Recitation 7 - Markow Jump Procenes and their Simulation

I-Markov Jump Procen

So far, we have been learning about Markov processes that one antinuous in both their state and in time. Markov Jump Processes are continuous -time processes with a discrete state space  $\mathcal{S} = \{P_i, P_2, \dots, Q_n\}$ . The number of states may be finite or infinite. At any point in time, the system occupies some state  $\mathcal{K}(t) \in \mathcal{S}$ . Hore are some examples:

- 1) A chemical reaction  $A \rightleftharpoons B$ . State space  $S = \{A, B\}, n = 2$ .
- 2) A random walk on a 1D-lattice State at any time is the particle position. Thus, S = Z,  $n = \infty$
- 3) A collection of particles hopping on a Anite ID talkies of size L, with birth, death, and hard-care exclusion.

You an think of this as a microtuble in a solution of microhubate-birding poteins

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Each stak is represented by a vector of accurpancy numbers  $x(t) = (n_1(t), n_2(t), ..., n_L(t))$ , where  $n_i = 0$  or f. Thus  $\mathcal{S} = \{0, f\}^L$ .

A trajectory of a Morkov procen is specified by a sequence of states {x; } and a corresponding sequence of dwell times:

$$trajectory: ((x_0, T_0), (x_1, T_1), ..., (x_p, T_n), t_p), \quad where x_i \in \mathcal{S}$$

....

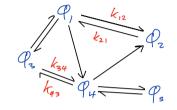
The "jumps" occur at times  

$$t_i \equiv t_{i-1} + T_{i-1}$$
, for i>0.  
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As discussed in lecture, the state occupation polabilities follow a matter equation  $Q_t P(Q_i, t) = \sum_{j \neq i} W(Q_j \rightarrow Q_i) P(Q_j, t) - W(Q_i \rightarrow Q_j) P(Q_i, t)$ where  $W(Q_i \rightarrow Q_i)$  is the transition rate from state  $Q_j$  to state  $Q_i$ , defined in terms of the conditional polability

$$\lim_{\Delta t \to 0} \frac{P(\varphi_i, t+\Delta t | \varphi_j, t)}{\Delta t} = \mathcal{W}(\varphi_j \to \varphi_i) \qquad \text{for } i \neq j$$

In the antext of chemical reachers, the state space is often regresented as a graph, with arraws denothing allowed transitions, and the concernating rates indicated on the arrows  $(k_{ij} \equiv W(q, \rightarrow q_i))$ .



- <u>2-Sinculating a Markov Jump Procen</u> <u>2-1</u> Dwell times and transviblen probabilities To simulate a Markov jump process requires a rule or choosing the dwell time and the rest state toonsviblen given the current state. Duell times to at use to be to be duell times and choosing the dwell time and
- Duell times: We showed in lecture that the dwell times are expandentially distributed according to the escape rate

$$\Gamma(q_i) = \sum_{j \neq i} W(q_i \rightarrow q_j),$$

so that

$$T_i \sim Exp(r(q_i)), \quad p(\tau_i) \propto r_i e^{-t}$$

Next shate : we showed in lecture that, upon escaping state i, the probability that the system lards in state ; is:

$$P_{i \to j} = \frac{\omega(\varphi_i \to \varphi_j)}{\Gamma(\varphi_i)}$$

Togethor, these define an algorithm to exactly sample the trajecturies of a jump process, often called the Gillespie algorithm (though it is much obser than Gillespie).

## 2.2 The Gillespie algorithm (Direct method)

The Gillespie algorithm, developed by Doob and others in ~ 1945 and popularized by Gillespie in 1976-77, generates the next state and jump time of a Markov jump procen as follows:

1) Initialize the time and state variables  $t = t_0$  and  $x = x_0$ . 2) Compute the escape rate  $r = \underset{q \neq x}{\underset{q \neq x}{\underset{p = \atop{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p \neq x}{\underset{p = \atop{p = \atop{p \neq x}{\underset{p = \atop{p = \atop{p = x}{\underset{p = \atop{p \neq x}{\underset{p = \atop{p \neq x}{\underset{p = \atop{p = \atop{p \neq x}{\underset{p = x}{\underset{p$ 

$$\Rightarrow$$
 For step (2+), deck that  $\frac{1}{r} \log(\frac{1}{4}) \sim \exp(r)$ .

