

Probing 3N-SRCs Using Inclusive and Exclusive Reactions

Misak Sargsian
Florida International University, Miami



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#motivations:

- High Density Fluctuations in Nuclei

Can such fluctuations be probed? – 2N, 3N SRCs

What they tell us about the nuclear forces at short distances

QCD/ Hadron transitions

Partons in the nuclear medium/EMC effects

Nature of the nuclear core

Nature of 3N forces

- Astrophysical Implications

Conceptually: How to probe nuclei at short nucleon separations

- Probe bound nucleons at large internal momenta
- Need high energy probes to resolve such nucleons in nuclei

#Theory of High Energy eA Scattering:

Emergence of High Energy Dynamics

- Short Range Nucleon Correlations(SRCs) are important feature of Nuclear Dynamics
- Transition from hadronic to quark-gluon degrees of freedom in cold-nuclei “should happen” through SRCs
- Internal momenta relevant to SRCs $p \sim M_N$ - Relativistic
- Such states are probed in high energy processes: high energy approximations can be applied in description of SRC dynamics

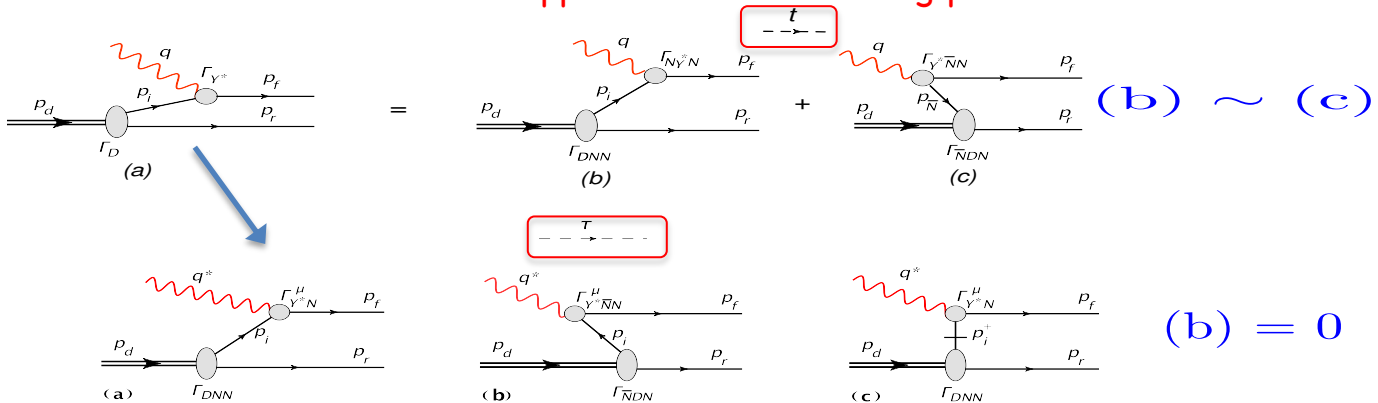
Light-Front wave function of the Nucleus

$$\Psi_{LF}(\mathcal{Z}_1, \mathcal{Z}_2, \mathcal{Z}_3, \dots, \mathcal{Z}_n, \tau)$$

- in the momentum space

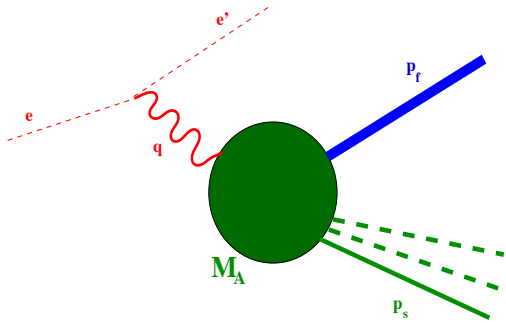
$$\Psi_{LF}(\alpha_1, p_{1\perp}; \alpha_2, p_{2\perp}; \alpha_3, p_{3\perp}; \dots, \alpha_n, p_{n\perp}) \quad \alpha_i = \frac{p_{i-}}{p_{A-}/A}$$

- How the LF wave function appears in the scattering process



Frank Vera, M.S. ArXiv 2018

High Energy Approximations:



$$|\vec{q}| = q_3 \sim p_{f3} \gg p \sim M_N$$

$$Q^2 \geq \text{few GeV}^2$$

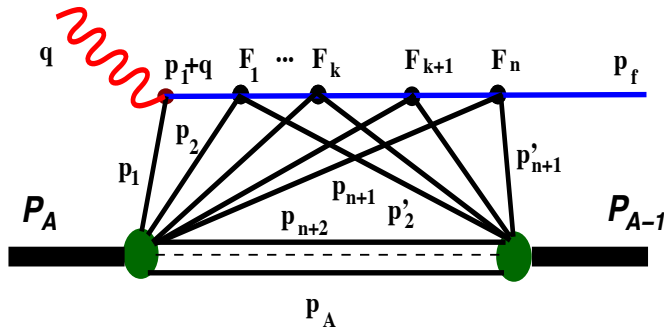
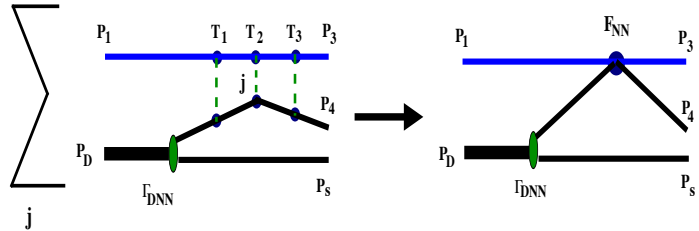
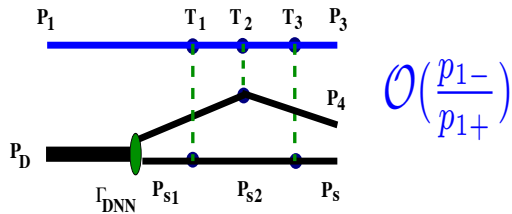
Both for QE/DIS

- Emergence of the small parameter

$$\frac{q_-}{q_+} = \frac{q_0 - q_3}{q_0 + q_3} \ll 1 \quad \mathcal{O}\left(\frac{q_-}{q_+}\right)$$

$$\frac{p_{f-}}{p_{f+}} = \frac{E_f - p_{f3}}{E_f + p_{f3}} \ll 1 \quad \mathcal{O}\left(\frac{p_{f-}}{p_{f+}}\right)$$

Emergence of "effective" theory

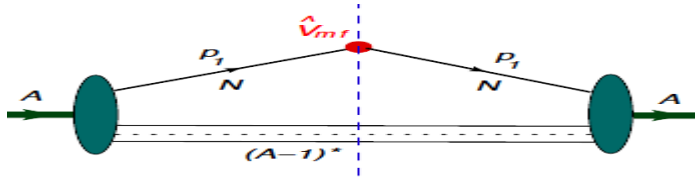


Effective Feynman
Diagrammatic Rules

M.S. IJMS 2001

Wave function?

Spectral Function Calculations



$$S_A^{MF} = -Im \int \chi_A^\dagger \Gamma_{A,N,A-1}^\dagger \frac{\not{p}_1 + m}{p_1^2 - m^2 + i\varepsilon} \hat{V}^{MF} \frac{\not{p}_1 + m}{p_1^2 - m^2 + i \times \varepsilon} \left[\frac{G_{A-1}(p_{A-1})}{p_{A-1}^2 - M_{A-1}^2 + i\varepsilon} \right]^{on} \Gamma_{A,N,A-1} \chi_A \frac{d^4 p_{A-1}}{i(2\pi)^4}$$

$$\hat{V}^{MF} = i a^\dagger(p_1, s_1) \delta^3(p_1 + p_{A-1}) \delta(E_m - E_\alpha) a(p_1, s_1)$$

$$\psi_{N/A}(p_1, s_1, s_A, E_\alpha) = \frac{\bar{u}(p_1, s_1) \Psi_{A-1}^\dagger(p_{A-1}, s_{A-1}, E_\alpha) \Gamma_{A,N,A-1} \chi_A}{(M_{A-1}^2 - p_{A-1}^2) \sqrt{(2\pi)^3 2E_{A-1}}}$$

$$S_A^{MF}(p_1, E_m) = \sum_{\alpha} \sum_{s_1, s_{A-1}} | \psi_{N/A}(p_1, s_1, s_A, E_\alpha) |^2 \delta(E_m - E_\alpha)$$

1.Short-Range Correlations

for large $k > k_{Fermi}$

$$n_A^{LC}(\alpha, p_t) \approx a_{NN}(A, z) n_{NN}^{LC}(\alpha, p_t)$$

Frankfurt, Strikman Phys.
Rep, 1988
Day, Frankfurt, Strikman,
MS, Phys. Rev. C 1993

