

(1)

Hermite:  $A = \text{randn}(n)$ ;  $\text{eig}(A + A^*)/2$

Laguerre:  $A = \text{randn}(m, n)$ ;  $\text{svd}(A)$

Jacobi:  $A = \text{randn}(n_1, p)$   
 $B = \text{randn}(n_2, p)$   $\text{gsvd}(A, B) \leftarrow \text{rot form}$

OR

$X = \text{randn}(n, p)$  ( $n = n_1 + n_2$ )

$[Y, U] = \text{qr}(X, U)$

$\text{svd}(Y(1:n_1, :)) \leftarrow \text{cosine form}$

Complex Versions:  $A = \text{randn}(n) + i * \text{randn}(n)$  etc.

Quaternion Versions:

Want:  $A = \text{randn}(n) + i * \text{randn}(n) + j * \text{randn}(n) + k * \text{randn}(n)$ ;

Workaround:  $X = \text{randn}(n) + i * \text{randn}(n)$ ;

$Y = \text{randn}(n) + i * \text{randn}(n)$ ;

$A = \begin{bmatrix} X & Y \\ -\bar{Y} & \bar{X} \end{bmatrix}$  "conj"

- These are inefficient computationally
- The better way computationally is also better theoretically in so many ways
- Began a whole study of Stochastic Differential Equations still with many open problems
- Began a whole study of beta ensembles

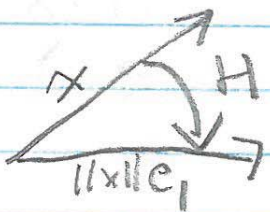
(2)

Outline: Numerical Linear Algebra: Sym eig problem  
orthogonal invariance for chi-squared  
real beta<sup>(2)</sup> ensembles  
general beta ensembles

1. Crash course on eigenvalue computation for sym matrices  
Characteristic Polynomial never used (too slow, too inaccurate)

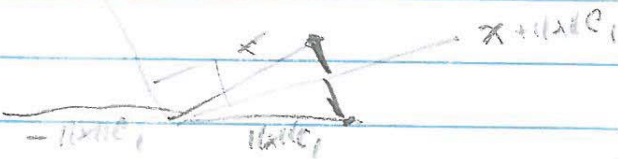
Sym Dense  $\xrightarrow[\text{finite, slow}]{\text{similar}}$  Sym Tridiagonal  $\xrightarrow[\text{infinite, fast}]{\text{similar}}$  Diagonal

→ Eigenvalues preserved



$H =$  Householder Reflector  
 $Hx = \|x\|e_1$       $e_1 = \begin{pmatrix} 1 \\ 0 \\ \vdots \end{pmatrix}$

$$H = I - \frac{2uu^T}{u^Tu} \quad u = x - \|x\|e_1$$



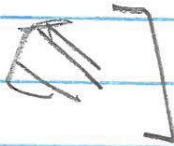
$I - \frac{u u^T}{u^T u}$  Projects into  $x \cos(\alpha) e_1$   
 $I - \frac{2 u u^T}{u^T u}$  Reflects



(3)

$$\begin{bmatrix} 1 & 0 \\ 0 & H(x) \end{bmatrix} \begin{bmatrix} \alpha & x^T \\ x & A \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & H(x) \end{bmatrix}$$

$$= \begin{bmatrix} \alpha & \|x\| & 0 \\ \|x\| & & \\ 0 & & A' \end{bmatrix}$$

Recur on  $A'$  together 

2)  $Q \cdot \text{rand}_n(n, 1) \equiv \text{rand}_n(n, 1)$  for any fixed  $Q$   
 $\text{norm}(\text{rand}_n(n, 1)) \equiv \chi_n$

Example:  $[z, R] = q(\text{rand}_n(n, 1))$

$$x = \text{rand}_n(n, 1)$$
$$H(x) \cdot x = \begin{bmatrix} \chi_n \\ \vdots \\ \vdots \end{bmatrix}$$

④

$$\begin{bmatrix} G & G & G & G \\ G & G & G & G \\ G & G & G & G \\ G & G & G & G \end{bmatrix} \rightarrow \left[ \begin{array}{c|ccc} \chi_4 & G & G & G \\ 0 & G & G & G \\ 0 & G & G & G \\ 0 & G & G & G \end{array} \right] \text{ etc}$$

$$|A| \sim \chi_n \chi_{n-1} \dots \chi_1$$

$$\mathbb{E}|A|^2 = n! \quad (\text{True more generally})$$

### Real Beta Ensembles

$$\left. \begin{array}{l} \text{Triagonalize } (A+A^T) \\ \text{diag. var} = 1 \\ \text{off diag. var} = \frac{1}{2} \end{array} \right\} \begin{array}{l} \text{diag: } \text{randn}(n, 1) \\ \text{off-diag: } \frac{1}{\sqrt{2}} [\chi_{n-1}, \chi_{n-2}, \dots, \chi_1] \end{array}$$

$$H_\beta \sim \frac{1}{\sqrt{2}} \begin{pmatrix} N(0, 2) & \chi_{n-1} \\ \chi_{n-1} & N(0, 2) \\ & & N(0, 2) \end{pmatrix} \quad \beta = 1$$

