# 8.08/8.S308 - Problem Set 1 - IAP 2025

# Due before January 10, 23:59

Problems that are referred to as "graduate" count as bonus problems for undergraduate students. This first problem set is essentially a reminder/test of the Mathematical prerequisites for the course.

# Problem 1—Probabilities

Consider a random variable X of probability density p. Then p(x)dx is the probability that the random variable X takes a values in [x, x + dx], as  $dx \to 0$ . The  $n^{\text{th}}$  moment of X is denoted by  $m_n = \langle X^n \rangle = \int dx x^n p(x)$ . The generating function of the moments of p is  $Z(h) = \langle e^{hX} \rangle$ . It satisfies

$$\langle X^n \rangle = \left. \frac{d^n Z}{dh^n} \right|_{h=0} \qquad \text{so that} \qquad Z(h) = \sum_{n \ge 0} \langle X^n \rangle \frac{h^n}{n!} \,.$$
 (1)

The function  $W(h) = \ln Z(h)$  is the generating function of the cumulants (also called "connected moments") of p. By definition, the  $n^{\text{th}}$  cumulant  $\kappa_n$  of p is  $\kappa_n = \frac{d^n}{dh^n} W(h) \Big|_{h=0}$ . We will use the notation  $\kappa_n = \langle X^n \rangle_c$ , where c stands for "cumulant" or "connected". One thus has  $W(h) = \sum_{n \ge 1} h^n \langle X^n \rangle_c / (n!)$ .

**1.1)** Determine the  $m_n$ 's and  $\kappa_n$ 's for  $p(x) = \exp(-|x|)/2$ .

**1.2)** For an arbitrary p(x), show that  $\kappa_1 = m_1$  and  $\kappa_2 = m_2 - m_1^2$ . Find similar relations for  $\kappa_3$  in terms of  $m_3$ ,  $m_2$  and  $m_1$ , and for  $\kappa_4$  in terms of  $m_4$ ,  $m_3$ ,  $m_2$  and  $m_1$ . For this question, you may want to us that  $\ln(1+u) \simeq u - \frac{u^2}{2} + \frac{u^3}{3} - \frac{u^4}{4} + \mathcal{O}(u^4)$  as  $u \to 0$ .

**1.3)** Show that, for an even p(x), the relationship between  $\kappa_4$  and the moments simplifies into  $\kappa_4 = m_4 - 3m_2^2$ . Conclude that, if p(x) is a Gaussian with vanishing mean,  $\langle X^4 \rangle = 3 \langle X^2 \rangle^2$ . The fourth cumulant is important in that it shows that cumulants are *not* the moments of the centered random variable  $X - \langle X \rangle$ .

#### Problem 2—Fourier transforms and series

Let  $f_n$  be a function defined on an N-site lattice, n = 1, ..., N, with N assumed to be even. We denote the lattice spacing by a so that L = Na is the total length of the lattice. We define

$$\tilde{f}_q = \sum_{n=1}^N e^{iqna} f_n \ . \tag{2}$$

**2.1)** Show that if  $q = \frac{2\pi k}{Na}$ , with  $k = -\frac{N}{2} + 1, \ldots, \frac{N}{2}$ , then  $f_n = \frac{1}{N} \sum_q \tilde{f}_q e^{-iqna}$ . Bonus: Show that  $\sum_q e^{iq(k-n)a} = N\delta_{k,n}$ . If short of time, feel free to use this identity without proving it.

Note that such Fourier transforms are defined up to an arbitrary normalization factor A through

$$\tilde{f}_q = \frac{1}{A} \sum_{n=1}^{N} e^{iqna} f_n; \quad \text{and} \quad f_n = \frac{A}{N} \sum_q e^{-iqna} \tilde{f}_q .$$
(3)

This is reflected in the diversity of conventions that are commonly found in the literature.

**2.2)** We denote x = na and take the  $N \to \infty$ ,  $a \to 0$  limits, with L = Na kept fixed. To this end, we adopt the convenient convention  $A = \frac{1}{a}$ . This is the limit of a continuous but finite interval. Express  $\tilde{f}_q$  as an integral involving f(x), using the convergence of the Riemann sum  $\sum_n ag_n \sim_{N\to\infty} \int dxg(x)$ . How does one obtain f(x) if  $\tilde{f}_q$  is given? What are the acceptable values of q?