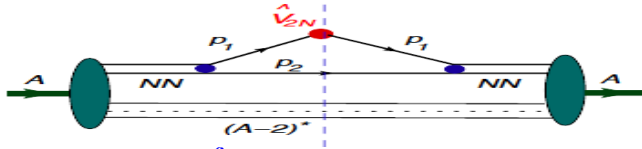
- Experimental observations

Egiyan et al, 2002, 2006
Fomin et al, 2011

$$n_A^{LC}(\alpha, p_t) \approx a_{pn}(A, z) n_d^{LC}(\alpha, p_t)$$

E. Piasetzky, MS, L. Frankfurt,
M. Strikman, J. Watson PRL, 2006

2N SRC model



$$\begin{aligned}
 P_{A,2N}^N(\alpha_1, p_{1,\perp}, \tilde{M}_N^2) &= \sum_{s_2, s_{NN}, s_{A-2}} \int \chi_A^\dagger \Gamma_{A \rightarrow NN, A-2} \chi_{A-2}(p_{A-2}, s_{A-2}) \\
 &\times \frac{\chi_{NN}(p_{NN}, s_{NN}) \chi_{NN}^\dagger(p_{NN}, s_{NN})}{p_{NN}^2 - \tilde{M}_{NN}^2} \Gamma_{NN \rightarrow NN}^\dagger \frac{u(p_1, s_1) u(p_2, s_2)}{p_1^2 - M_N^2} \\
 &\times \left[2\alpha_1^2 \delta(\alpha_1 + \alpha_2 + \alpha_{A-2} - A) \delta^2(p_{1,\perp} + p_{2,\perp} + p_{A-2,\perp}) \delta(\tilde{M}_N^2 - \tilde{M}_N^{(2N),2}) \right] \frac{\bar{u}(p_1, s_1) \bar{u}(p_2, s_2)}{p_1^2 - M_N^2} \\
 &\times \Gamma_{NN \rightarrow NN} \frac{\chi_{NN}(p_{NN}, s_{NN}) \chi_{NN}^\dagger(p_{NN}, s_{NN})}{p_{NN}^2 - \tilde{M}_{NN}^2} \chi_{A-2}^\dagger(p_{A-2}, s_{A-2}) \Gamma_{A, NN, A-2} \chi_A \\
 &\times \frac{d\alpha_2}{\alpha_2} \frac{d^2 p_{2,\perp}}{2(2\pi)^3} \frac{d\alpha_{A-2}}{\alpha_{A-2}} \frac{d^2 p_{A-2,\perp}}{2(2\pi)^3}.
 \end{aligned} \tag{1}$$

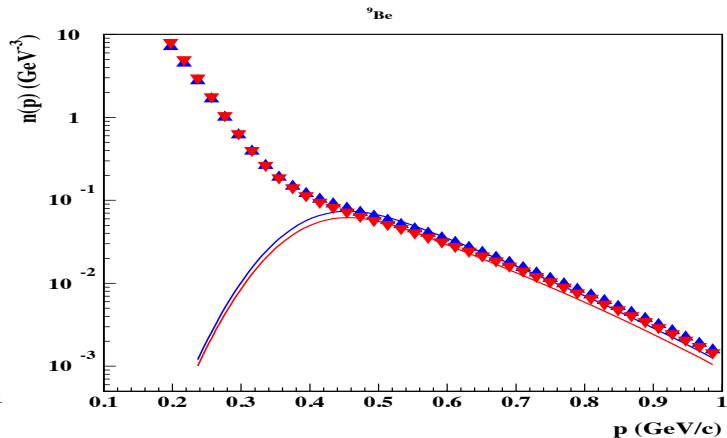
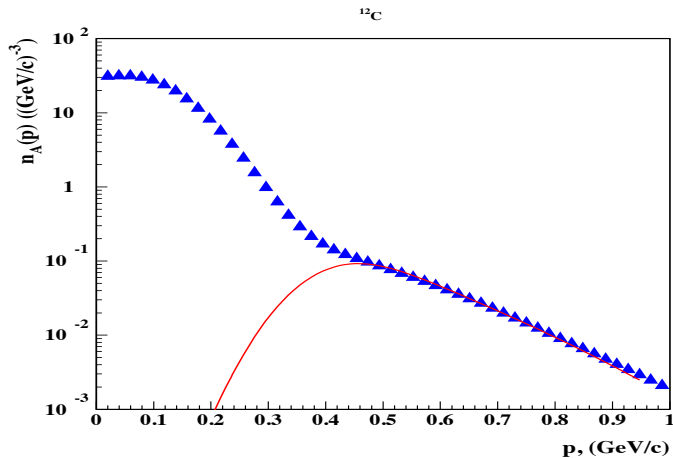
O. Artiles & M.S. Phys. Rev. C
2016

$\psi_{pn}^{LF}(\alpha, p_\perp) \approx C \psi_d^{LF}(\alpha, p_\perp)$

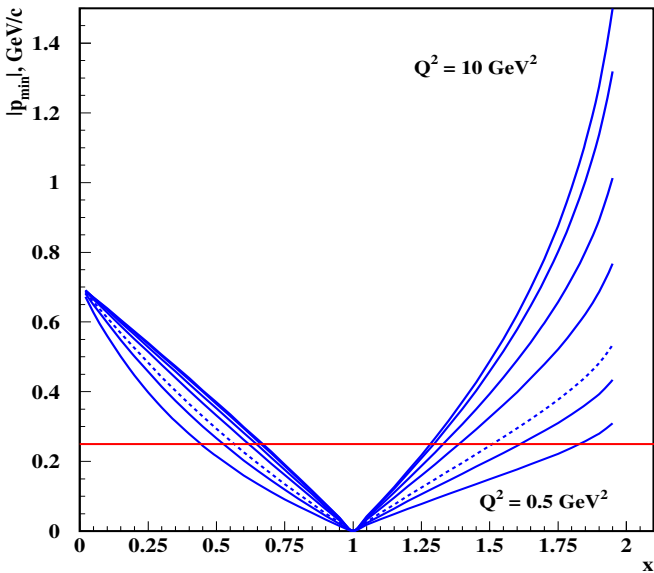
$$\psi_{2N}^{s_{NN}}(\beta_1, k_{1,\perp}, s_1, s_2) = -\frac{1}{\sqrt{2(2\pi)^3}} \frac{\bar{u}(p_1, s_1) \bar{u}(p_2, s_2) \Gamma_{NN \rightarrow NN} \cdot \chi_{NN}(p_{NN}, s_{NN})}{\frac{1}{2} [M_{NN}^2 - 4(M_N^2 + k_1^2)]}$$

$$\psi_{CM}(\alpha_{NN}, k_{NN,\perp}) = -\frac{1}{\sqrt{\frac{A-2}{2}}} \frac{1}{\sqrt{2(2\pi)^3}} \frac{\chi_{NN}^\dagger(p_{NN}, s_{NN}) \chi_{A-2}^\dagger(p_{A-2}, s_{A-2}) \Gamma_{A \rightarrow NN, A-2} \chi_A^{s_A}}{\frac{2}{A} [M_A^2 - s_{NN, A-2}(k_{CM})]}$$

2N SRC model Non Relativistic Approximation

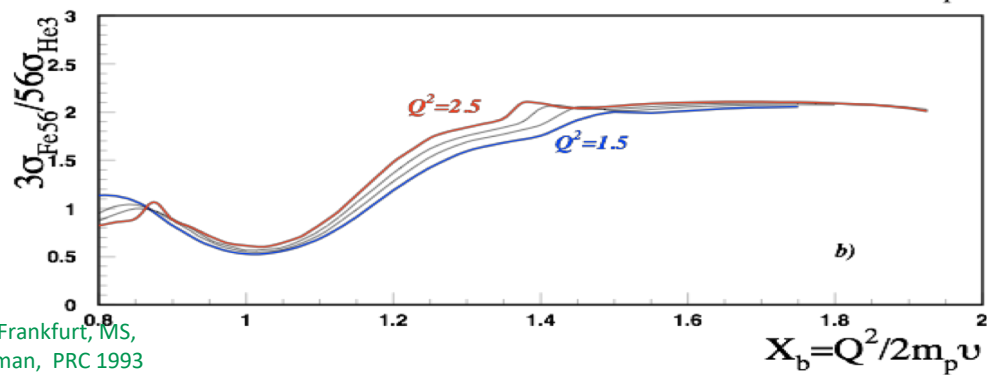
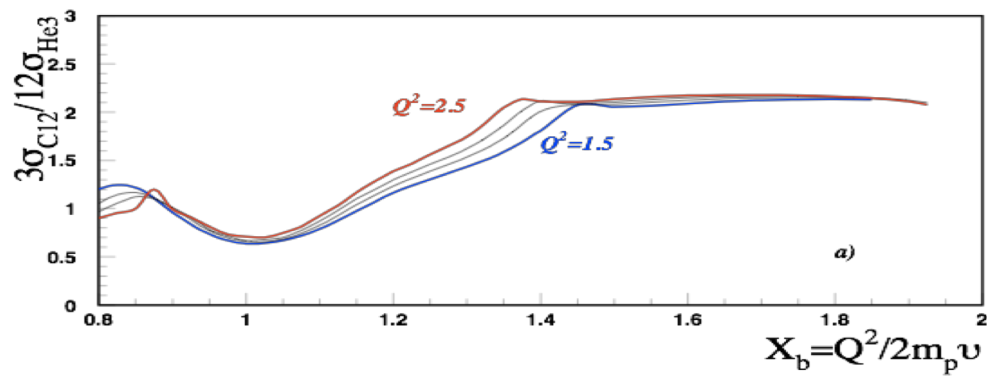


