## Scattering theory of transport, III

## 1. Wigner's scattering time.

In this problem we consider scattering on a resonance in the time domain. For simplicity, we shall focus on the one-channel system discussed in Lecture 4: a delta-function potential  $U(x) = \alpha \delta(x-a)$  on a half-line x > 0 with a hard-wall boundary condition at x = 0. Suppose that this system has a quasi-bound state with the complex energy  $E = E_0 - \frac{i}{2}\gamma$ .

(i) Consider an incident wavepacket with energy  $\vec{E}$  close to  $E_0$  and narrow dispersion  $\Delta \epsilon \ll \gamma$ (this corresponds to large spatial width  $\Delta x \gg \hbar v_F/\gamma$ ). Show that after scattering the wavepacket comes out with a time lag

$$\Delta \tau = 2\hbar (d\delta(E)/dE)_{E=\bar{E}} \tag{1}$$

where  $\delta(E)$  is scattering phase defined by  $e^{2i\delta} = -\frac{E-E_0-\frac{i}{2}\gamma}{E-E_0+\frac{i}{2}\gamma}$ .

Analyze the time lag as a function of energy for this resonance model. What is the maximal value of  $\Delta \tau$ ?

(ii) Now consider a wide wavepacket,  $\Delta E \gg \gamma$ , which corresponds to narrow spatial width  $\Delta x = \hbar v / \Delta \epsilon \ll \hbar v_F / \Gamma$ . Show that the reflected wavepacket consists of a non-delayed and delayed parts. Find the probability for the scattering particle to be delayed by a time  $\tau$ . Interpret the result.

## 2. Multichannel transport

For an electron system having two external leads with N channels in each lead, left and right, show that the scattering matrix can be written as

$$S = \begin{pmatrix} r_{11} & t_{21} \\ t_{12} & r_{22} \end{pmatrix} \tag{2}$$

where  $r_{11}$ ,  $r_{22}$ ,  $t_{21}$ ,  $t_{12}$  are  $N \times N$  matrices describing reflection and transmission among N channels of each lead.

(i) Evaluate electric current between the leads as a difference of the left-to-right and right-to-left currents. Show that the total current is equal to

$$I = \frac{e}{\hbar} \int \frac{d\epsilon}{2\pi} \left[ f_1(\epsilon) \operatorname{Tr} \left( t_{21} t_{21}^{\dagger} \right) - f_2(\epsilon) \operatorname{Tr} \left( t_{12} t_{12}^{\dagger} \right) \right]$$
(3)

where  $f_{1,2}(\epsilon)$  are particle energy distributions in reservoirs. (Note that  $\text{Tr}AA^{\dagger} = \sum_{i,j=1...N} |A_{ij}|^2$ .)

(ii) Apply the result (3) to current between two reservoirs which supply equilibrium distributions, with external voltage controlling the chemical potential difference,  $\mu_L = \mu_R + eV$ . Use unitarity of S to show that  $\text{Tr}(t_{21}t_{21}^{\dagger}) = \text{Tr}(t_{12}t_{12}^{\dagger})$ . Take the limit  $V \to 0$  to obtain Ohm's law I = GV with conductance given by the multi-channel Landauer formula

$$G = \frac{e^2}{h} \operatorname{Tr} \left( t_{21} t_{21}^{\dagger} \right)$$
 (4)

What are the maximal and minimal possible values of G?

## 3. Drude-Lorentz conductivity of a Fermi gas

Consider electrons at density n in a 2D metal at zero temperature in the presence of randomly placed point-like scatterers,  $U(r) = \sum_{i} u\delta(\mathbf{r} - r_i)$ . The concentration of scatterers is  $n_{imp}$ .

a) Treating scattering in the Born approximation, estimate the transition rate  $w(\mathbf{p}', \mathbf{p})$  in the collision integral of the Boltzmann equation. Find the mean free path and evaluate Drude conductivity.

b) In the presence of a finite magnetic field, find the components of the resistivity tensor  $\rho_{xx}(B)$  and  $\rho_{xy}(B)$ .

See: C. W. J. Beenakker and H. van Houten review cond-mat/0412664 "Quantum transport in semiconductor nanostructures" pages 8-11.