**2.3)** We now consider  $N \to \infty$  with L/N = a fixed. This is the limit of an infinite lattice. Show that, in this limit,  $f_n = a \int_{-\pi/a}^{\pi/a} \frac{dq}{2\pi} \tilde{f}_q e^{-iqna}$ . (We are back to the convention A = 1.)

**2.4)** Let  $f(\tau)$  be a periodic function with period  $\beta$ , prove that  $f(\tau) = \sum_{n \in \mathbb{Z}} \tilde{f}_{\omega_n} e^{-i\omega_n \tau}$  where  $\omega_n$  and where  $\tilde{f}_{\omega_n}$  will be given in terms of f.

### Problem 3—Gaussian integrals

**3.1)** We consider a > 0. Compute  $I(a, 0)^2$ , where

$$I(a,b) = \int_{-\infty}^{\infty} dx \, e^{-\frac{1}{2}ax^2 - bx}$$
(4)

*hint:* you may want to write  $I^2$  as a two-dimensional integral that can be computed in polar coordinates.

**3.2)** Compute I(a, b).

**3.3)** By taking derivates of *I*, compute

$$\int_{-\infty}^{\infty} dx \, x^2 e^{-\frac{1}{2}ax^2 - bx} \tag{5}$$

The rest of this problem counts as a *Graduate* problem. Let  $\mathbf{x} = (x_1, \ldots, x_n)$  and  $\mathbf{h} = (h_1, \ldots, h_N)$  be *n*-component vectors. We define

$$Z(\mathbf{h}) = \int d\mathbf{x} e^{-\frac{1}{2}x_i \Gamma_{ij} x_j + h_i x_i},\tag{6}$$

where  $\Gamma$  is, for now, a positive definite  $n \times n$  matrix. We use the notation  $\frac{1}{2}x_i\Gamma_{ij}x_j - h_ix_i$  for  $\frac{1}{2}\mathbf{x}\cdot(\Gamma x) - \mathbf{h}\cdot\mathbf{x}$ , *i.e.* we implicitly sum over repeated indices. We also defined  $p(\mathbf{x}) = \frac{1}{Z(\mathbf{0})}e^{-\frac{1}{2}\mathbf{x}\cdot(\Gamma\mathbf{x})}$  and the angular brackets mean  $\langle \ldots \rangle \equiv \int d\mathbf{x} \ldots P(\mathbf{x})$ .

**3.4)** Check that  $\langle e^{\mathbf{h}\cdot\mathbf{x}} \rangle = Z(\mathbf{h})/Z(\mathbf{0}).$ 

**3.5)** Why can we always restrict our analysis to the case where  $\Gamma$  is symmetric? This property will be assumed in the rest of this problem.

**3.6)** Prove that

$$\frac{1}{2}\mathbf{x}\cdot(\Gamma\mathbf{x}) - \mathbf{h}\cdot\mathbf{x} = \frac{1}{2}(\mathbf{x}-\Gamma^{-1}\mathbf{h})\cdot[\Gamma(\mathbf{x}-\Gamma^{-1}\mathbf{h})] - \frac{1}{2}(\Gamma^{-1}\mathbf{h})\cdot[\Gamma(\Gamma^{-1}\mathbf{h})].$$
(7)

Show then that

$$\langle e^{\mathbf{h} \cdot \mathbf{x}} \rangle = e^{\frac{1}{2}\mathbf{h} \cdot (\Gamma^{-1}\mathbf{h})} . \tag{8}$$

This is the generalization of the computation of I(a, b) to N variables,  $b^2/2a$  has now became a matrix relation. **3.7)** With our choice for P above,  $P(\mathbf{x}) = P(-\mathbf{x})$ , so that  $\langle \mathbf{x} \rangle = 0$ . Consider instead the new probability density obtained by including the field term  $\mathbf{h} \cdot \mathbf{x}$ :  $\tilde{P}(\mathbf{x}) \propto e^{-\frac{1}{2}x_i\Gamma_{ij}x_j + h_ix_i}$ . Show, using symmetry consideration, that  $\langle \mathbf{x} \rangle = \Gamma^{-1}\mathbf{h}$ .

