Implications Of EFT For Two-Body Coordinate-Space Nuclear Distributions

Quantitative Challenges in SRC and EMC Research







Joel E. Lynn

December 2, 2016

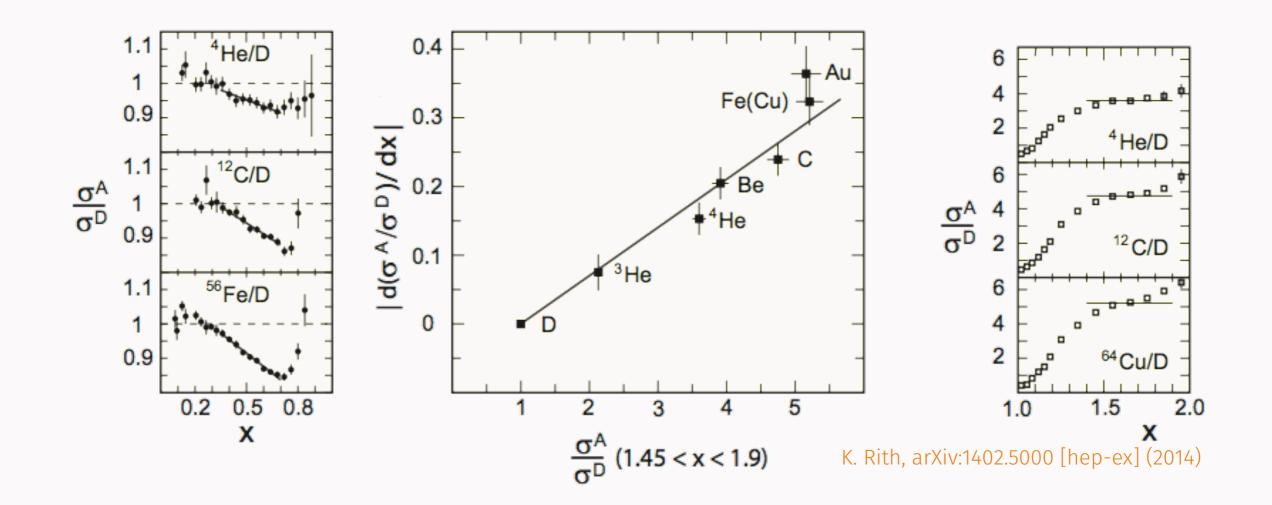
TECHNISCHE

UNIVERSITÄT

DARMSTADT

SRC scaling factor
$$a_2(A, x) \equiv \frac{2\sigma^A}{A\sigma^d}|_{1.5 < x < 2}$$
.

 $dR_{\rm EMC}/dx \propto a_2$



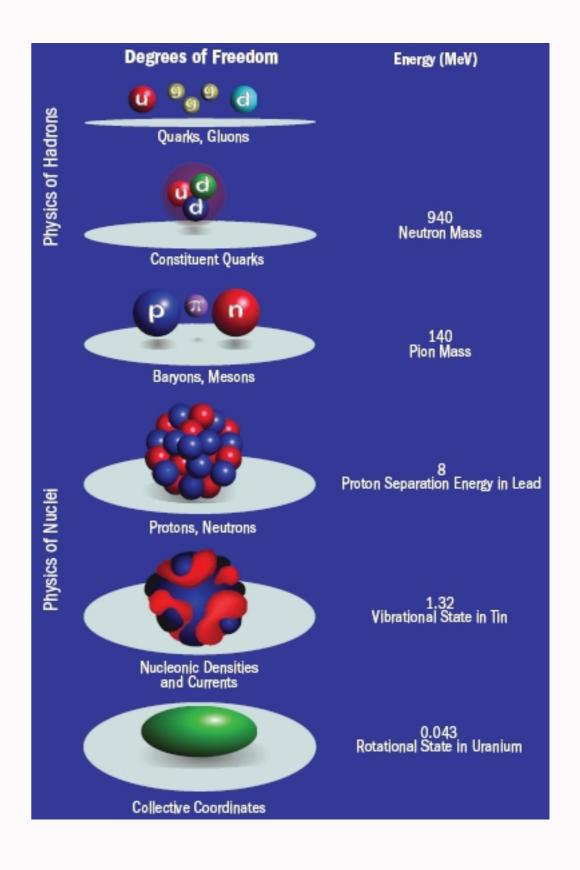
J.-W. Chen & W. Detmold, Phys. Lett. B **625**, 165 (2005):