=> For step (S), we wish to pick the next state with a probability poportional to the rate of jumping from x to that state. See the below picture:

$$w(x \rightarrow 1) \qquad w(x \rightarrow 2) \qquad w(x \rightarrow 3) \qquad w(x \rightarrow 4) \qquad w(x \rightarrow 5)$$

=) Why are the dwell time and vext reaching independent? Go back to taday's lecture notes and check that the joint density  $P(\tau,j)$  is  $P(\tau,j) = W(x \rightarrow j) \exp\left[\frac{1}{2} - w(x \rightarrow e)\tau\right]$ 

This algorithm is exact and is sometimes called the "direct" method. It is "rejection free", in the sense that every random number generation is associated with a state transition. It is an example of a "kinetic Mante Carlo" algorithm.

In many situations (e.g. large reaction network, many particles on a lattice), the number of possible state transitions is very large. This means the escape rates are very large, and the dwell times are very short, so very many iterations of the algorithm are needed to elapse a reasonably long time interval. It also means that computation of the escape rates is

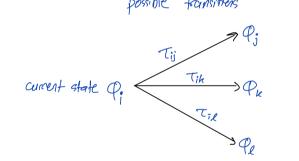
itself costly as it takes a time which scales linearly with the number of possible transitions. There are some optimizations that can be done when many of the reachians are independent (e.g. for-oway particles on a lattice). These rely an an alternative description of a Markov jump procen. 2.3 "First-reachen" and "Next reachen" methods In the above, we characterize a Markov jump poten in terms of the distributions governing the dwell time and the next state transition choice. An equivalent description is the following: suppose that the system is in state Q. at some time. To each possible escape reacher Q; -> Q; is associated a stopwortch, and a waithing time

 $\tau_{ii} \sim \exp(w(\varphi_i \rightarrow \varphi_i))$ 

is drawn. The stop watches are starfed. The first time a stopwatch hits Tij Br some j, reachen  $\rho \rightarrow \rho$  is executed. The waiting times are then drown onew, and the stopwortdaes restarted.

Schematically:

possible transitions



Dwell time: min (Tij) Next state:  $\phi_p$ , where  $p = \arg(min(T_{ij}))$ 

To see that this is equivalent to the above, the Allowing must be shown: Consider a collection of n independent exponentially-distributed radion variables with rates  $k_1, k_2, ..., k_n$ ; that is,  $z \tau, ..., \tau_n z$  with  $\tau_i \sim Exp(k_i)$ (I) The minimum  $T = \min \{T_i\}$  is Erp(r) distributed with  $r = E_i k_i$ .

(2) The probability that T: is the smallest among the ET; 3 is:

$$\mathbb{P}\left[\operatorname{argmin}_{i} \{t_{j}\} = i\right] = k_{i/r}$$

For (1), it suffices to prove the statement for n=2, because it may then be applied pairwise  $\min(\tau_1, \tau_2, ..., n) = \min(\min(\tau_1, \tau_2), \tau_3, ..., \tau_n) = \min(\min(\min(\min(\tau_1, \tau_2), \tau_3), \tau_4, ..., \tau_h))$ and so on. For n=2, we have  $P(\min(\tau_1, \tau_2) > t) = P(\tau_1 > t) P(\tau_2 > t) = \int_t^{\infty} ds \ k, \ e^{-k_1 s} \int_t^{\infty} ds' \ k_2 \ e^{-k_2 s'}$   $= e^{-k_1 t - k_2 t}$  $= 1 - P(\min(\tau_1, \tau_2) < t)$ 

we thus identify the CDF as that of an Exp(ki+kz) R.V.

$$\mathcal{P}\left[\operatorname{argmin}_{j} \mid t_{j} \right] = i = \left\langle \prod_{p \neq i} \mathbb{1}\left(\tau_{k} > \tau_{i}\right) \right\rangle$$

where I(A) = I if event A occurs and zero othernise.