**3.8)** Since  $\Gamma$  is symmetric, it can be diagonalized into a matrix D such that  $\Gamma = QDQ^T$ , with  $Q^T = Q^{-1}$ . By changing variables from  $\mathbf{x}$  to  $\mathbf{y} = Q^{-1}\mathbf{x}$ , compute  $Z(\mathbf{0})$  to show that

$$\int dx e^{-\frac{1}{2}\mathbf{x}\cdot\Gamma\mathbf{x}} = \frac{(2\pi)^{N/2}}{\sqrt{\det\Gamma}} \tag{9}$$

This is the generalization of the derivation of I(a, 0) to N variables.

#### <u>Problem 4—Dirac distribution</u>

The Dirac distribution  $\delta(x)$  can be defined from

$$\int_{-\infty}^{\infty} f(x)\delta(x) = f(0) \tag{10}$$

**4.1)** Show that  $\delta(ax) = \frac{1}{|a|}\delta(x)$ .

**4.2)** Compute the Fourier transform of  $\delta(x)$ ,  $\hat{\delta}(k) = \int_{-\infty}^{\infty} dx \, e^{ikx} \delta(x)$  and show that

$$\delta(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk \, e^{-ikx} \tag{11}$$

**4.3)** Consider a probability density p(x) and define the average of an observable  $f(x, x_0)$  with respect to x as  $\langle f(x, x_0) \rangle_x = \int dx f(x, x_0) p(x)$ . Show that  $p(x) = \langle \delta(x - x_0) \rangle_{x_0}$ .

**4.4)** Consider two Gaussian random variables x and y of averages  $\bar{x}$  and  $\bar{y}$  and of variance  $\sigma_x^2$  and  $\sigma_y^2$ . Using the result of question 4.3, compute  $p(z = \alpha x + \beta y)$ . *Hint*: you may want to use Eq.(11) to turn  $\delta(z - z_0)$  into a more useful expression.

# Problem 5—Functional derivatives (Graduate)

Let q(t) be a function of t and let S[q] be a functional of q (i.e. an application for the space of function q(t) into the field of real or complex numbers). The functional derivative of S with respect to  $q(t_0)$  is defined as follows. Let  $q_{\epsilon,t_0}(t) = q(t) + \epsilon \delta(t - t_0)$ , then

$$\frac{\delta S}{\delta q(t_0)} = \lim_{\epsilon \to 0} \frac{1}{\epsilon} (S[q_{\epsilon, t_0}] - S[q]) .$$
(12)

An equivalent definition is to say that when  $q \to q + \delta q$  (meaning that the trajectory q(t) is perturbed by  $\delta q(t)$ ), the functional changes from S to  $S + \delta S$ , with

$$\delta S = \int \frac{\delta S}{\delta q(t')} \delta q(t') dt' \tag{13}$$

to first order in  $\delta q$ . This relation defines the functional derivative  $\delta S/\delta q(t')$ , which is a functional of q and a function of t'.



Figure 1: A profile separating high-density and low-density regions

**5.1)** Compute  $\frac{\delta q(t_1)}{\delta q(t_2)}$ .

**5.2)** If S can be written in the form  $S[q] = \int_0^\infty dt L(q(t), \dot{q}(t))$ , where L is a function of q(t) and  $\dot{q}(t)$ , prove that

$$\frac{\delta S}{\delta q(t_0)} = \frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} , \qquad (14)$$

were everything is evaluated at  $t = t_0$ .

**5.3)** Consider a free energy  $\mathcal{F}[\rho] = \int dx \Big( f(\rho(x)) + \frac{\kappa}{2} [\partial_x \rho(x)]^2 \Big)$ . Show that the free energy is extremalized (minimized, really), by a profile that satisfies

$$\kappa \partial_{xx} \rho(x) = f'(\rho(x)) \tag{15}$$

Show that, for the phase-separated profile shown in Fig. 1, the free-energy density f is equal in the coexisting phases, i.e.  $f(\rho_g) = f(\rho_\ell)$