Structure functions factorize: $F_2^A(x)/A = F_2^N(x) + g_2(A, \Lambda)f_2(x, \Lambda)$ $g_2(A, \Lambda) = \frac{1}{A} \langle A | (N^{\dagger}N)^2 | A \rangle_{\Lambda}$

> J.-W. Chen, W. Detmold, JEL, A. Schwenk, arXiv:1607.03065 [hep-ph] (2016):

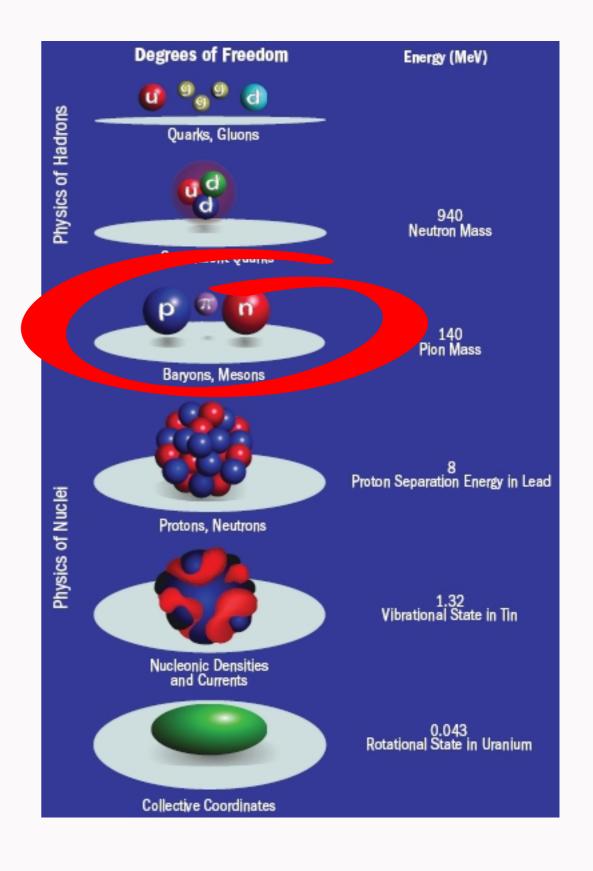
$$a_2(A, x > 1) = \frac{g_2(A, \Lambda)}{g_2(2, \Lambda)} \Rightarrow \frac{dR_{EMC}}{dx} \propto a_2.$$

Quantum chromodynamics (QCD) is ultimately responsible for strong interactions.



Quantum chromodynamics (QCD) is ultimately responsible for strong interactions.

Nucleons are the relevant degrees of freedom for low-energy nuclear physics.



Nuclei are strongly interacting many-body systems.

1. How do we solve the many-body Schrödinger equation?

 $H\left|\Psi\right\rangle = E\left|\Psi\right\rangle$

 $2^{A}\binom{A}{Z}$ coupled differential equations in 3A - 3 variables.

 $^{4}\text{He} \rightarrow 96 \text{ equations in 9 variables}$ $^{12}\text{C} \rightarrow 3784704 \text{ equations in 33 variables}$

2. What is the Hamiltonian?

Nuclei are strongly interacting many-body systems.

1. How do we solve the many-body Schrödinger equation? QMC methods!

 $H\left|\Psi\right\rangle = E\left|\Psi\right\rangle$

 $2^{A}\binom{A}{Z}$ coupled differential equations in 3A - 3 variables.

 $^{4}\text{He} \rightarrow 96 \text{ equations in 9 variables}$ $^{12}\text{C} \rightarrow 3784704 \text{ equations in 33 variables}$

2. What is the Hamiltonian? \rightarrow Chiral EFT!

QMC methods in two lines:

$$H |\Psi\rangle = E |\Psi\rangle$$
$$\lim_{\tau \to \infty} e^{-H\tau} |\Psi_T\rangle \to |\Psi_0\rangle$$

QMC methods in more than two lines:

J. Carlson et al, RMP **87**, 1067 (2015).

QMC Methods - Variational Monte Carlo (VMC) Method

- 1. Guess a trial wave function Ψ_T and generate a random position: $\mathbf{R} = \mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A$.
- 2. Use the Metropolis algorithm to generate new positions **R'** based on the probability $P = \frac{|\Psi_T(\mathbf{R'})|^2}{|\Psi_T(\mathbf{R})|^2}$. (Yields a set of "walkers" distributed according to $|\Psi_T|^2$).
- 3. Invoke the variational principle: $E_T = \frac{\langle \Psi_T | H | \Psi_T \rangle}{\langle \Psi_T | \Psi_T \rangle} > E_0$.