Nuclear Scaling in QE Inclusive $A(e,e')X$ reaction at $x > 1$ region



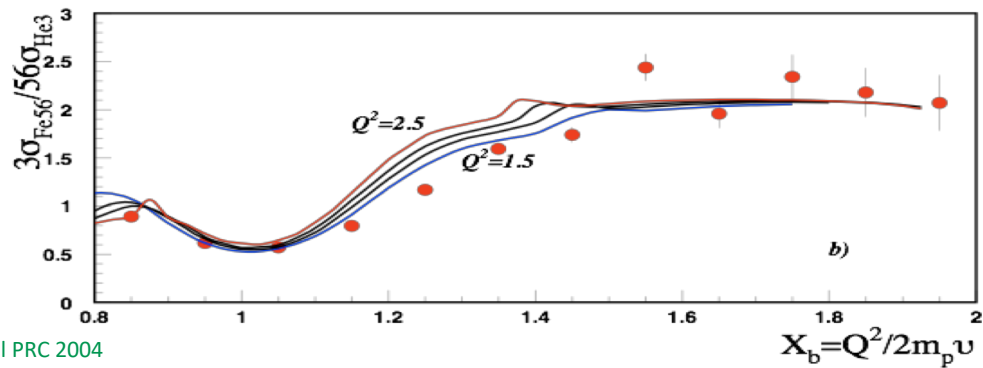
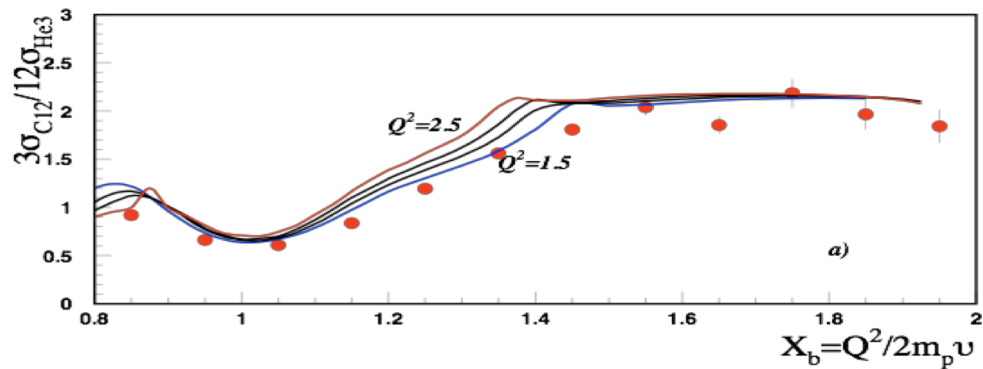
Day, Frankfurt, MS,
Strikman, PRC 1993

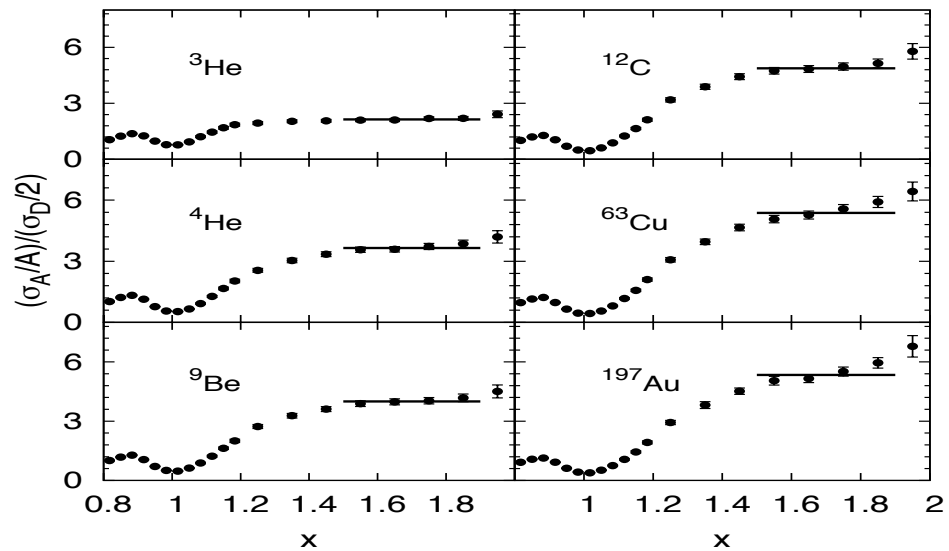
$$A(e,e')$$



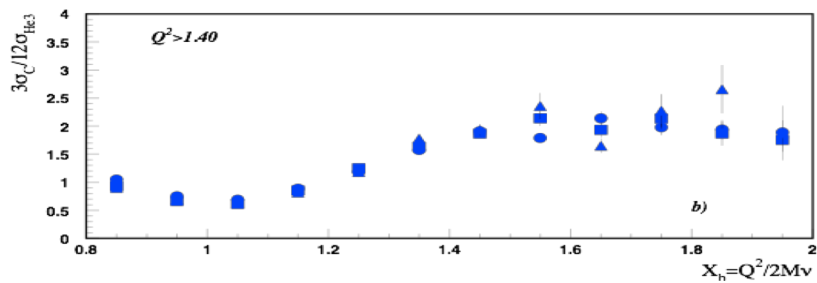
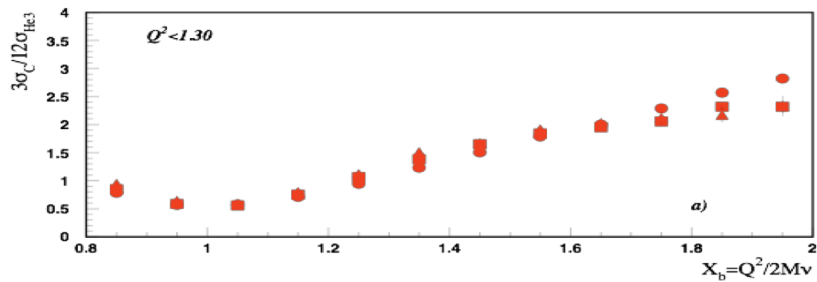
Day, Frankfurt, MS,
Strikman, PRC 1993

$$A(e,e')$$





$A(e,e')$



Meaning of the scaling values

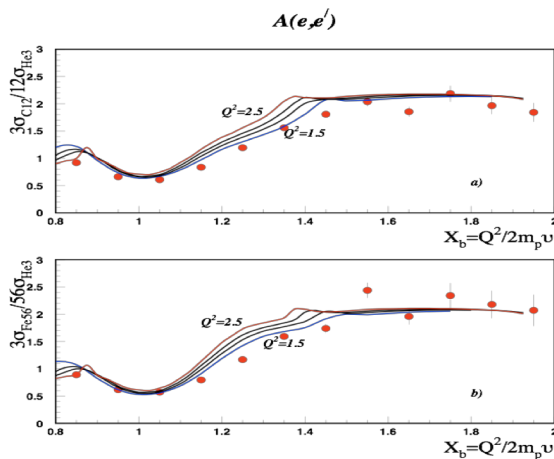
Day, Frankfurt, MS,
Strikman, PRC 1993

Frankfurt, MS, Strikman,
IJMP A 2008

Fomin et al PRL 2011

$$R = \frac{A_2 \sigma[A_1(e, e')X]}{A_1 \sigma[A_2(e, e')X]}$$

$$\text{For } 1 < x < 2 \quad R \approx \frac{a_2(A_1)}{a_2(A_2)}$$



Egiyan, et al PRL 2006, PRC 2004

a2's as relative probability of 2N SRCs

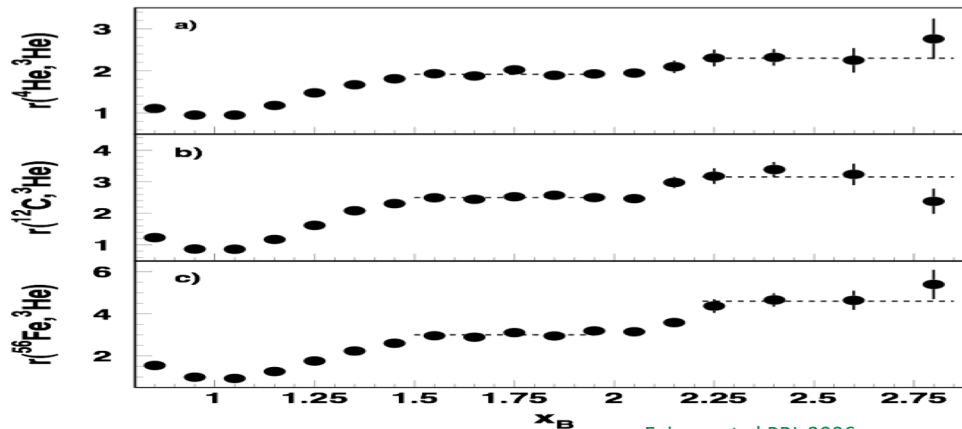
Table 1: The results for $a_2(A, y)$

A	y	This Work	Frankfurt et al	Egiyan et al	Famin et al
^3He	0.33	2.07 ± 0.08	1.7 ± 0.3		2.13 ± 0.04
^4He	0	3.51 ± 0.03	3.3 ± 0.5	3.38 ± 0.2	3.60 ± 0.10
^9Be	0.11	3.92 ± 0.03			3.91 ± 0.12
^{12}C	0	4.19 ± 0.02	5.0 ± 0.5	4.32 ± 0.4	4.75 ± 0.16
^{27}Al	0.037	4.50 ± 0.12	5.3 ± 0.6		
^{56}Fe	0.071	4.95 ± 0.07	5.6 ± 0.9	4.99 ± 0.5	
^{64}Cu	0.094	5.02 ± 0.04			5.21 ± 0.20
^{197}Au	0.198	4.56 ± 0.03	4.8 ± 0.7		5.16 ± 0.22

