

Erasing correlations, destroying entanglement and other new challenges for quantum information theory

quant-ph/0511219

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QIP, 19 Jan 2006

outline

- General rules for reversing protocols
- Coherent erasure of classical correlations
- Disentangling power of quantum operations

Everything is a resource

qubit $[q \rightarrow q]$ $|0\rangle^A |0\rangle^B$ and $|1\rangle^A |1\rangle^B$

ebit $[qq]$ the state $(|0\rangle^A |0\rangle^B + |1\rangle^A |1\rangle^B) / \sqrt{2}$

~~cbit $[c \rightarrow c]$ $|0\rangle^A |0\rangle^B |0\rangle^E$ and $|1\rangle^A |1\rangle^B |1\rangle^E$~~

cobit $[q \rightarrow qq]$ $|0\rangle^A |0\rangle^A |0\rangle^B$ and $|1\rangle^A |1\rangle^A |1\rangle^B$

resource inequalities

super-dense coding:

$$[q \rightarrow q] + [qq] \geq 2[c \rightarrow c] = 2 [q \rightarrow qq]$$

In fact, $[q \rightarrow q] + [qq] = 2 [q \rightarrow qq]$

Undoing things is also a resource

reversal

meaning

$$[q! \ q]^y = [q\tilde{A}q]$$

(relation between time-reversal and exchange symmetry)

$$[qq]^y = -[qq]$$

(disentangling power)

$$[q! \ qq]^y = [q\tilde{A}qq] (?)$$

$|0i^A \ 0i^B\rangle \ |0i^A \ \text{and} \ 1i^A \ 1i^B\rangle \ |1i^A$
(coherent erasure??)

What good is coherent erasure?

$$\alpha|0\rangle^A + \beta|1\rangle^A \quad \alpha|0\rangle^A|0\rangle^B + \beta|1\rangle^A|1\rangle^B \quad (\text{using } [q \rightarrow qq])$$

$$\quad \rightarrow \quad \alpha|0\rangle^B + \beta|1\rangle^B \quad (\text{using } [qq \rightarrow q])$$

$$[q \rightarrow qq] + [qq \rightarrow q] \overset{=}{\cancel{[q \rightarrow q]}}$$

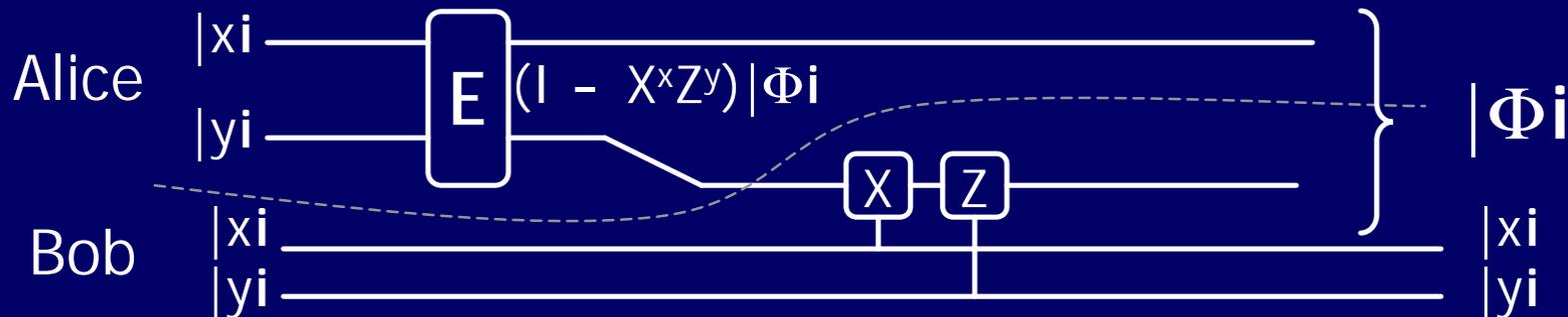
$$[qq \rightarrow q] \overset{=}{\cancel{[q \rightarrow q]}} - [q \rightarrow qq]$$

$$= [q \rightarrow qq] - [qq]$$

$$= ([q \rightarrow q] - [qq]) / 2$$

} entanglement-assisted communication only

In fact, these are all equalities! (Proof: reverse SDC.)



application to unitary gates

U is a bipartite unitary gate (e.g. CNOT)

Known:

$$U \succ C[c! \ c] \text{ implies } U \succ C[q! \ qq]$$

Time reversal means:

$$\begin{aligned} U^y &\succ C[q\tilde{A}qq] \\ &= C[qq\tilde{A}q] - C[qq] \end{aligned}$$

Corollary: If entanglement is free then $C_!^E(U) = C_{\tilde{A}}^E(U^y)$.

