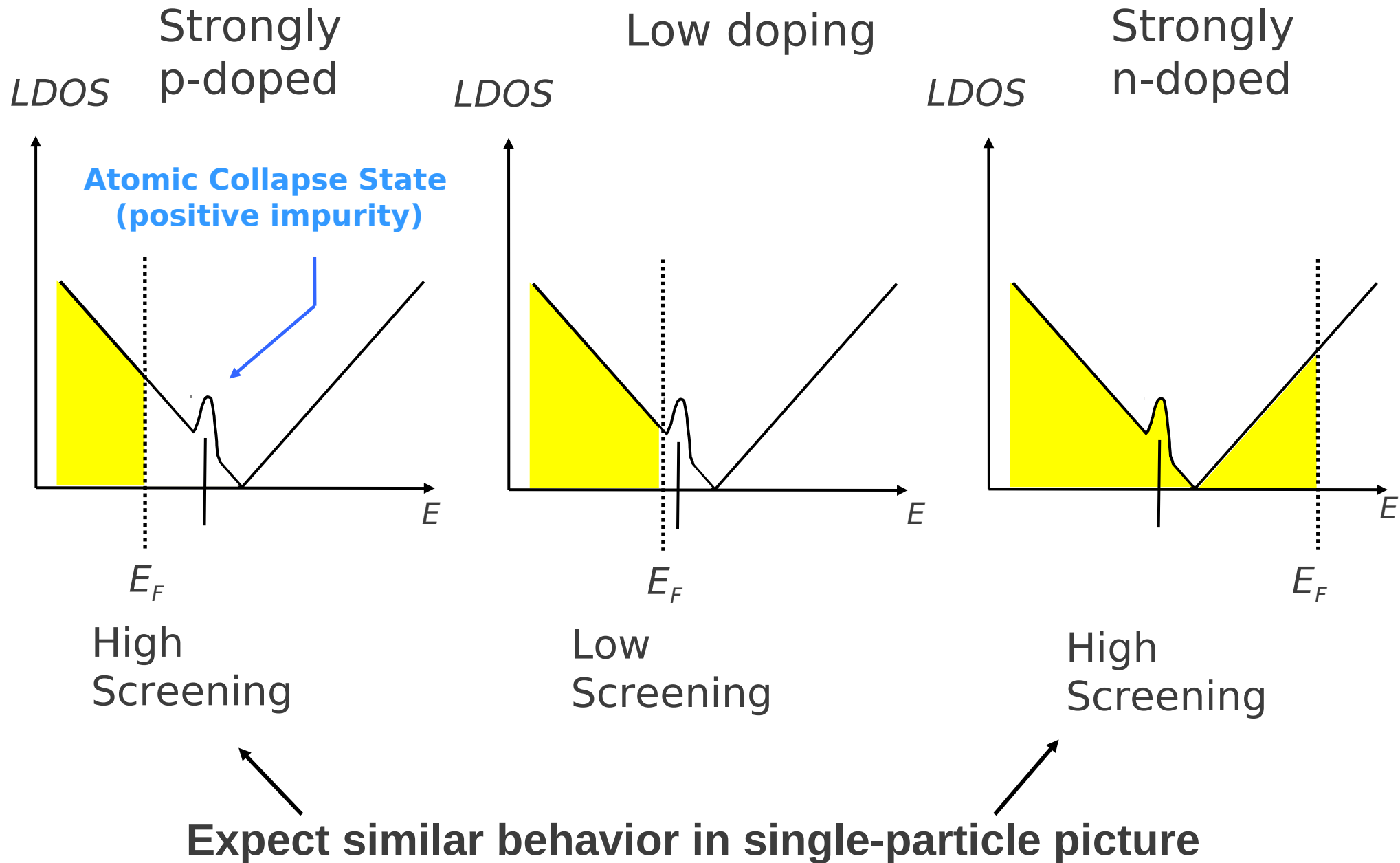
Compare to Simulated Behavior



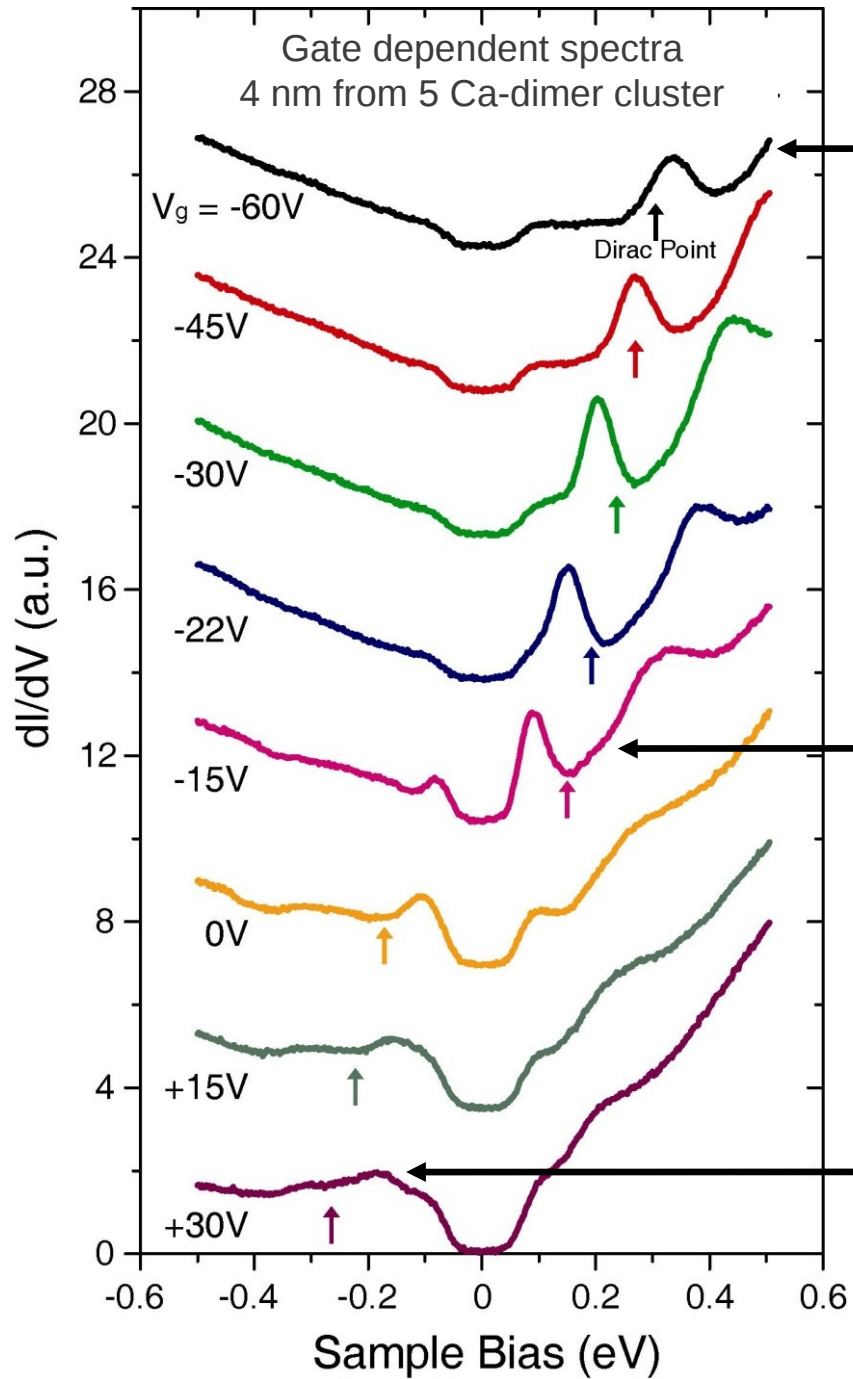
Atomic Collapse Spatial Dependence n=5 Cluster



Dependence on Electron Occupation



