

The Wiener Process

Recap

A Wiener Process or Standard Brownian Motion can be defined in two ways:

- ① $\left\{ \begin{array}{l} \text{A process } \{W(t) : t \geq 0\} \text{ having} \\ \text{i) Continuous paths} \\ \text{ii) Stationary, independent increments} \\ \text{iii) } W(t) \sim N(0, t) \quad \forall t \geq 0 \end{array} \right.$

or

- ② $\left\{ \begin{array}{l} \text{A Gaussian process with continuous paths, mean zero and correlation function:} \\ \langle W(t_1) W(t_2) \rangle_c = \min(t_1, t_2) \end{array} \right.$

Some properties of the Wiener Process

Thm: Suppose W is a Wiener process and let $c > 0$. Then the process

$$X(t) = c^{-1/2} W(ct)$$

is a Wiener process.

Pf: Follows immediately from ②: $X(t)$ is clearly Gaussian.

$$\begin{aligned} \langle X(t_1) X(t_2) \rangle_c &= c^{-1} \langle W(ct_1) W(ct_2) \rangle_c \\ &= c^{-1} \min(ct_1, ct_2) = \min(t_1, t_2) \end{aligned} \quad \square$$

Thm: The process $X(t) = W(c+t) - W(c)$ is a Wiener process independent of $W(t \leq c)$, for any $c > 0$.

Pf: Follows immediately from ① ii and ②. \square

This is the sense in which Brownian motion is Markovian and stationary - it is constantly restarting.

Thm: The process $X(t) = t W(1/t)$ is a Wiener process.

Pf: $X(t)$ is Gaussian and continuous (ignore technicalities at $t=0$). Covariance function:

$$\langle t_1 W(1/t_1) t_2 W(1/t_2) \rangle_c = t_1 t_2 \min(1/t_1, 1/t_2) = \frac{t_1 t_2}{\max(t_1, t_2)} = \min(t_1, t_2).$$

General Brownian motions

We may generalize the Wiener process by adding a drift and a scale to $W(t)$.

A process $X(t)$ defined as

$$X(t) = X(0) + \mu t + \sigma W(t)$$

is called a (μ, σ^2) -Brownian motion (though this terminology is rarely used in physics).

I bring this up because of an important fact:

Thm: If a stochastic process X has continuous paths and stationary independent increments, then X is a Brownian motion as defined above.

The Gaussianity of increments thus comes "for free". This can be intuited from the central limit theorem: If increments are stationary and independent, then any one increment can be partitioned as the sum of many smaller increments that are i.i.d.

Some pathologies of Brownian Paths

1) For any $\varepsilon > 0$, $W(t)$ crosses zero infinitely many times on the interval $(0, \varepsilon)$.

One way to understand this is to use that $tW(1/t)$ is a Wiener process.

It is believable that a Wiener process crosses 0 infinitely many times on the interval $(1/\varepsilon, \infty)$ (we may come back to this when we talk about recurrence of random walks). But if $sW(1/s) = 0$ for infinitely many $s \in (1/\varepsilon, \infty)$, then $W(t) = 0$ for infinitely many $t \in (0, \varepsilon)$.

2) $W(t)$ is nowhere differentiable, with prob 1.

This may be understood from the independence of increments: $W(t+\delta) - W(t)$ and $W(t-\delta) - W(t)$ are independent $\forall t, \delta > 0$.

Another way: consider $\Delta W = W(t+\delta) - W(t)$.