QMC Methods - Diffusion Monte Carlo Method

- The wave function is imperfect: $|\Psi_T\rangle = \sum_{i=0}^{\infty} \alpha_i |\Psi_i\rangle$.
- Propagate in imaginary time to project out the ground state $|\Psi_0\rangle$.

$$\begin{aligned} |\Psi(\tau)\rangle &= \mathrm{e}^{-(H-E_{T})\tau} |\Psi_{T}\rangle \\ &= \mathrm{e}^{-(E_{0}-E_{T})\tau} [\alpha_{0} |\Psi_{0}\rangle + \sum_{i\neq 0} \alpha_{i} \mathrm{e}^{-(E_{i}-E_{0})\tau} |\Psi_{i}\rangle]. \end{aligned}$$

QMC Methods - Diffusion Monte Carlo Method

- The wave function is imperfect: $|\Psi_T\rangle = \sum_{i=0}^{\infty} \alpha_i |\Psi_i\rangle$.
- Propagate in imaginary time to project out the ground state $|\Psi_0\rangle$.

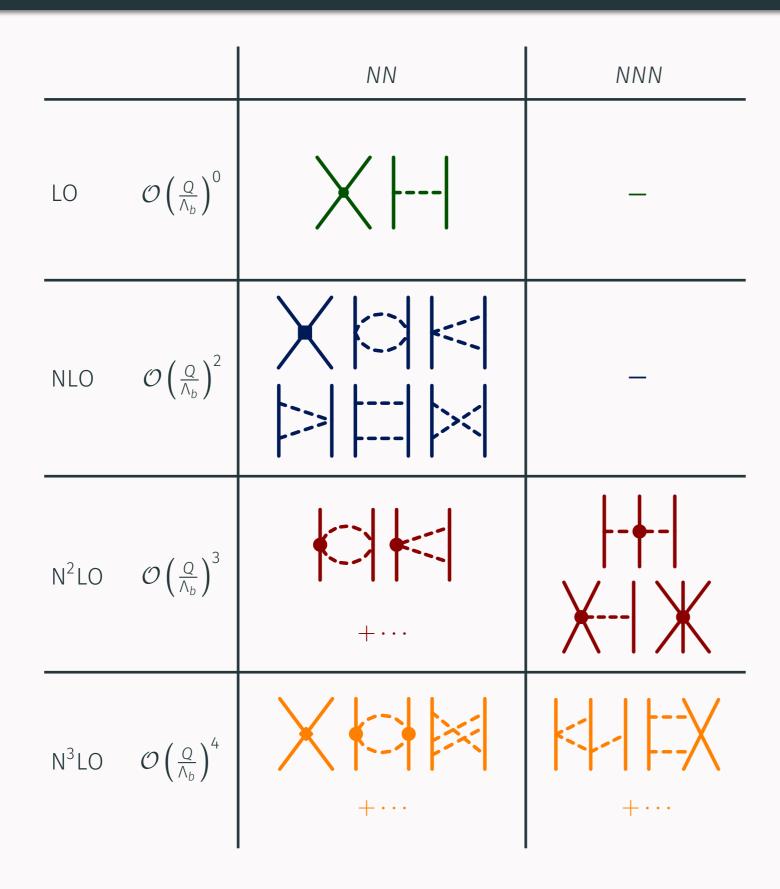
$$\begin{split} \left| \Psi(\tau) \right\rangle &= e^{-(H - E_T)\tau} \left| \Psi_T \right\rangle \\ &= e^{-(E_0 - E_T)\tau} \left[\alpha_0 \left| \Psi_0 \right\rangle + \sum_{i \neq 0} \alpha_i e^{-(E_i - E_0)\tau} \left| \Psi_i \right\rangle \right]. \\ \left| \Psi(\tau) \right\rangle \xrightarrow{\tau \to \infty} \left| \Psi_0 \right\rangle. \end{split}$$

Of course, the nuclear Hamiltonian is complicated.

$$H = \sum_{i=1}^{A} \frac{\mathbf{p}_{i}^{2}}{2m_{i}} + \sum_{i< j}^{A} V_{ij} + \sum_{i< j< k}^{A} V_{ijk} + \cdots$$

Where should it come from?