Towards Three Nucleon Short Range Correlations

Looking for the Plateau in Inclusive Cross Section Ratios $x > 2$

For $1 < x < 2$ $R \approx \frac{a_2(A_1)}{a_2(A_2)}$ For $2 < x < 3$ $R \approx \frac{a_3(A_1)}{a_3(A_2)}$

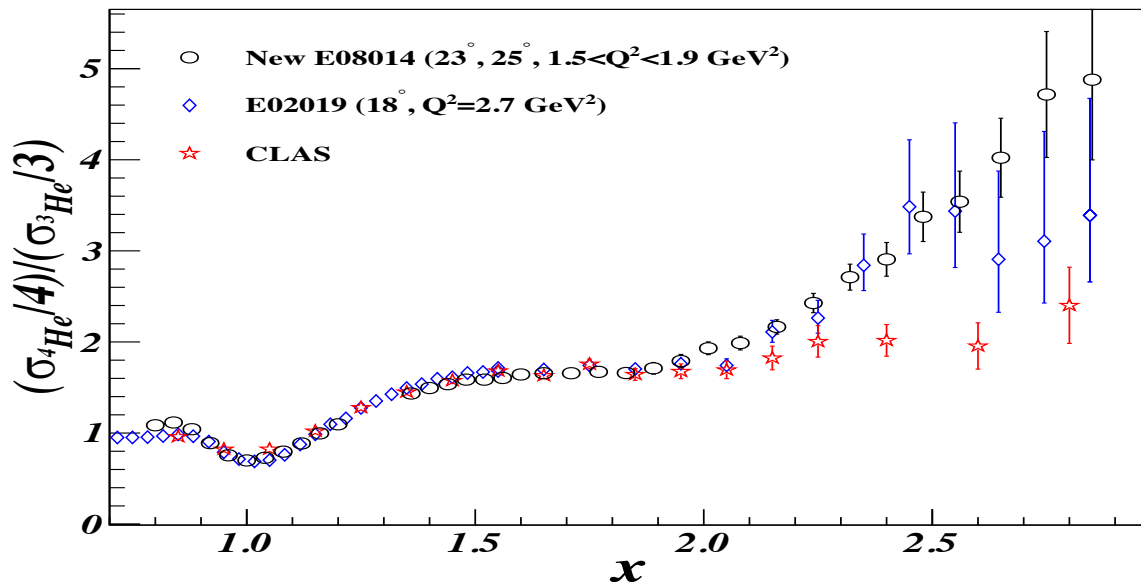


Egiyan, et al PRL 2006

3N SRCs

Z. Ye, et al, 2017

Looking for the Plateau in Inclusive Cross Section Ratios $x > 2$



Egiyan, et al PRL 2006

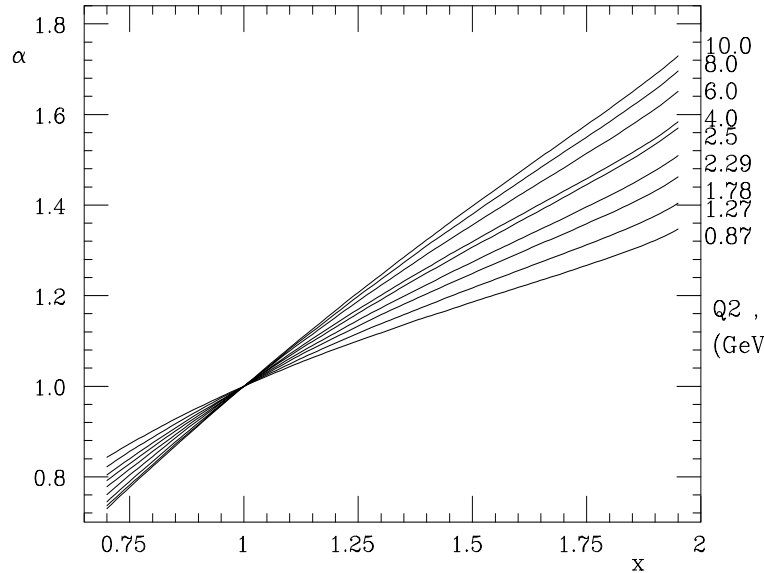
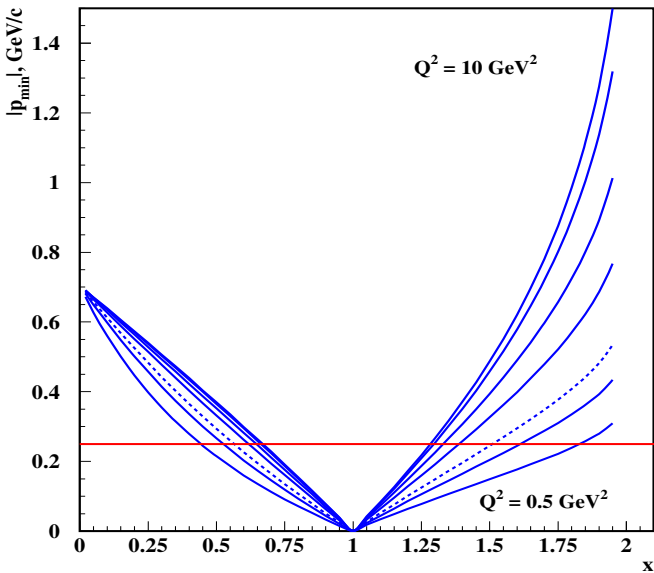
Egiyan, et al PRC 2004

2N SRCs:

Proper Variables of 2N SRC are

- the Light Front Momentum Fraction: $\alpha = \frac{p_N^+}{p_{NN}^+}$
- transverse momentum: $p_\perp$

Nuclear Scaling in QE Inclusive $A(e,e')X$ reaction at $x > 1$ region



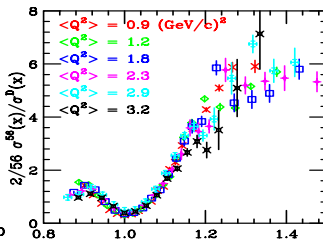
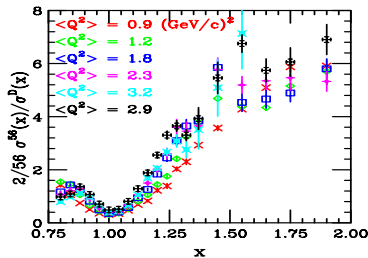
Day, Frankfurt, MS,
Strikman, PRC 1993

Back to inclusive $A(e,e')X$ scattering

$$\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha) \quad \text{where} \quad \rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$$

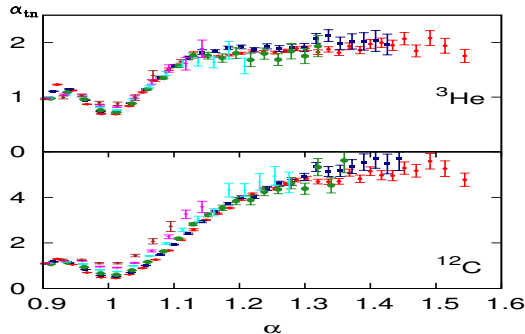
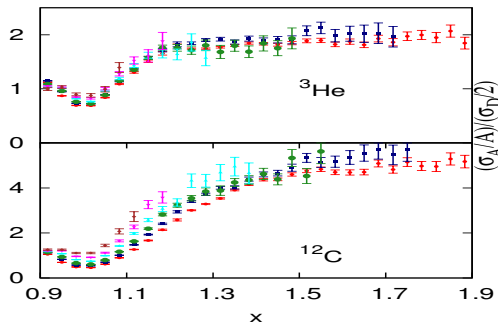
$$1.3 \leq \alpha_{2N} \leq 1.5$$

$$\alpha_{2N} = 2 - \frac{q_- + 2m_N}{2m_N} \left(1 + \frac{\sqrt{W_{2N}^2 - 4m_N^2}}{W_{2N}} \right)$$



$$\alpha \left| \begin{array}{l} Q^2 \rightarrow \infty \rightarrow x \\ x \rightarrow 1 \rightarrow 1 \end{array} \right.$$

