$$\lim_{n \to \infty} \underset{i}{\overset{k}{\approx}} M_{i-1} \Delta t^2 = 0, \qquad \lim_{n \to \infty} \underset{i}{\overset{k}{\approx}} \mathcal{M}_{i-1} dW^n = 0 \quad \text{for } n > 2, \quad \text{etc}$$

Let's check (1). Define $I_n = \stackrel{\circ}{\underset{i=1}{\overset{\circ}{\leftarrow}}} M_{i-1} \left(\Delta W_i^2 - \Delta t \right)$ Notice n t

$$\lim_{n \to \infty} I_n = \lim_{n \to \infty} \bigoplus_{i=1}^n \mathcal{M}_{i-1} \Delta \mathcal{W}_i^2 - \int_0^t ds \mathcal{M}(s)$$

So to check equation F, it suffices to show that In conveges to zero. We will show this in the mean square sense: i.e. we will show

$$\lim_{n \to \infty} \langle I_n \rangle = 0, \quad \lim_{n \to \infty} \langle I_n^2 \rangle = 0$$

Indeed:

$$\langle \mathbf{I}_{n} \rangle = \sum_{i}^{n} \langle \mathcal{M}_{i-1} \rangle \langle \Delta \mathcal{W}_{i}^{2} - \Delta t \rangle = 0$$

$$\langle \mathbf{I}_{n}^{2} \rangle = \left\langle \left[\sum_{i=1}^{n} \mathcal{M}_{i-1} \left(\Delta \mathcal{W}_{i}^{2} - \Delta t \right) \right]^{2} \right\rangle$$

$$= \sum_{i=1}^{n} \left\langle \mathcal{M}_{i-1}^{2} \left(\Delta \mathcal{W}_{i}^{2} - \Delta t \right)^{2} \right\rangle + \sum_{i>j}^{n} 2 \left\langle \mathcal{M}_{i-1} \mathcal{M}_{j-1} \left(\Delta \mathcal{W}_{j}^{2} - \Delta t \right) \left(\Delta \mathcal{W}_{i}^{2} - \Delta t \right) \right\rangle$$

$$= \sum_{i=1}^{n} \left\langle \mathcal{M}_{i-1}^{2} \right\rangle \left\langle \left(\Delta \mathcal{W}_{i}^{2} - \Delta t \right)^{2} \right\rangle$$

$$= \sum_{i=1}^{n} \left\langle \mathcal{M}_{i-1}^{2} \right\rangle \left\langle \left(\Delta \mathcal{W}_{i}^{2} - \Delta t \right)^{2} \right\rangle$$

Now use $\langle (\Delta w;^2 - \Delta t)^2 \rangle = \langle \Delta w;^4 \rangle - 2 \langle \Delta w;^2 \rangle \langle \Delta t \rangle + \Delta t^2 = 3 \Delta t^2 - 2 \Delta t^2 + \Delta t^2 = 2 \Delta t^2$ Thus: $\langle In^2 \rangle = 2 \stackrel{n}{\underset{i=1}{\overset{n}{=}}} \langle M_{ini}^2 \rangle \frac{t^2}{n^2} \stackrel{n \to \infty}{\longrightarrow} 0$ The other rules can be prown similarly. Using these rules, we can manipulate infinitesmals like dW and dt and Taylor expand with respect to them, knowing when to terminate the socies.

Let
$$dx = \mu(x)dt + \tau(x)dW$$

Consider $f(x)$ for some smooth f . Then
 $df = f'(x)dx + f''(x)\frac{dx^2}{2} + \frac{f'''(x)}{6}\frac{dx^3}{6} + \cdots$
 $= f'(x)dx + \frac{f''}{2}(\mu(x)dt + \tau(x)dW)^2 + \cdots$

$$= f'(x) dx + \frac{\mu''}{2} (\mu(x)^2 dt^2 + 2\mu(x) dt dW + \sigma^2 dt)$$

Dividing by dt:

$$\dot{f} = f'(x)\dot{x} + \frac{\tau^2}{2}f''$$

we have recovered Ito's Granub.