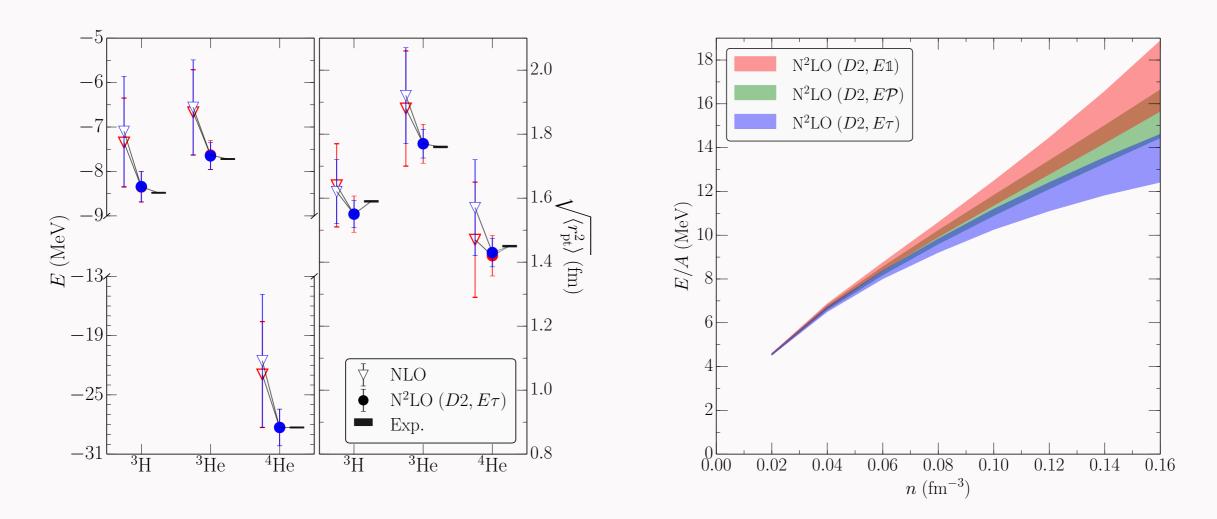
Chiral EFT



- Chiral EFT: Expand in powers of Q/Λ_b . $Q \sim m_{\pi} \sim 100$ MeV $\Lambda_b \sim 500$ MeV
- Long-range physics: π exchanges.
- Short-range physics: Contacts × LECs.
- Many-body forces & currents enter systematically.

Results

A simultaneous description of properties of light nuclei, *n*-α scattering and neutron matter is possible. Uncertainty analysis as in E. Epelbaum et al, EPJ **A51**, 53 (2015).



JEL et al, PRL **116**, 062501 (2016)

Uncertainty Estimates

Define

$$\begin{split} X^{(i)} &= X^{(0)} + \Delta X^{(2)} + \dots + \Delta X^{(i)}, \\ \Delta X^{(2)} &= X^{(2)} - X^{(0)}, \ \Delta X^{(i)} = X^{(i)} - X^{(i-1)}, \ i > 2. \end{split}$$

Expected size

$$\Delta X^{(i)} = \mathcal{O}(Q^i X^{(0)}).$$

Then,

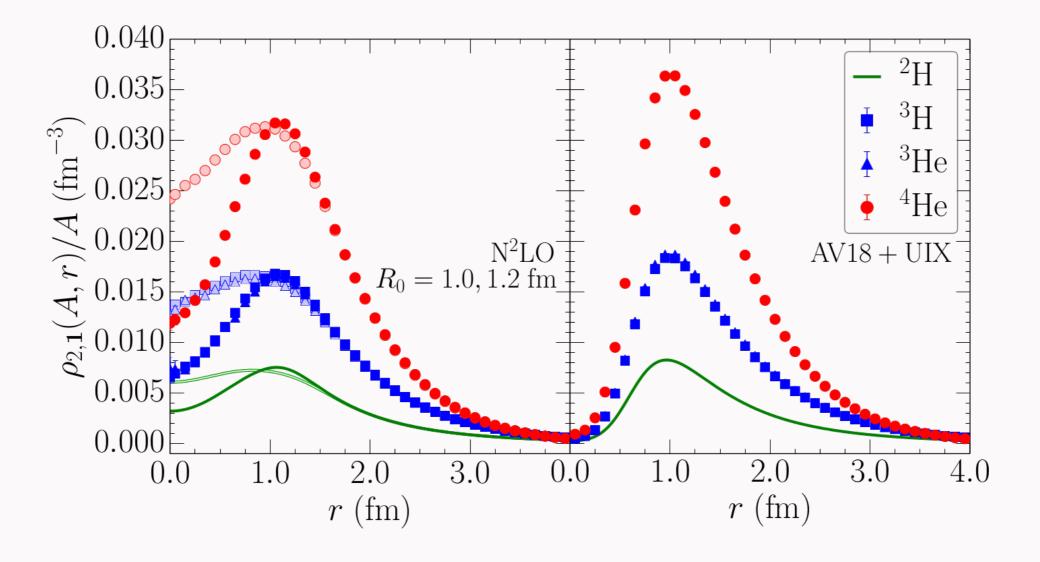
 $\delta X^{(0)} = Q^2 |X^{(0)}|, \ \delta X^{(i)} = \max(Q^{i+1} |X^{(0)}|, Q^{i+1-j} |\Delta X^{(j)}|), \ 2 \leq j \leq i$