J. Arrington, D. Higinbotham
G. Rosner, M.S. Prog. PNP 2012

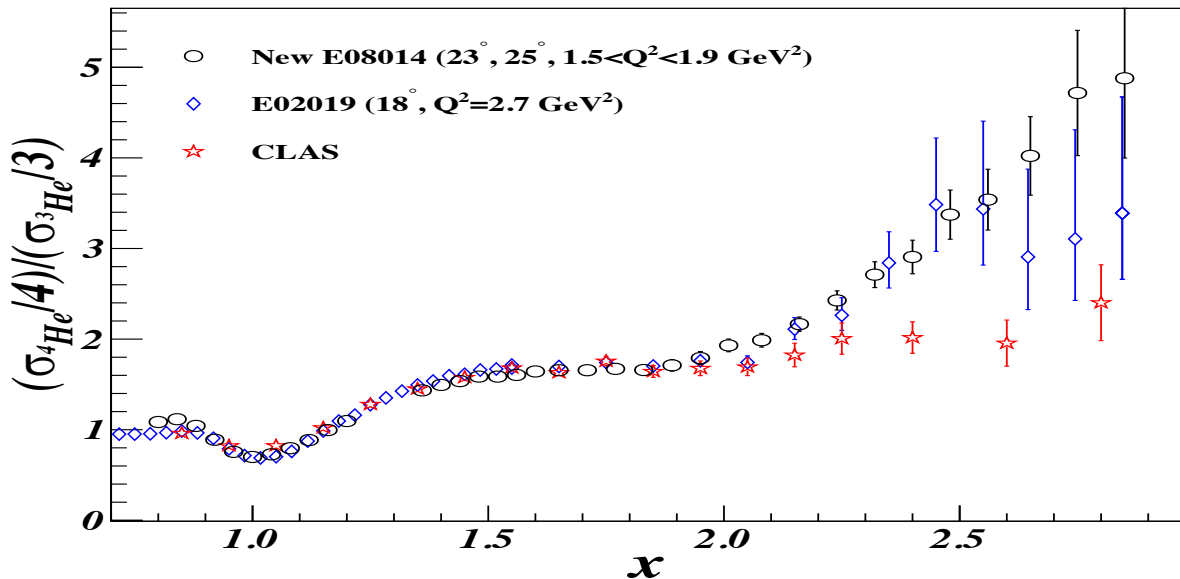


N. Fomin, D. Higinbotham
M.S., P. Sovignon ARNPS, 2017

Three Nucleon Short Range Correlations

Z. Ye, et al, 2017

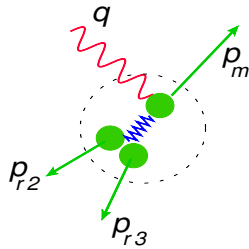
Looking for the Plateau in Inclusive Cross Section Ratios $x > 2$



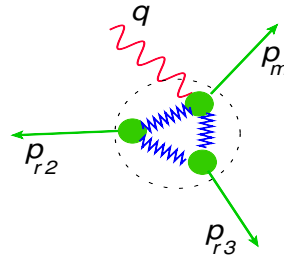
3N SRCs:

Proper Variables of 3N SRC are

- the Light Front Momentum Fraction: $\alpha = \frac{p_N^+}{p_{3N}^+}$
- transverse momentum: $p_\perp$

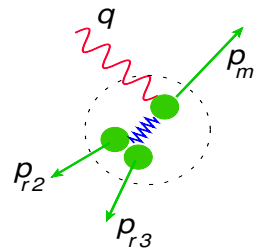
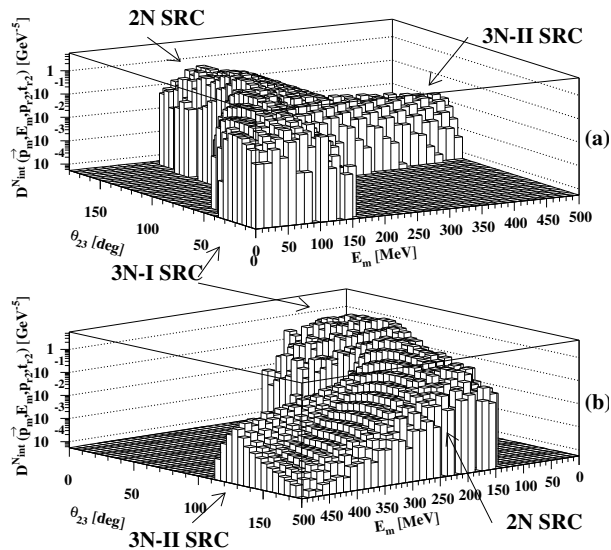


(a)

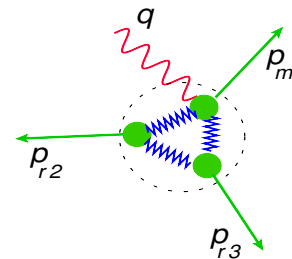


(b)

3N SRCs:



(a)



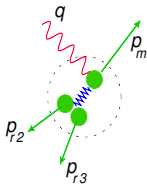
(b)

3N SRC model $\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha)$ where $\rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$

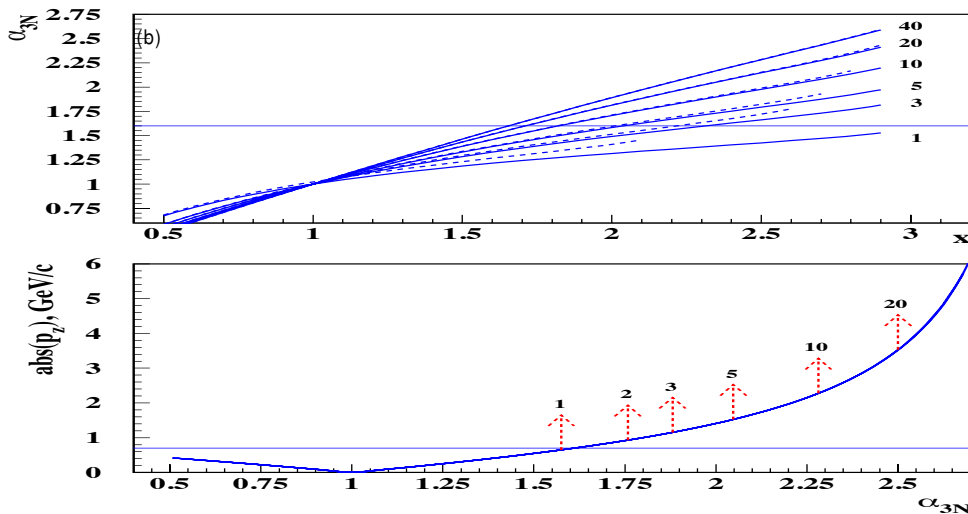
$$1.6 \leq \alpha_{3N} < 3$$

$$q + 3m_N = p_f + p_s$$

$$= 3 - \frac{q_- + 3m_N}{2m_N} \left[1 + \frac{m_S^2 - m_N^2}{W_{3N}^2} + \sqrt{\left(1 - \frac{(m_S + m_n)^2}{W_{3N}^2}\right) \left(1 - \frac{(m_S - m_n)^2}{W_{3N}^2}\right)} \right] (1)$$



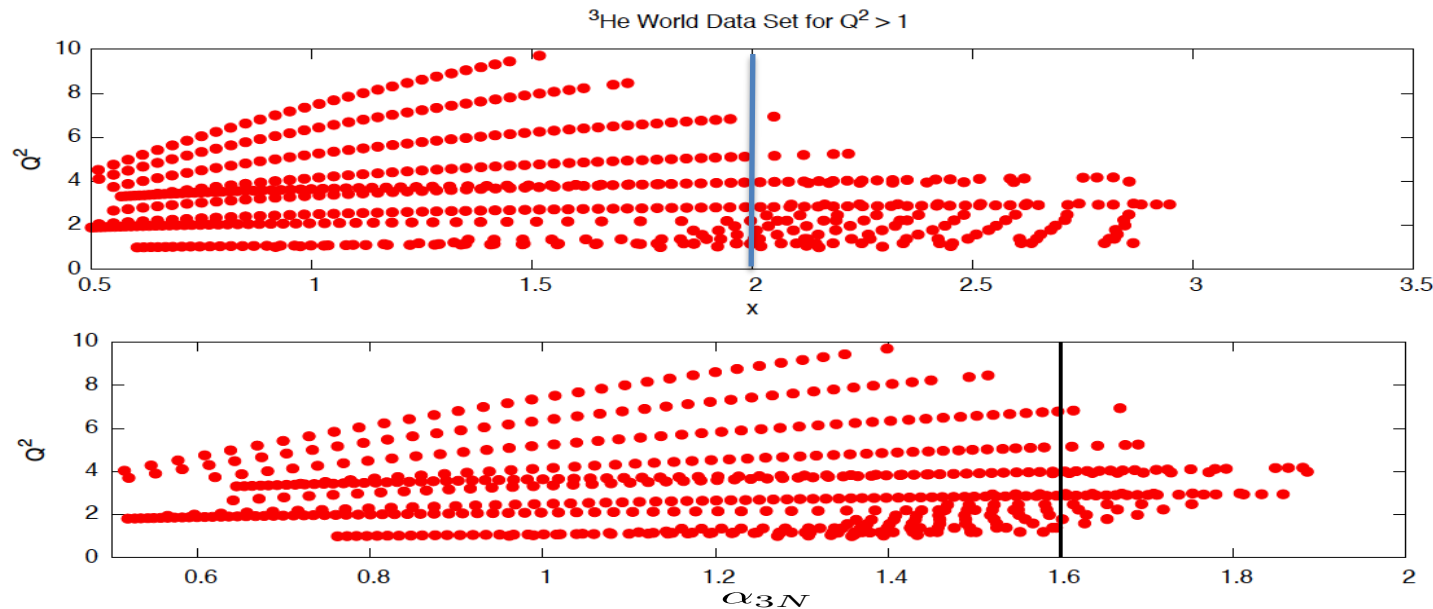
(a)