The quest for asymmetric unitary gate capacities

Problem: If U is nonlocal, it has nonzero quantum capacities in both directions. Are they equal?

Yes, if U is 2×2 .

No, in general, but for a dramatic separation we will need a gate that violates time-reversal symmetry.

the construction:

(U_m acts on $2^m \times 2^m$ dimensions)

$$U_m |x\rangle^A |0\rangle^B = |x\rangle^A |x\rangle^B \quad \text{for } 0 \leq x < 2^m$$

$$U_m |x\rangle^A |y\rangle^B = |x\rangle^A |y-1\rangle^B \quad \text{for } 0 < y \leq x < 2^m$$

$$U_m |x\rangle^A |y\rangle^B = |x\rangle^A |y\rangle^B \quad \text{for } 0 \leq x < y < 2^m$$

et voilà l'asymétrie!

$$U_m |x i^A| 0 i^B = |x i^A| x i^B \quad \text{for } 0 \leq x < 2^m$$

$$U_m |x i^A| y i^B = |x i^A| y^{-1} i^B \quad \text{for } 0 < y \leq x < 2^m$$

$$U_m |x i^A| y i^B = |x i^A| y i^B \quad \text{for } 0 \leq x < y < 2^m$$

$$U_m > m [q! \quad qq]$$

Upper bound by simulation:

$$m [q! \quad qq] + O((\log m)(\log m/\varepsilon)) ([q! \quad q] + [q \tilde{A} q]) \& U_m$$

Similarly, $U_m^y > m [q \tilde{A} qq]$ and

$$m [q \tilde{A} qq] + O((\log m)(\log m/\varepsilon)) ([q! \quad q] + [q \tilde{A} q]) \& U_m^y$$

Meaning: $U_m \approx \frac{1}{4} m [q! \quad qq]$

and $U_m^y \approx \frac{1}{4} m [q \tilde{A} qq]$ (almost worthless w/o ent. assistance!)

disentanglement

clean resource inequalities:

$$\alpha \stackrel{\text{clean}}{\geq} \beta$$

means that α^{-n} can be asymptotically converted to β^{-n} while discarding only $o(n)$ entanglement.

(equivalently: while generating a sublinear amount of local entropy.)

Example: $[q! q] \stackrel{\text{clean}}{>} [qq]$ and $[q! q] \stackrel{\text{clean}}{>} -[qq]$

Example: $U_m \stackrel{\text{clean}}{>} m[qq]$, but can only destroy $O(\log^2 m)$ $[qq]$

$U_m^y \stackrel{\text{clean}}{>} -m[qq]$, but can only create $O(\log^2 m)$ $[qq]$

You can't just throw it away

Q: Why not?

A: Given unlimited EPR pairs, try creating the state

$$\frac{1}{\sqrt{2}} \left(|00\rangle_{AB}^{\otimes n} + |\Phi_{-}\rangle_{AB}^{\otimes n} \right)$$

Hayden & Winter [quant-ph/0204092] proved that this requires $\frac{1}{4}n$ bits of communication.

more relevant examples

Entanglement dilution:

$|\psi\rangle_{AB}$ is partially entangled. $E = S(\psi^A)$.

$|\Phi\rangle_{nE+o(n)} \rightarrow |\psi\rangle_{nE}$

Even $|\Phi\rangle_{nE-1} \rightarrow |\psi\rangle_{nE}$ requires $\Omega(n^{1/2})$ cbits (in either direction).

OR a size $O(n^{1/2}/\epsilon)$ embezzling state [q-ph/0205100, Hayden-van Dam]

Quantum Reverse Shannon Theorem for general inputs

[Bennett,
Devetak,
Harrow,
Shor,
Winter]

Input ρ_{AB} requires $I(A;B)_\rho$ [cbits] + $H(N(\rho))$ [qq].

Superpositions of different ρ_{AB} mean consuming superpositions of different amounts of entanglement: we need either extra cbits, embezzling, or another source of **entanglement spread**.

summary

- new ideas

- coherent erasure
- clean protocols
- entanglement spread

READ ALL ABOUT IT!

[quant-ph/0511219](#)

- new results

- asymmetric unitary gate capacities
- QRST and other converses

- new directions

- formalizing entanglement spread
- clean protocols involving noisy resources (cbits?)