The terms within the product one not independent: If I tell you  $T_{k} > T$ ; for some k, this fells you comething about the smallness of T; so that  $(T_{m} > T;)$  is now more likely. To resolve this, we condition on T; and use the former rule E(X) = E[E[X|Y]]:

$$\begin{aligned} \mathcal{P}\left[ \underset{j \neq i}{\operatorname{argmin}} \left\{ \tau_{j} \right\} = i \right] &= \left\langle \left\langle \tau_{p \neq i} \right| \left\{ \tau_{p > \tau_{i}} \right\} \left| \tau_{i} \right\rangle \right\rangle_{\tau_{i}} \\ &= \left\langle \tau_{p \neq i} \right| \left\{ \left\{ \tau_{p > \tau_{i}} \right\} \right\}_{\tau_{i}} \\ &= \left\langle \tau_{p \neq i} \right| \left\{ e^{-k_{p} \tau_{i}} \right\}_{\tau_{i}} \\ &= k_{i} \cdot \int_{p \neq i} \left\{ e^{-k_{p} \tau_{i}} \right\}_{\tau_{i}} \\ &= k_{i} \cdot \int_{p \neq i} \left\{ e^{-k_{p} \tau_{i}} \right\}_{\tau_{i}} \end{aligned}$$

2.3.1 = "First reachian" method

Based a the above Gillespit proposed the following algorithm: at each step draw an  $E + p(w(x \rightarrow \varphi))$  variable  $T_{\varphi}$  for every  $\varphi \in \mathcal{S}$ . Set  $t \rightarrow t + \min \{\tau_{\varphi}\}$ and  $x \longrightarrow \arg \min_{\varphi} \{\tau_{\varphi}\}$ . This algorithm is obviously less efficient because it draws up to  $|\mathcal{S}|$  reaction variables at each step and identifies their minimum. Its usefulnum is thus mostly perlogging.

图

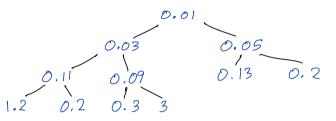
2.3.2: "Next reachien" method (Gibson and Bruck, 2000).

The idea here is to draw all the possible transition times of t; 3 and to store them in an efficient data structure that maintains their order.

At each step, the smallest T; transition is executed, and only those reactions which are affected are redrawn.

The original implementation is specialized to chemical reactions, so we list here only the menin idea:

⇒ The transition times are stored in a min heap data structure (usually indexed)



=> The rest reacher is always the root and can be found in O(1) time.

 $\Rightarrow$  An additional data structure keeps track of the dependency structure between the elements (e.g. nearest neighbor sites an a lathic).

Example: spins on a lattice (e.g. a kinetit Ising model):

Consider a model with spin  $s = \pm 1$  on a lattice  $\mathbb{Z}^d$ . Spin i Plips at a rate which is a function of s; and the nearest neighbors of i, denoted  $\sigma(i)$ :

 $Aip \ rate = w(\S_{S}: j \in \sigma(i)).$ 

This can be simulated as follows:

- 1) Initialize the spin lattice. For each  $i \in L^d$ , draw the flip time of spin i according to  $\tau_i \sim \exp(w(\bar{z}s_j:j\in\tau_{(i)}))$
- 2) Organize the L<sup>d</sup> parts (i, T;) into a min heap sorted according to T;. Creak a data structure of pointers to the tree for each i.
- 3) Set the simulation time to  $t = \tau_j$ , where  $j = \operatorname{orgmin} \xi \tau_i \xi$ .
- 4) Flip spin j.
- 5) For each  $i \in \sigma(j)$  and for j itself, redraw  $\tau_i \Rightarrow \tau_i \rightarrow t + \Delta \tau_i$ , where  $\Delta \tau_i \sim E_{tp}(w(S_h: h \in \sigma(i)))$ . Re-organize the heap. Return to step 3.

The clowest step is the time to reorganize the heap, which is  $O(\log N)$  instead of O(N). This is a big improvement over the direct method

2.4 Disorete-time (rejection) algorithms