(5)

Theorem: The joint eigenvalue density of  $H_1$  is the same as the GOE, GUE, GSE namely

$$C_\beta \prod_{i < j} |\lambda_i - \lambda_j|^\beta e^{-\sum_{i=1}^n \frac{\lambda_i^2}{2}} \quad \text{when } \beta = 1, 2, 4$$

$$C_\beta = \frac{2^n \pi^{n^2/2}}{\prod_{j=1}^n \Gamma(1 + \frac{\beta}{2})}$$

These joint densities are already known & remain unchanged.

Theorem: For all  $\beta > 0$ , the joint density is correct.

Proof:

Lemma:  $T = \begin{pmatrix} a_1 & b_1 & & \\ & \ddots & & \\ & & a_n & \\ & & & b_n \end{pmatrix}$

$$Q^T T Q = \Lambda \quad \begin{matrix} \lambda_{11} & \dots & \lambda_{nn} \\ + q \end{matrix}$$
$$q = \varphi(1, i)$$



*Proof.* To obtain the Jacobian, we will study the transformation from GOE to 1-Hermitic ensemble (see Figure 2). Note that  $J$  does *not* depend on  $\beta$ ; hence computing the Jacobian for this case is sufficient.

Let  $T$  be a 1-Hermitic matrix. We know from Section II.A that the eigenvalues of  $T$  are distributed as the eigenvalues of a symmetric GOE matrix  $A$ , from which  $T$  can be obtained via tridiagonal reduction ( $T = HAH^T$  for some orthogonal  $H$ , which is the product of the consecutive reflections described in Section II.A).

The joint element distribution for the matrix  $T$  is

$$\mu(a, b) = c_{a,b} e^{-\frac{1}{2} \sum_{i=1}^n a_i^2} \prod_{i=1}^n b_i^{i-1} e^{-\sum_{i=1}^n b_i^2}, \quad \text{where} \quad c_{a,b} = \frac{2^{n-1}}{(2\pi)^{n/2} \prod_{i=1}^{n-1} \Gamma(\frac{i}{2})}.$$

Let

$$da = \wedge_{i=1}^n da_i, \quad db = \wedge_{i=1}^{n-1} db_i, \quad d\lambda = \wedge_{i=1}^n \lambda_i,$$

and  $dq$  be the surface element of the  $n$ -dimensional sphere. Let  $\mu(a(q, \lambda), b(q, \lambda))$  be the expression for  $\mu(a, b)$  in the new variables  $q, \lambda$ . We have that

$$\mu(a, b) da db = J \mu(a(q, \lambda), b(q, \lambda)) dq d\lambda \equiv \nu(q, \lambda) dq d\lambda. \quad (17)$$

We combine Properties 1 and 2 of Section II.A to get the joint p.d.f.  $\nu(q, \lambda)$  of the eigenvalues and first eigenvector row of a GOE matrix, and rewrite it as

$$\nu(q, \lambda) dq d\lambda = n! c_H^1 \frac{2^{n-1} \Gamma(\frac{n}{2})}{\pi^{n/2}} \Delta(\lambda) e^{-\frac{1}{2} \sum_i \lambda_i^2} dq d\lambda.$$

We have introduced the  $n!$  and removed the absolute value from the Vandermonde, because the eigenvalues are ordered. We have also included the distribution of  $q$  (as mentioned in Property 2, it is uniform, but only on the all-positive  $2^{-n}$ th of the sphere because of the condition  $q_i \geq 0$ ).

Since orthogonal transformations do not change the Frobenius norm  $\|A\|_F = \sum_{i,j=1}^n a_{ij}^2$  of a matrix  $A$ , from (17), it follows that

$$J = \frac{\nu(q, \lambda)}{\mu(a, b)} = \frac{n! c_H^1 \frac{2^{n-1} \Gamma(\frac{n}{2})}{\pi^{n/2}} \Delta(\lambda)}{c_{a,b} \prod_{i=1}^n b_i^{i-1}}.$$

All constants cancel, and by Lemma 2.7 we obtain

$$J = \frac{\prod_{i=1}^{n-1} b_i}{\prod_{i=1}^n q_i}.$$

Note that we have not expressed  $\mu(a, b)$  in terms of  $q$  and  $\lambda$  in the above, and have thus obtained the expression for the Jacobian neither in the variables  $q$  and  $\lambda$ , nor  $a$  and  $b$ , solely; but rather in a mixture of the two sets of variables. The reason for this is that of simplicity.  $\square$