Example: Rom Stratonouich to Ito Suppose I have a stratonouich SDE $dx = \mu(x)dt + \sigma(x) \circ dW$, wiresponding to the integral $x(t) = x_0 + \int_0^t ds \,\mu(x(s)) + \int_0^t dW(s) \circ \sigma(x(s))$ where follow denotes the stratonouich integral, defined by

$$\int_{0}^{t} dW(s) \circ \sigma(x(s)) \equiv \lim_{n \to \infty} \sum_{i=1}^{n} \sigma\left(\frac{x_{i} + x_{i-1}}{2}\right) \Delta W_{i}$$

1 wish to rewrite this as on I to SDE

$$dx = \tilde{\mu}(x)dt + \tilde{\sigma}(x)dW$$

where $\tilde{\mu}, \tilde{\sigma}$ are to be given in terms of μ, σ . Let write $x_i = x_{i+1} + D \pi$; and expand the gradient integral w.r.t $D x_i$:

$$\sum_{i=1}^{n} \sigma\left(\frac{x_{i}+x_{i-1}}{2}\right) \Delta W_{i} = \sum_{i=1}^{n} \left[\sigma(x_{i})+\sigma'(x_{i})\frac{\Delta x_{i}}{2}+\sigma''(x_{i})\frac{\Delta x_{i}^{2}+\cdots}{4}\right] \Delta W_{i}$$

$$\sum_{i=1}^{n} \sigma\left(\frac{x_{i} + x_{i-1}}{2}\right) \Delta W_{i} = \sum_{\substack{i=1 \\ j \neq i}}^{n} \sigma(x_{i}) \Delta W_{i} + \sum_{\substack{i=1 \\ j \neq i}}^{n} \sigma'(x_{i}) \left(\tilde{\mu}(x_{j}) \Delta t + \tilde{\sigma}(x_{i}) \Delta W_{i}\right) \Delta W_{i}$$

Taking $n \rightarrow \infty$, the ferm $\Delta t \Delta W' \sim \mathcal{O}(\Delta t^{3/2})$ is subleading, $\Delta W'^2$ becomes $\partial W^2 = dt$, so that $\lim_{n \to \infty} \hat{\nabla}(x_i) \tilde{\nabla}(x_i) \Delta w_i^2 = \int_{-\infty}^{t} ds \nabla'(x(s)) \tilde{\nabla}(x(s))$ Thus

$$\int_{a}^{t} dW(s) \circ \sigma(x(s)) = \int dW(s) \sigma(x(s)) + \frac{1}{2} \int ds \sigma'(x(s)) \widetilde{\sigma}(x(s))$$

Altogether, the Ito expremsion for $x(t)$ is
$$x(t) = x_{o} + \int_{a}^{t} ds \left[\mu(x(s)) + \frac{1}{2} \widetilde{\sigma}(x(s)) \partial_{x} \sigma(x(s)) \right] + \int_{a}^{t} dW(s) \sigma(x(s))$$

Martchning terms, we see $\tilde{\sigma} = \sigma$

$$\tilde{\mu} = \mu + \frac{1}{2} \sigma \partial_x \sigma$$

In summary:

The stratonovich SDE $dx = \mu(x) dt + \sigma(x) \circ dw$ (9) is the same as the Ito SDE $dr = (\mu(r) + \frac{1}{2}\sigma\sigma'(r)) dt + \sigma(r) dw$. (b) Notice that for additive noise $\sigma'(x) = 0$, there are the same. We may take the above as the <u>definition</u> of the Stratonovich SDE, and study its properties. For example, we can prove that strateneousch SDEs doey the standard chart rule: Let f(x) be some smooth Runchlen, with x obeying (a). Apply Ito's lemma to (b): $\hat{f} = \hat{f}' \dot{x} + \frac{1}{2} \sigma^2 f''$ $df = f' \mu dt + \frac{1}{2} f' \sigma \sigma' dt + f' \sigma dw + \frac{1}{2} \sigma^2 f'' dt$ $= \left[f'_{\mu} + \frac{1}{2} \left(f'_{\sigma} \sigma' + \sigma^2 f'' \right) \right] dt + f'_{\sigma} dw$ (\mathbf{v}) This last result is on Ito equality. To convert to Strato, we use Ito SDE: df = milidt + J(f) dw equivalent shalp: $df = (\mu(f) - \frac{1}{2} \sigma_{\overline{p}} \sigma_{\overline{f}}'(f)) dt + \sigma_{\overline{p}}(f) \circ dw$ So that is above how the strato firm $df = \left[f'_{\mu} + \frac{1}{2}(f'_{\sigma}\sigma' + \sigma^{2}f'') - \frac{1}{2}f'_{\sigma}\partial_{p}(f'_{\sigma})\right]dt + f'_{\sigma}(x) \circ dw \quad (\forall k)$ Notice: $\frac{1}{p}\left[f'(x)\sigma(x)\right] = \frac{1}{p}\left[f'(f'(f))\sigma(f'(f))\right]$

$$= \nabla f'' \frac{\partial f^{-}(f)}{\partial f} + f' \nabla' \frac{\partial f^{-}(f)}{\partial f}$$
$$= \frac{1}{P'} \left(\nabla f''_{f} f' \nabla'_{f} \right)$$

Plugging into ** above, we get

 $df = f' \mu \, dt + f' \sigma dw \Rightarrow \dot{f} = f' \dot{x}$

Thus the stratonoxich quarter obeys the standard char rule.