 $Q = \max(p/\Lambda_b, m_{\pi}/\Lambda_b).$

Two-Body Distribution Functions (g_2)

$$g_2(A,\Lambda) = \rho_{2,1}(A,r=0)/A, \ \rho_{2,1}(A,r) \equiv \frac{1}{4\pi r^2} \left\langle \Psi_0 \right| \sum_{i < j} \delta(r - r_{ij}) \left| \Psi_0 \right\rangle$$

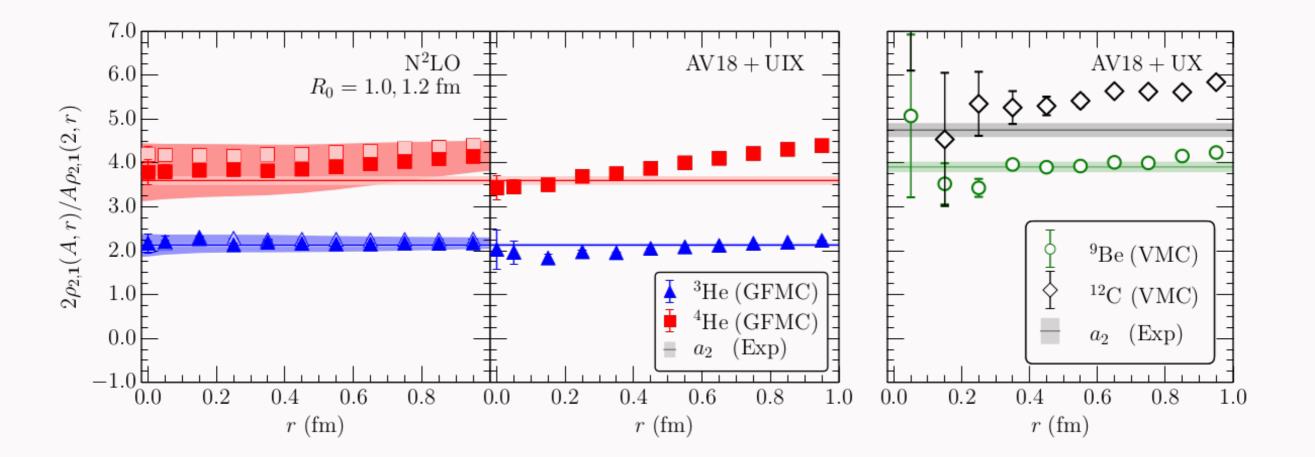
Scale and scheme dependent



SRC Correlation Factors

$$a_2 \equiv \lim_{r \to 0} \frac{2\rho_{2,1}(A, r)}{A\rho_{2,1}(2, r)}$$

Scale and scheme *independent*!



Detailed comparison of experiment and theory for a_2 .

N²LO (
$$R_0 = 1.0 - 1.2 \text{ fm}$$
)AV18+UIXExp³H2.1(2) - 2.3(3)2.0(4)³He2.1(2) - 2.1(3)2.0(4)⁴He3.8(7) - 4.2(8)3.4(3)

- EFT explains the linear relationship between dR_{EMC}/dx and a_2 , and further implies that a_2 is calculable in "traditional nuclear physics".
- Two-body distributions are scheme and scale dependent, but ratios are observable.
- Our results suggest that EFT can shed light on the existence of a_3 .

What does "short-range" in SRC mean?

Should *a*² be called the SRC factor?