This has variance $\langle \Delta W^2 \rangle = \delta$, hence $\lim_{\delta \rightarrow 0} \frac{\Delta W}{\sqrt{\delta}} \sim \frac{1}{\sqrt{\delta}}$

Computing things: First passage times

Let $\{W(t)\}$ be a Wiener process. Define the first passage time for some $b > 0$:

$$\begin{aligned}\tau_b &= \inf \{t : W_t \geq b\} \\ &= \inf \{t : W_t = b\}.\end{aligned}$$

Let us compute the cumulative distribution function $F(t) = \mathbb{P}[\tau_b \leq t]$.

$$\begin{aligned}\text{Then } \mathbb{P}[\tau_b \leq t] &= \mathbb{P}[\tau_b \leq t, W_t < b] + \mathbb{P}[\tau_b \leq t, W_t > b] \\ &= \mathbb{P}[W_t < b | \tau_b \leq t] \mathbb{P}[\tau_b \leq t] + \mathbb{P}[W_t > b]\end{aligned}$$

The second term simply follows from $W_t \sim N(0, t)$:

$$\mathbb{P}[W_t > b] = \int_b^\infty dx \frac{1}{\sqrt{2\pi t}} e^{-x^2/2t} = \frac{1}{2} \text{Erfc}\left(\frac{b}{\sqrt{2t}}\right).$$

How about the first term? First convince yourself that, by symmetry,

$$\mathbb{P}[W_t < b | \tau_b \leq t] = 1/2 \quad \text{"Reflection Principle"}$$

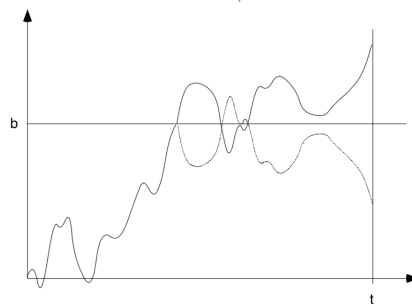
i.e. once you condition on hitting b at some time $\tau_b \leq t$, $W(t)$ is equally likely to be above or below b at subsequent times. Thus:

$$\boxed{\mathbb{P}[\tau_b < t] = 2 \mathbb{P}[W_t > b] = \text{erfc}\left(\frac{b}{\sqrt{2t}}\right)}.$$

($\text{erfc} \equiv 1 - \text{erf}$). The corresponding probability density is

$$\boxed{P_{\tau_b}(t) = \frac{b e^{-b^2/2t}}{\sqrt{2\pi t^3}}} \quad (\text{supported on } \mathbb{R}^+).$$

The reflection principle:



This is called a Levy distribution. It has diverging mean:

$$\langle \tau_b \rangle = \infty$$

But τ_b is finite with prob 1: $\lim_{t \rightarrow \infty} \mathbb{P}[\tau_b < t] = \text{erfc}(0) = 1$.

We say that a Brownian motion in 1D is "null-recurrent": if it crosses zero, it comes back almost surely, but the time to so is infinite on average.

Conditional distributions

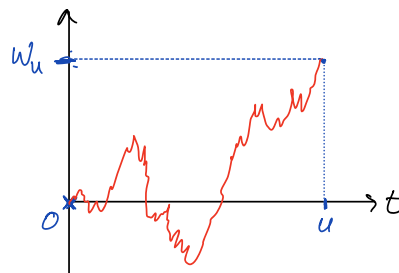
Consider a Wiener Process $W(t)$. Suppose I give you $W(u)$ for some $u > 0$. What is the conditional distribution of $(W(t \leq u) | W(u))$?

The following fact is useful:

Lemma: $W_t - (t/u)W_u$ is independent of W_u for $t \leq u$.

PF: Clearly the quantity is Gaussian, so it suffices to show that the covariance is zero.

$$\begin{aligned} \langle (W_t - \frac{t}{u}W_u) W_u \rangle_c &= \min(t, u) - \frac{t}{u} \langle W_u^2 \rangle_c \\ &= t - \frac{t}{u}u = 0. \end{aligned}$$



This makes life easy:

$$\begin{aligned} \langle W_t | W_u \rangle &= \langle W_t - (t/u)W_u | W_u \rangle + \frac{t}{u}W_u \\ &= \langle W_t - \frac{t}{u}W_u \rangle + \frac{t}{u}W_u \\ &= \boxed{\frac{t}{u}W_u} \end{aligned}$$

This makes sense: the mean is just obtained by linearly interpolating between the two points. The variance is:

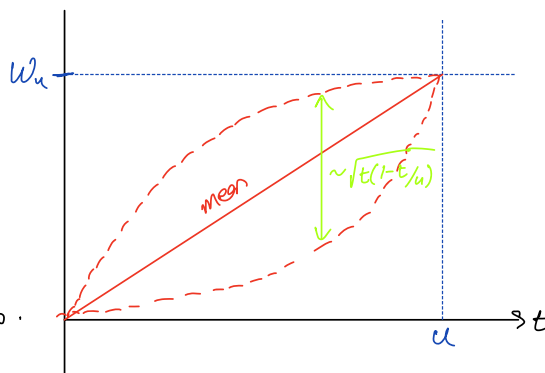
$$\begin{aligned} \langle W_t^2 | W_u \rangle_c &= \langle (W_t - \langle W_t | W_u \rangle)^2 | W_u \rangle \\ &= \langle (W_t - \frac{t}{u}W_u)^2 | W_u \rangle \\ &\stackrel{\text{by lemma}}{=} \langle W_t - \frac{t}{u}W_u \rangle^2 \\ &= t - 2\frac{t}{u}\min(t, u) + \frac{t^2}{u^2}u \\ &= \boxed{t(1 - t/u)} \end{aligned}$$

As expected, variance is smallest near t or u , and maximized at $t = u/2$.

Strangely, this is independent of W_u !

As W_u is made larger, the relative width vanishes:

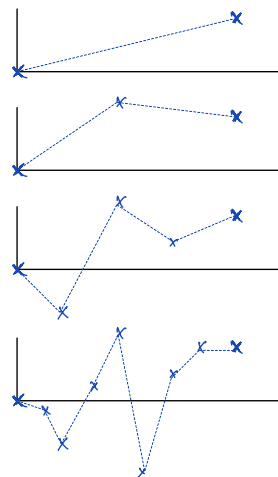
$$\frac{\sqrt{\langle W_t^2 | W_u \rangle_c}}{\langle W_t | W_u \rangle} \leq \frac{u/2}{t/u W_u} \rightarrow 0 \text{ as } W_u \rightarrow \infty.$$



(Maximum variance is always $u/4$)

Note that the result above gives a procedure for sampling a Wiener process on the interval $(0, u)$

- 1) Set $w(0) = 0$.
- 2) Set $w(u)$ by drawing a $\mathcal{N}(0, u)$ number.
- 3) Set $w(u/2)$ by drawing a $\mathcal{N}(\frac{1}{2}w(u), u/4)$
- 4) Repeat for the intervals $(0, u/2)$ and $(u/2, u)$, continue recursively.



The conditional correlation function is, for $s \leq t \leq u$

$$\langle w(s)w(t) | w(u) \rangle_c = \frac{s(u-t)}{u}$$

I leave the proof as an exercise.

The Brownian Bridge

In the special case $u=1$ and $w_u = 0$, the resulting conditional distribution is called a Brownian Bridge. This can be uniquely defined as follows

Def: A standard Brownian Bridge is a Gaussian process $\{X(t): t \in [0, 1]\}$ with continuous paths, mean zero, and $\langle X(s)X(t) \rangle_c = s(1-t)$ for $0 \leq s \leq t \leq 1$.

You can verify that if w is a Wiener process, the following give Brownian bridges:

$$X(t) = w(t) - t w(1)$$

$$Y(t) = (1-t) w\left(\frac{t}{1-t}\right)$$

You can interpret this as a Wiener process "pinned" at $w(1)=0$.

You may verify that a Brownian bridge has the following Fourier representation:

$$B_t = \sum_{k=1}^{\infty} \frac{Z_k}{k} \frac{\sqrt{2} \sin(k\pi t)}{\pi}$$

where the $\{Z_k\}$ are i.i.d $\mathcal{N}(0, 1)$.

The spectral density then falls as $1/k^2 \Rightarrow$ this is like a simple Gaussian field theory with Hamiltonian $\mathcal{H}[\psi] = \int dx (\nabla \psi)^2$ (don't worry if this is meaningless to you).