Observed Density Dependence NOT Symmetric



High p-doped
 $n \sim 3 \times 10^{12} \text{ cm}^{-2}$

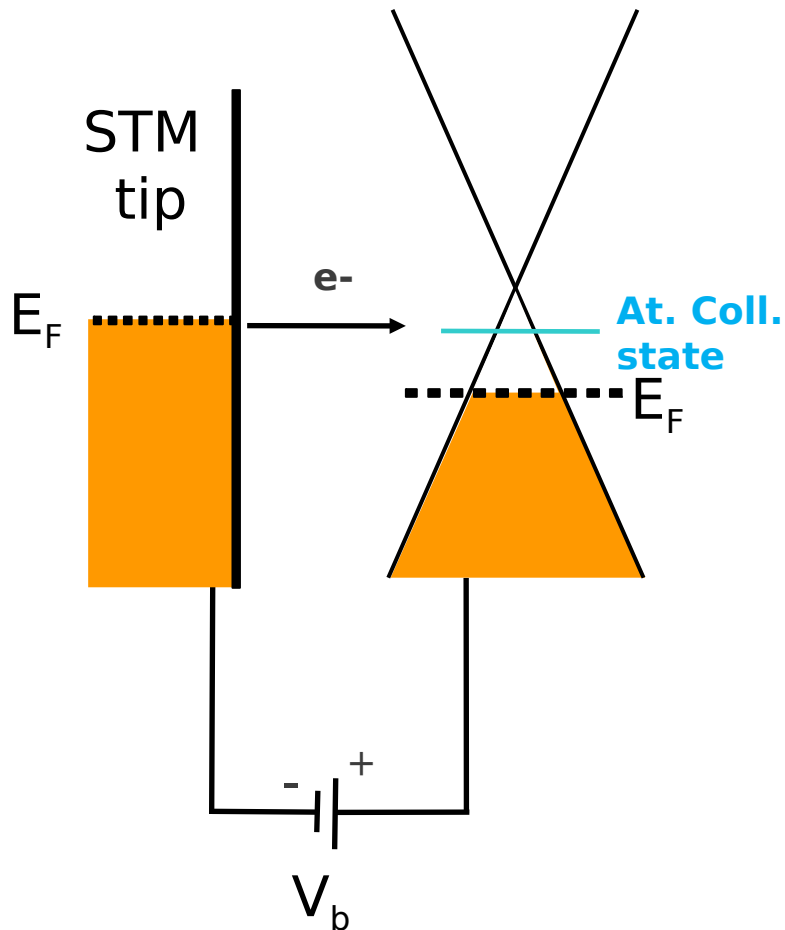
Low p-doped
 $n \sim 4 \times 10^{11} \text{ cm}^{-2}$