3N SRC model $\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha)$ where $\rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$

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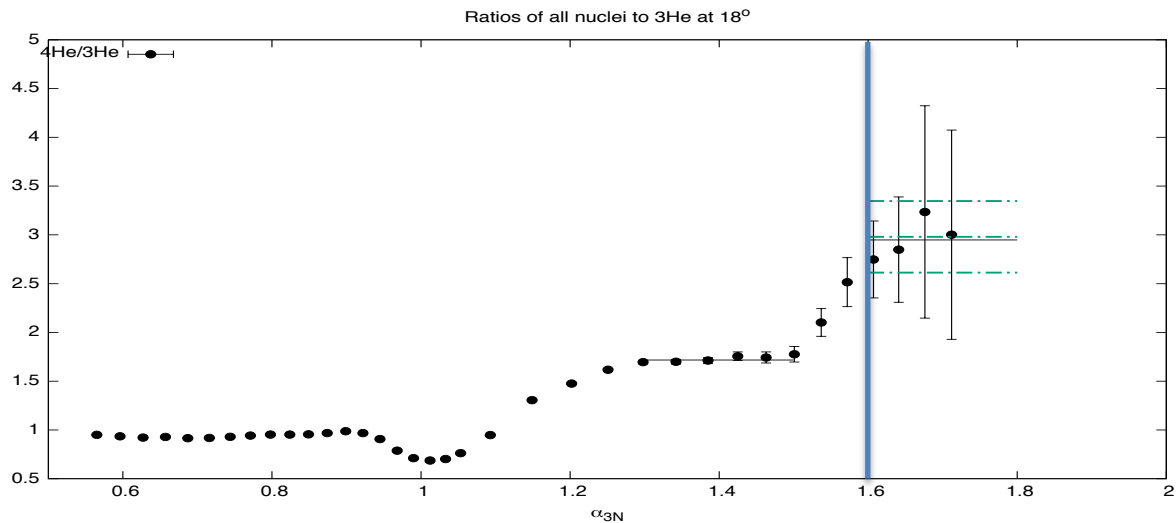
Donal Day, M.S., Frankfurt, Strikman 2018



3N SRCs

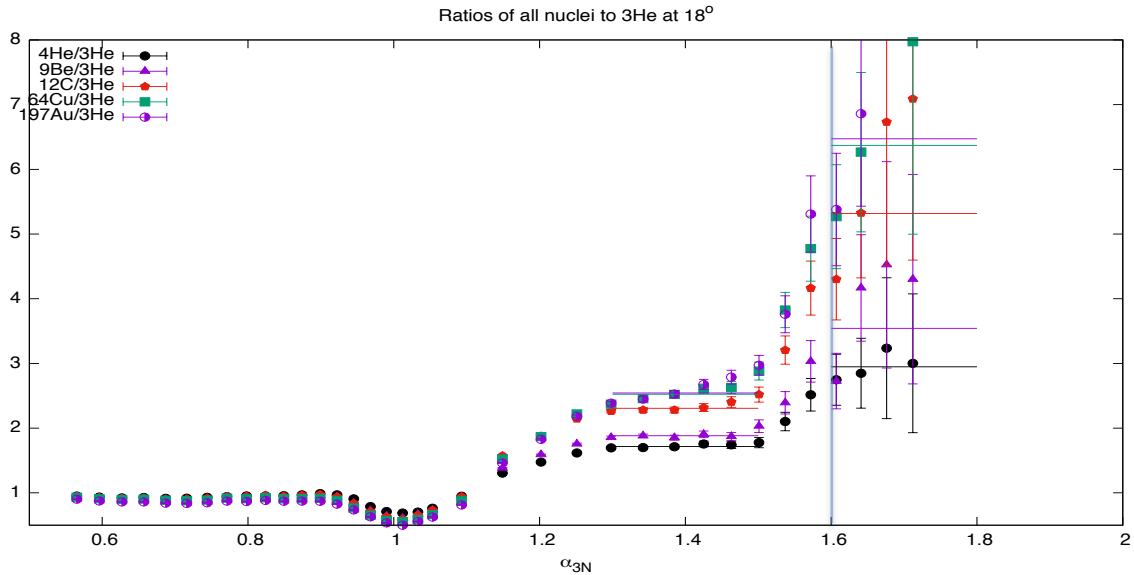
$$\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha) \quad \text{where} \quad \rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$$

$$1.6 \leq \alpha_{3N} < 3$$



JLab - E02019 - Data

3N SRC model $\sigma_{eA} = \sum \sigma_{eN} \cdot \rho_A(\alpha)$ where $\rho_A(\alpha) = \int \rho_A(\alpha, p_\perp) d^2 p_\perp$
 $1.6 \leq \alpha_{3N} < 3$

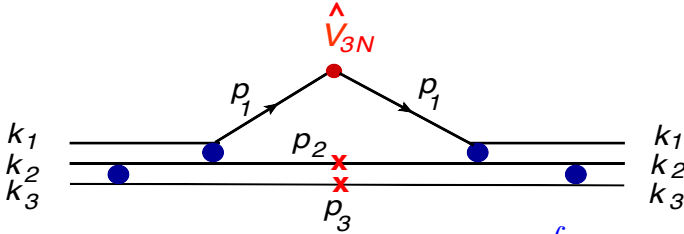


JLab - E02019 - Data

3N SRC: Light-Cone Momentum Fraction Distribution

A. Freese, M.S., M. Strikman, Eur. Phys. J 2015

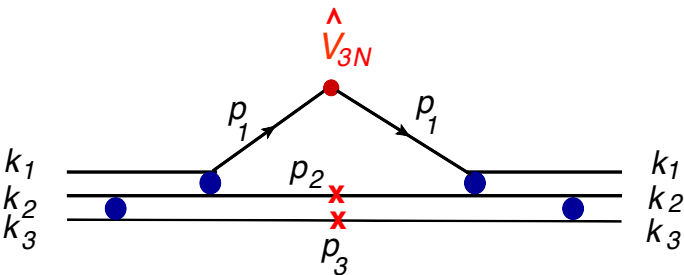
O. Artiles M.S. Phys. Rev. C 2016



$$\begin{aligned}
 P_{A,3N}^N(\alpha_1, p_{1,\perp}, s_1, \tilde{M}_N) = & \sum_{s_2, s_3, s_{2'}, \tilde{s}_{2'}} \int \bar{u}(k_1) \bar{u}(k_2) \bar{u}(k_3) \Gamma_{NN \rightarrow NN}^\dagger \frac{u(p_{2'}, \tilde{s}_{2'}) \bar{u}(p_{2'}, \tilde{s}_{2'})}{p_{2'}^2 - M_N^2} \\
 & \times \Gamma_{NN \rightarrow NN}^\dagger \frac{u(p_1, s_1)}{p_1^2 - M_N^2} u(p_2, s_2) \left[2\alpha_1^2 \delta(\alpha_1 + \alpha_2 + \alpha_3 - 3) \delta^2(p_{1\perp} + p_{2\perp} + p_{3\perp}) \delta(\tilde{M}_N^2 - M_N^{3N,2}) \right] \\
 & \times \bar{u}(p_2, s_2) \frac{\bar{u}(p_1, s_1)}{p_1^2 - M_N^2} \Gamma_{NN \rightarrow NN} \frac{u(p_{2'}, s_{2'}) \bar{u}(p_{2'}, s_{2'})}{p_{2'}^2 - M_N^2} u(p_3, s_3) \bar{u}(p_3, s_3) \Gamma_{NN \rightarrow NN}^\dagger u(k_1) u(k_2) u(k_3) \\
 & \times \frac{d\alpha_2}{\alpha_2} \frac{d^2 p_{2\perp}}{2(2\pi)^3} \frac{d\alpha_3}{\alpha_3} \frac{d^2 p_{3\perp}}{2(2\pi)^3}, \tag{1}
 \end{aligned}$$

$$\begin{aligned}
 P_{A,3N}^N(\alpha_1, p_{1,\perp}, \tilde{M}_N) = & \int \frac{3 - \alpha_3}{2(2 - \alpha_3)^2} \boxed{\rho_{NN}(\beta_3, p_{3\perp}) \rho_{NN}(\beta_1, \tilde{k}_{1\perp})} 2\delta(\alpha_1 + \alpha_2 + \alpha_3 - 3) \\
 & \delta^2(p_{1\perp} + p_{2\perp} + p_{3\perp}) \delta(\tilde{M}_N^2 - M_N^{3N,2}) d\alpha_2 d^2 p_{2\perp} d\alpha_3 d^2 p_{3\perp}, \tag{1}
 \end{aligned}$$

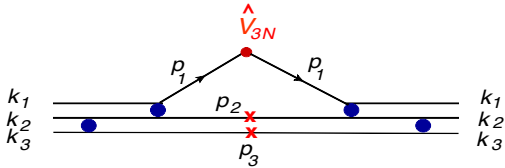
3N SRC: Light-Cone Momentum Fraction Distribution



$$\begin{aligned}
 \rho_{3N}(\alpha_1) = & \int \frac{1}{4} \left[\frac{3 - \alpha_3}{(2 - \alpha_3)^3} \boxed{\rho_{pn}(\alpha_3, p_{3\perp}) \rho_{pn}\left(\frac{2\alpha_2}{3 - \alpha_3}, p_{2\perp} + \frac{\alpha_1}{3 - \alpha_3} p_{3\perp}\right)} + \right. \\
 & \left. \frac{3 - \alpha_2}{(2 - \alpha_2)^3} \boxed{\rho_{pn}(\alpha_2, p_{2\perp}) \rho_{pn}\left(\frac{2\alpha_3}{3 - \alpha_2}, p_{3\perp} + \frac{\alpha_1}{3 - \alpha_2} p_{2\perp}\right)} \right] \delta\left(\sum_{i=1}^3 \alpha_i - 3\right) \\
 & d\alpha_2 d^2 p_{2\perp} d\alpha_3 d^2 p_{3\perp},
 \end{aligned} \tag{1}$$

$$\boxed{\rho_{pn}(\alpha, p_{\perp}) \approx a_2(A) \rho_d(\alpha, p_{\perp})}$$

3N SRC: Light-Cone Momentum Fraction Distribution



$$-\rho_{3N} \sim a_2(A, z)^2$$

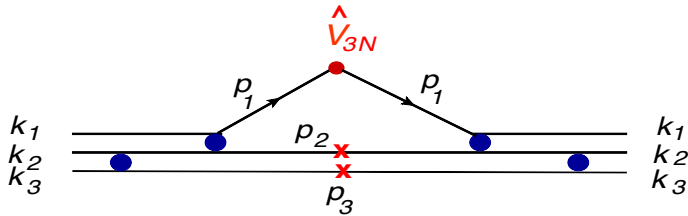
- For $A(e, e')$ X reactions: $\sigma_{eA} = \sum_N \sigma_{eN} \rho_{3N}(\alpha_{3N})$

- Defining: $R_3(A, Z) = \frac{3\sigma_{eA}}{A\sigma_{e^3He}} \big|_{\alpha_{3N} \geq \alpha_{3N}^0}$

- We predict: $R_3(A, Z) = \frac{9}{8} \frac{(\sigma_{ep} + \sigma_{en})/2}{(2\sigma_{ep} + \sigma_{en})/3} \left(\frac{a_2(A, Z)}{a_2(^3He)} \right)^2 = \frac{9}{8} \frac{(\sigma_{ep} + \sigma_{en})/2}{(2\sigma_{ep} + \sigma_{en})/3} R_2^2(A, Z),$

- Where: $R_2(A, Z) = \frac{3\sigma_{eA}}{A\sigma_{e^3He}} \big|_{1.3 \leq \alpha_{3N} \leq 1.5}$ where: $\alpha_{3N} \approx \alpha_{2N}$

3N SRC: Light-Cone Momentum Fraction Distribution



$$-\rho_{3N} \sim a_2(A, z)^2$$

- ppp and nnn strongly suppressed compared with ppn or pnn
- pp/nn recoil state is suppressed compared with pn

$$R_3(A, Z) = \frac{9}{8} \frac{(\sigma_{ep} + \sigma_{en})/2}{(2\sigma_{ep} + \sigma_{en})/3} \left(\frac{a_2(A, Z)}{a_2(^3He)} \right)^2 =$$

$$\frac{9}{8} \frac{(\sigma_{ep} + \sigma_{en})/2}{(2\sigma_{ep} + \sigma_{en})/3} R_2^2(A, Z),$$

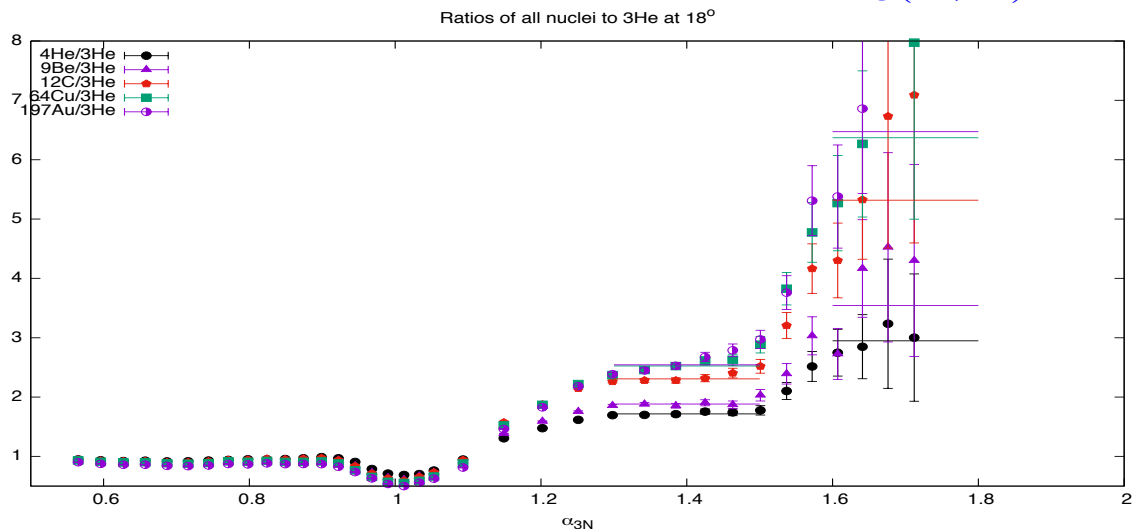
3N SRC model

$$R_2 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.3 \leq \alpha_{3N} \leq 1.5$$

$$1.6 \leq \alpha_{3N} < 3$$

$$R_3 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.6 \leq \alpha_{3N} \leq 1.8$$

$$R_3(A, Z) \approx R_2(A, Z)^2$$

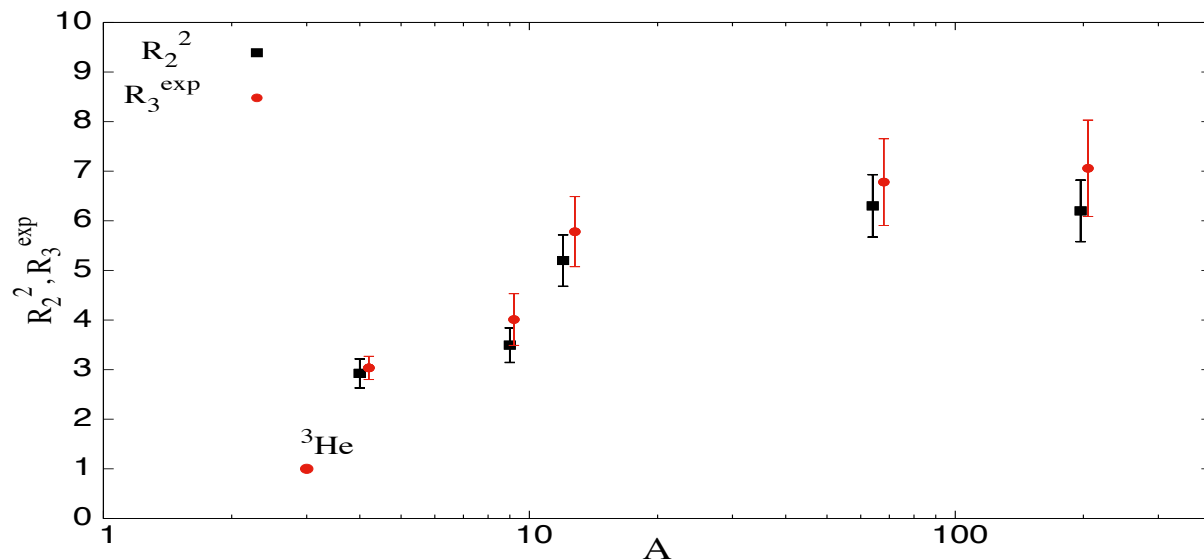


3N SRC model

$$R_2 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.3 \leq \alpha_{3N} \leq 1.5 \quad 1.6 \leq \alpha_{3N} < 3$$

$$R_3 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.6 \leq \alpha_{3N} \leq 1.8 \quad R_3(A) = R_2(A)^2$$

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3N SRC model

Defining: $a_3(A, Z) = \frac{3}{A} \frac{\sigma_{eA}}{(\sigma_{e^3He} + \sigma_{e^3H})/2}$

One relates: $a_3(A, Z) = \frac{(2\sigma_{ep} + \sigma_{en})/3}{(\sigma_{ep} + \sigma_{en})/2} R_3(A, Z)$