High n-doped
 $n \sim 3 \times 10^{12} \text{ cm}^{-2}$

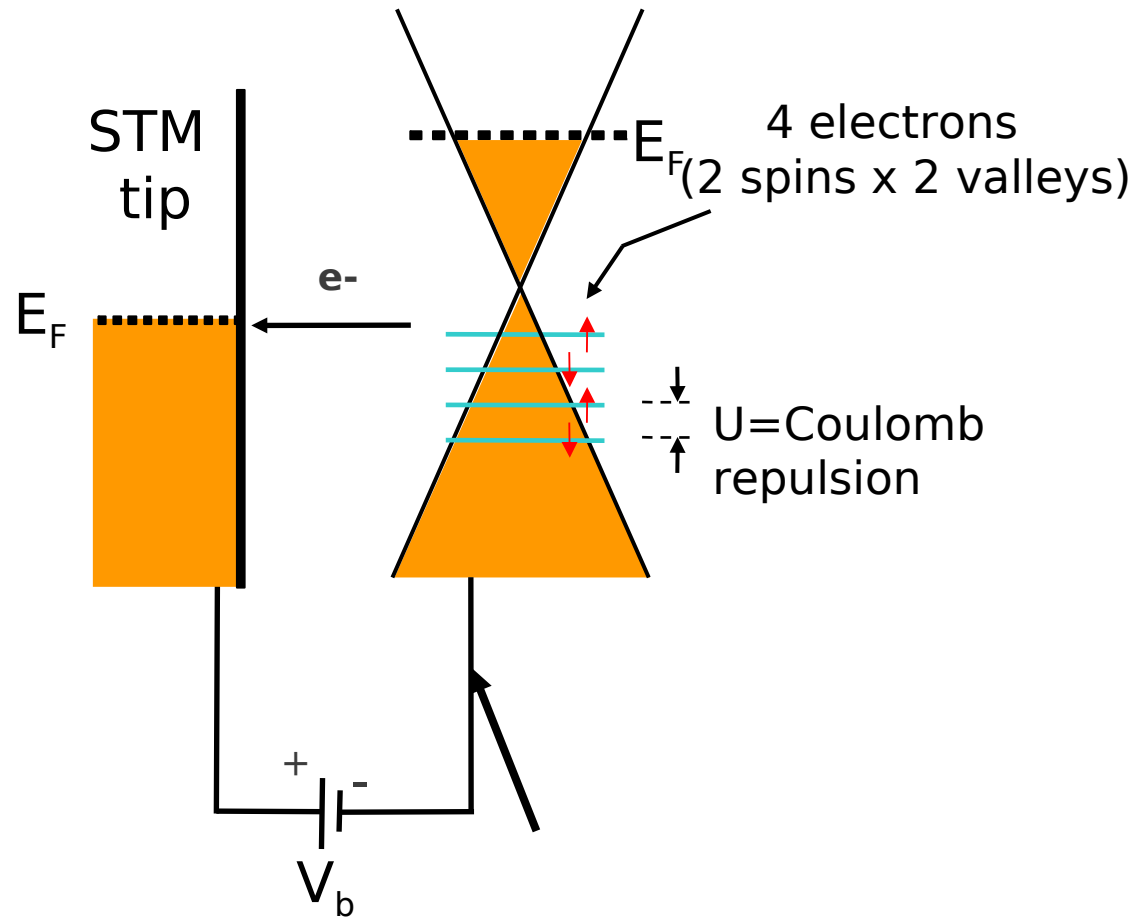
Why the Strong Suppression for Occupied Regime?

Possible explanation: electron-electron interactions

p-doped case



n-doped case (electron occupied)



Suppression of Single-Particle Spectral Density

Atomic collapse in graphene

1. Graphene: relativistic high-energy physics in a condensed matter system. Atomic collapse near charge impurities, $Z \sim 1$.
2. Manifestations: formation of resonances and Dirac vacuum polarization. STM experiments with $Z \sim 1$ impurities. Collapse observed on artificial nuclei (few-impurity clusters)
3. Use Aharonov-Bohm solenoid (B field or pseudo-B field) to bring collapse threshold from $Z \sim 1$ down to $Z \ll 1$. Exact solution for screening cloud near critical point.

Future:

- (1) New electron-electron interaction effects
- (2) Impurity-impurity interactions
- (3) Supercriticality in other systems? (Topol. insulators; $\text{mass} \neq 0$)
- (4) Spin Effects