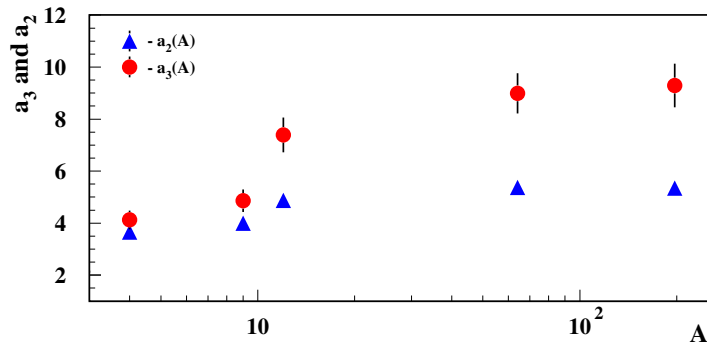
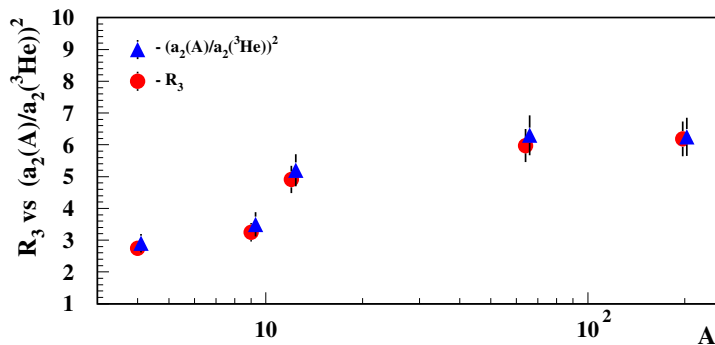
A	a ₂	R ₂	R ₂ ^{exp}	R ₂ ²	R ₃ ^{exp}	a ₃
3	2.14 ±0.04	NA	NA	NA	NA	1
4	3.66 ±0.07	1.71 ±0.026	1.722 ±0.013	2.924 ±0.29	3.034 ±0.23	4.55 ± 0.35
9	4.00 ±0.08	1.84 ±0.027	1.878 ±0.018	3.38 ±0.38	4.01 ±0.52	6.0 ± 0.78
12	4.88 ±0.10	2.28 ±0.027	2.301 ±0.021	5.2 ±0.5	5.78 ±0.71	8.7 ± 1.1
27	5.30 ±0.60	NA	NA	NA	NA	NA
56	4.75 ±0.29	NA	NA	NA	NA	NA
64	5.37 ±0.11	2.51 ±0.027	2.502 ±0.024	6.3 ±0.63	6.780 ±0.875	10.2 ± 1.3
197	5.34 ±0.11	2.46 ±0.028	2.532 ±0.026	6.05 ±0.6	7.059 ±0.970	10.6 ± 1.5

D.Day, L.Frankfurt,M.S, M.Strikman
ArXiv 2018

3N SRC model

$$R_2 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.3 \leq \alpha_{3N} \leq 1.5 \quad 1.6 \leq \alpha_{3N} < 3$$

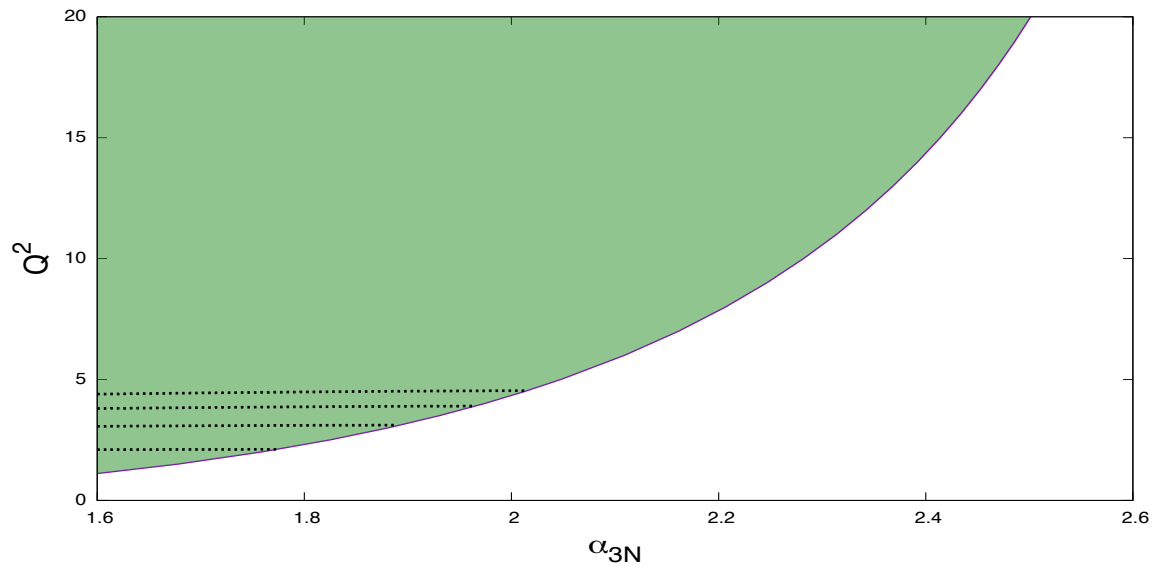
$$R_3 = \frac{3\sigma_{eA}(\alpha_{3N})}{A\sigma_{3A}(\alpha_{3N})} \quad 1.6 \leq \alpha_{3N} \leq 1.8 \quad R_3(A) = R_2(A)^2$$



3N SRC Summary & Outlook

- Proper variable for studies of 2N and 3N SRS are
Light-Cone momentum fractions: α_{2N} , α_{3N}
- It seems we observed first signatures of 3N SRCs in the form of the “scaling”
- Existing data in agreement with the prediction of: $R_3(A, Z) \approx R_2(A, Z)^2$
- Unambiguous verification will require larger Q^2 data to cover larger α_{3N} region
- Reaching $Q^2 > 5 \text{ GeV}^2$ will allow to reach: $\alpha_{3N} > 2$

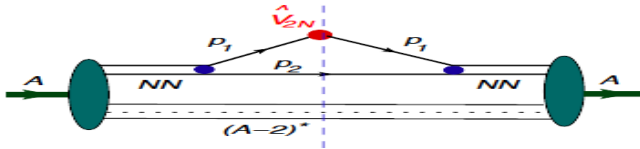
3N SRC Outlook



3N SRCs in Exclusive Processes: $A(e,e'pp)X$

Genuine pp-2NSRC vs pnp 3NSRC

Considering $A(e,e'pp)X$ process: looking for apparent
pp short range correlations

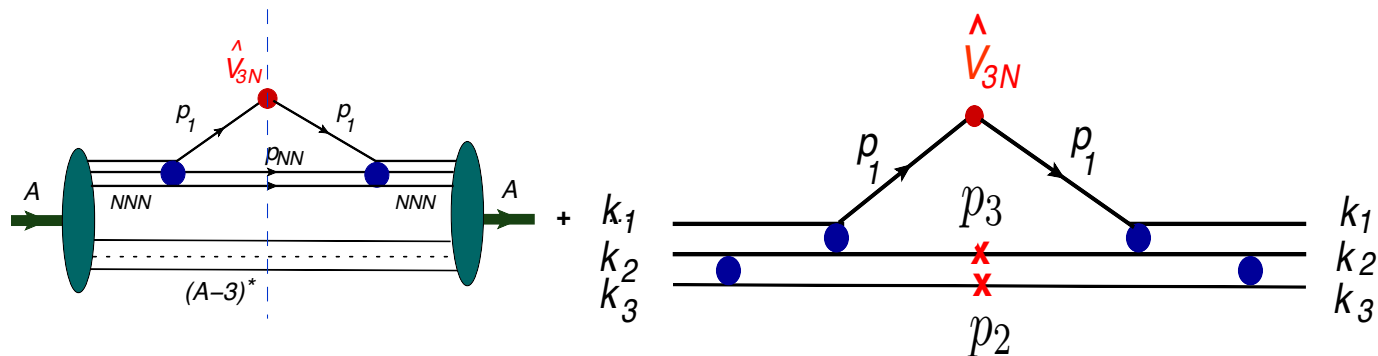


$$D_{A,2N}(\alpha_1, p_{1\perp}, \alpha_2, p_{2\perp}) = \rho_{pp}(\beta_1, k_{1\perp}) \rho_{CM}(\alpha_{cm}, p_{cm,\perp}) \frac{\alpha_2}{\alpha_{cm}}$$

$$\beta_1 = \frac{2\alpha_1}{\alpha_{cm}} \quad \beta_2 = \frac{2\alpha_2}{\alpha_{cm}}$$

$$k_{1\perp} = p_{1\perp} - \frac{\beta_1}{2} p_{cm,\perp}$$

Considering $A(e,e'pp)X$ process: looking for apparent
pp short range correlations



$$D_{A,3N}(\alpha_1, p_{1\perp}, \alpha_2, p_{2\perp}) = \frac{3-\alpha_2}{2(2-\alpha_3)^2} \rho_{pn}(\beta'_3, k'_{3\perp}) \rho_{pn}(\beta_1, k_{1,\perp}) \alpha_3$$

$$\alpha_2 = 3 - \alpha_1 - \alpha_3$$

$$\beta'_3 = 2 - \beta_2$$

$$\beta_1 = \frac{2\alpha_1}{3-\alpha_2}$$

$$\beta_2 \approx \alpha_2$$

$$k_{1\perp} = p_{1\perp} + \frac{\beta_1}{2} p_{2\perp}$$

$$k'_3 \approx -p_{2\perp}$$

